



McDONNELL CENTER
FOR THE SPACE SCIENCES

Zee-Burst: Non-Standard Interactions in IceCube

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In collaboration with
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Sudip Jana (MPI)
arXiv:1908.02779



Scalars in the Zee model

A. Zee Phys. Lett.95B,461(1980)



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$$(SU(3) \times SU(2)_L \times U(1)_Y)$$

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$$H_1 = \begin{pmatrix} G^+ \\ \frac{1}{\sqrt{2}}(v + H_1^0 + iG^0) \end{pmatrix}, \quad H_2 = \begin{pmatrix} H_2^+ \\ \frac{1}{\sqrt{2}}(H_2^0 + iA) \end{pmatrix}$$

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Due to the structure of scalar potential, H_2^+ will mix with η^+

$$\begin{aligned} h^+ & = \cos \varphi \eta^+ + \sin \varphi H_2^+, \\ H^+ & = -\sin \varphi \eta^+ + \cos \varphi H_2^+ \end{aligned}$$

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As for the Yukawa sector, we have:

$$-\mathcal{L}_Y \supset f_{\alpha\beta} L_\alpha^i L_\beta^j \epsilon_{ij} \eta^+ + \tilde{Y}_{\alpha\beta} \tilde{H}_1^i L_\alpha^j \ell_\beta^c \epsilon_{ij} + Y_{\alpha\beta} \tilde{H}_2^i L_\alpha^j \ell_\beta^c \epsilon_{ij} + \text{H.c.}$$



Neutrino Mass

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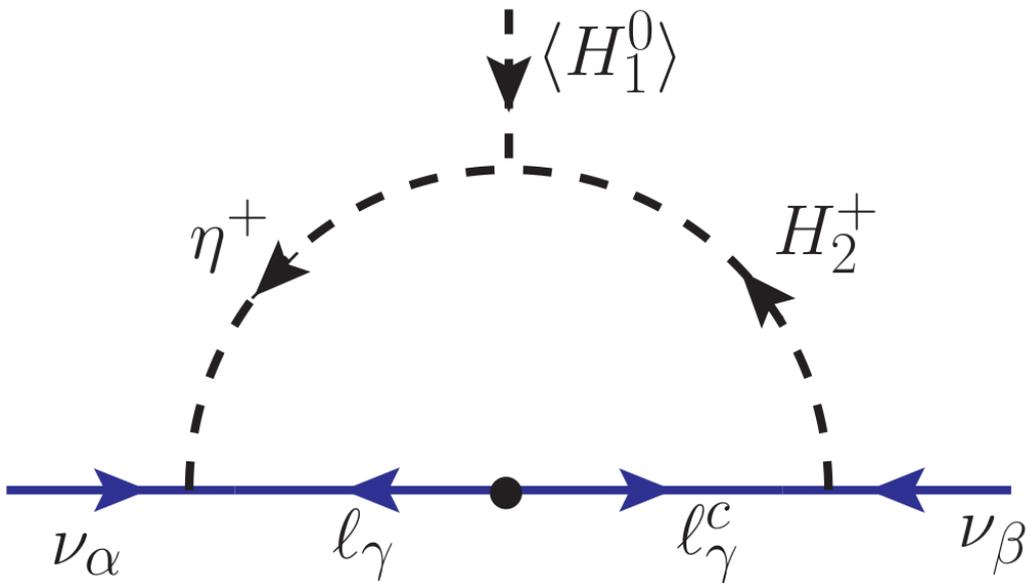
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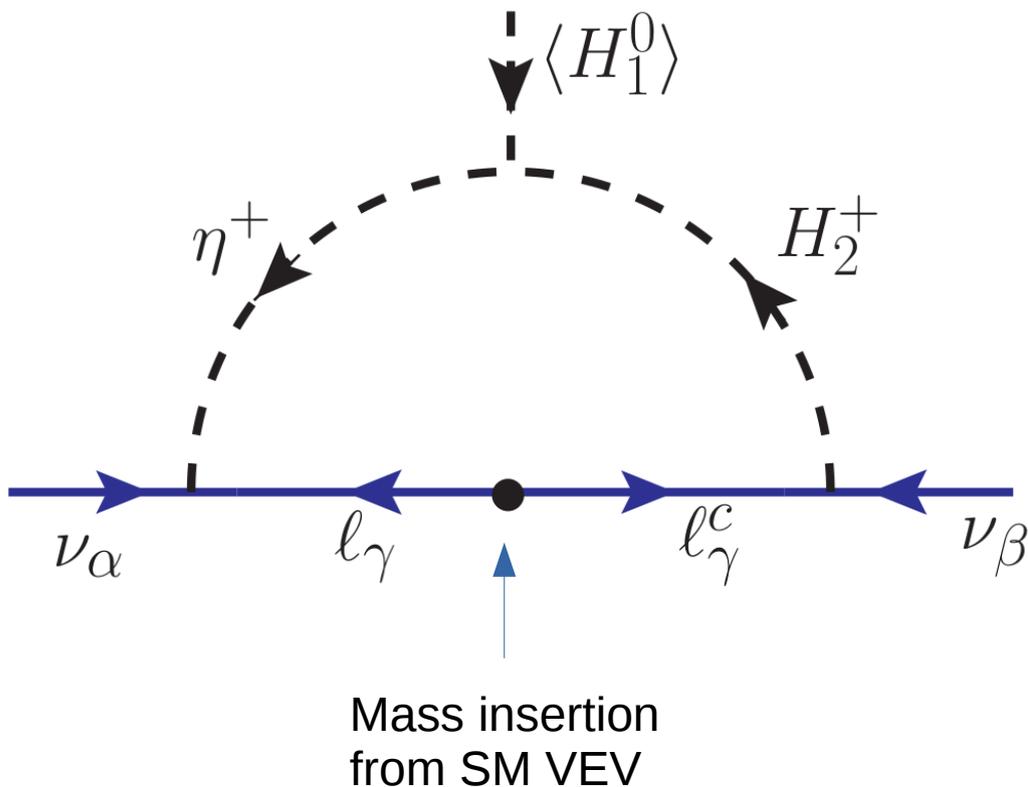
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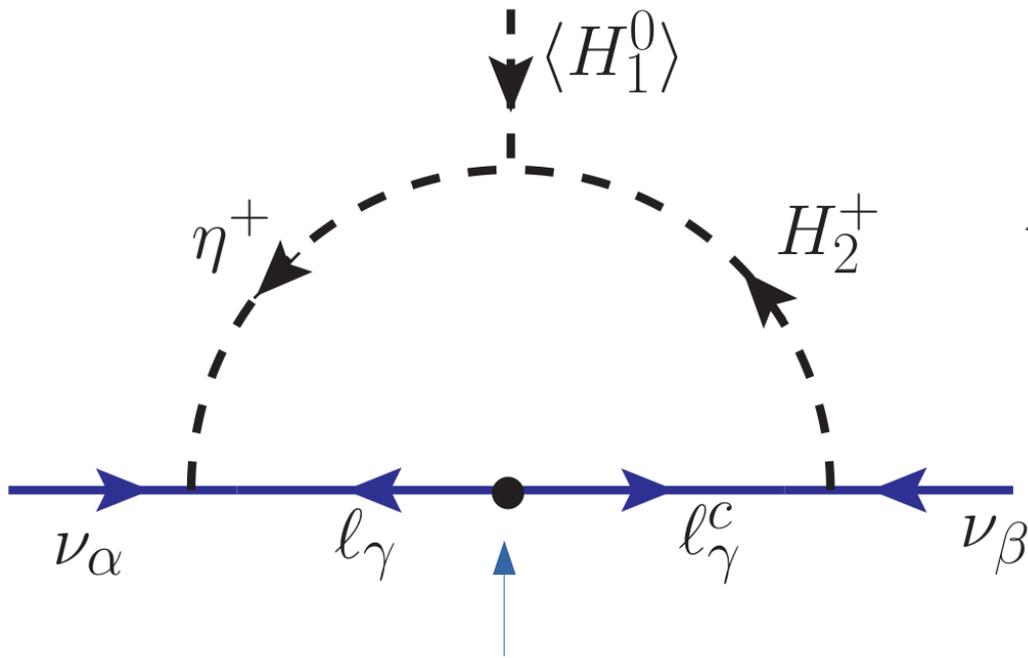
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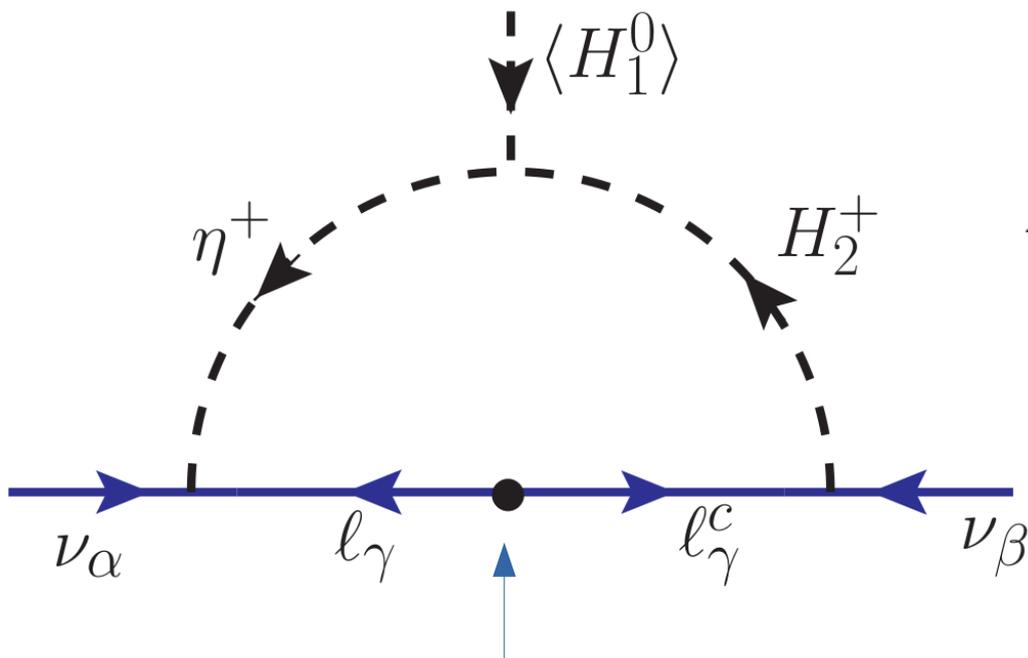
$$M_\nu = \kappa (f M_\ell Y + Y^T M_\ell f^T)$$

Mass insertion
from SM VEV

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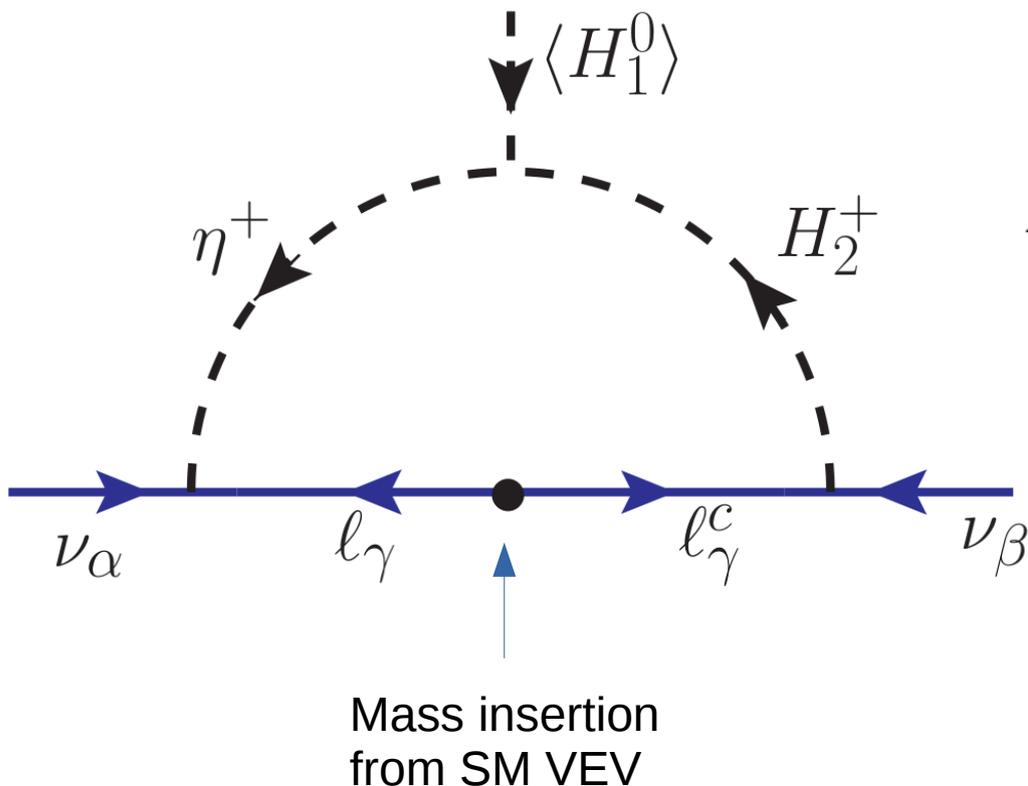
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Charged Lepton
Mass Matrix

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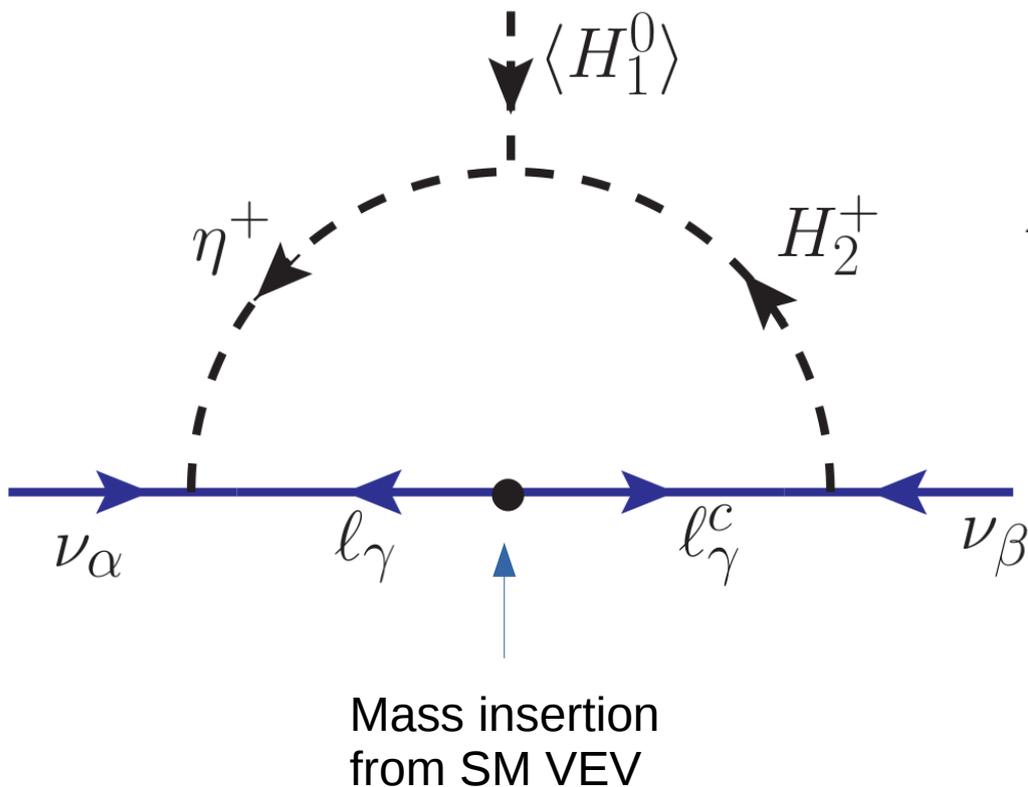
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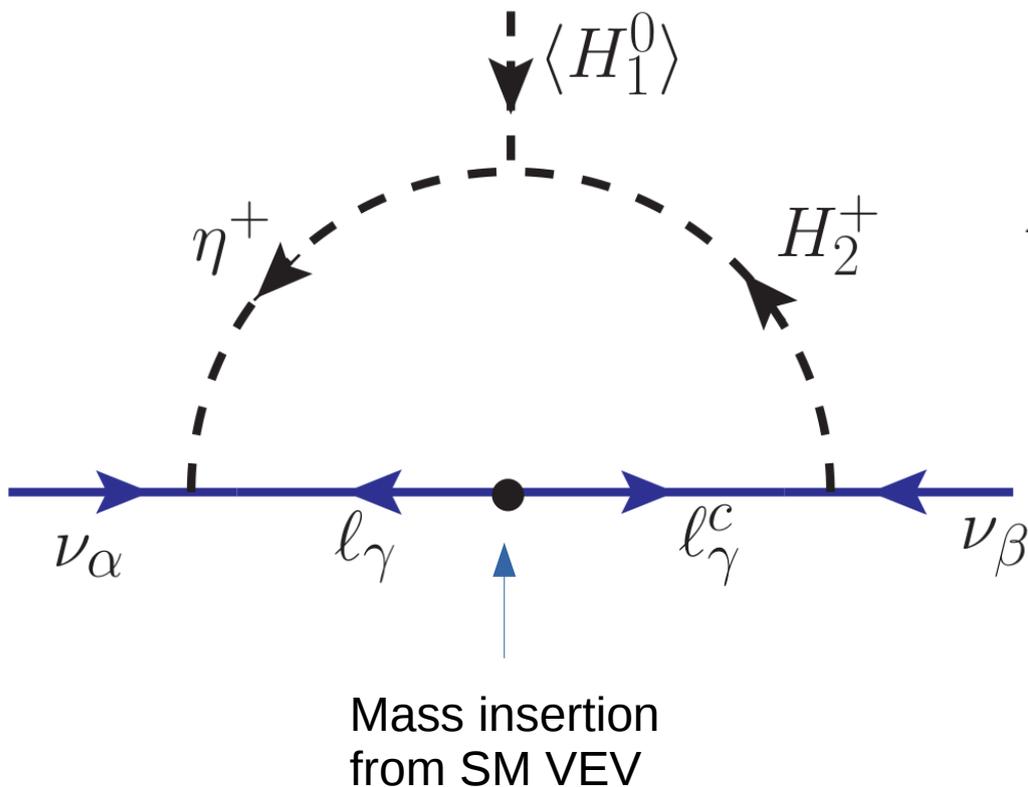
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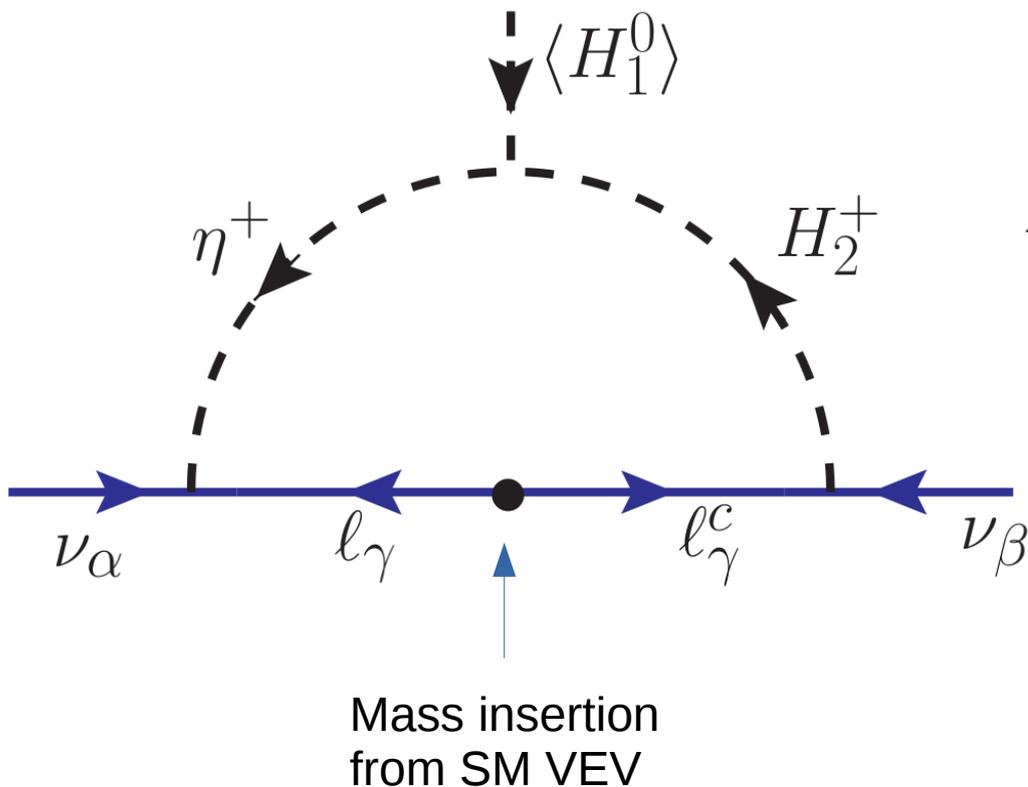
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Herrero-Garcia, Ohlsson, Riad, Wiren, 2017'



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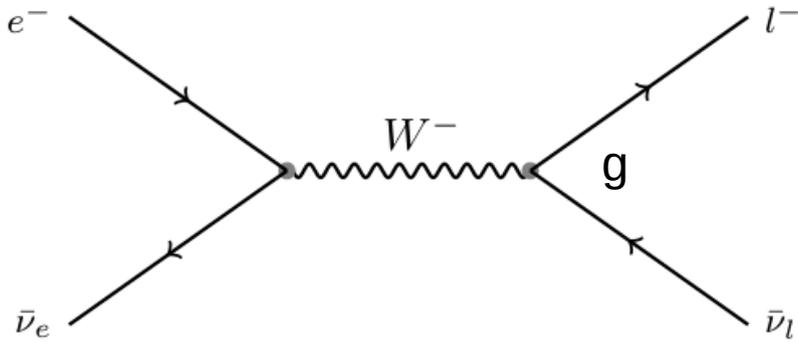
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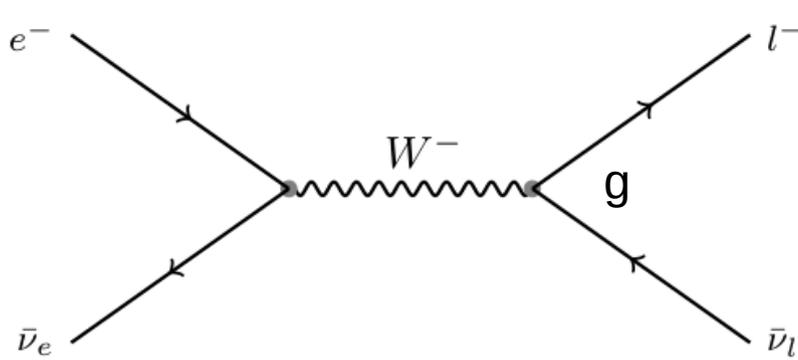
Glashow-Like Signatures

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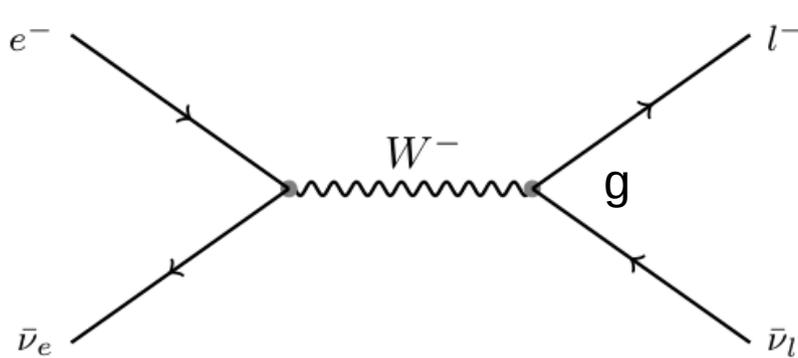
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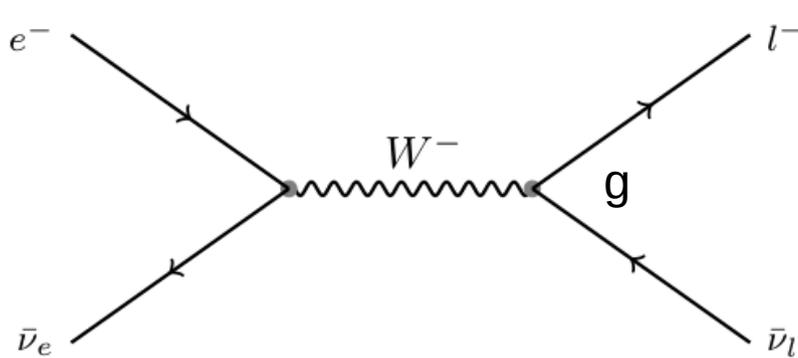


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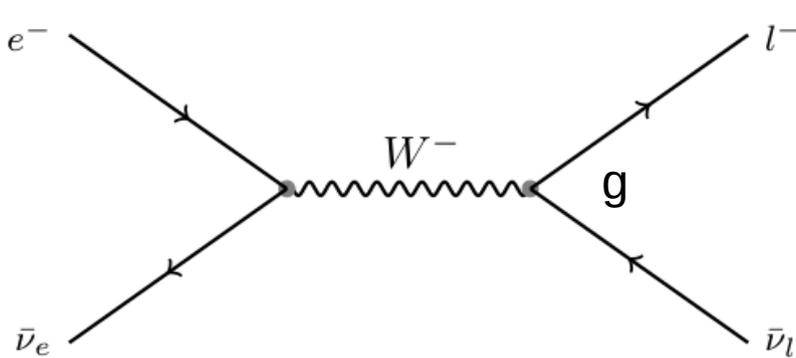


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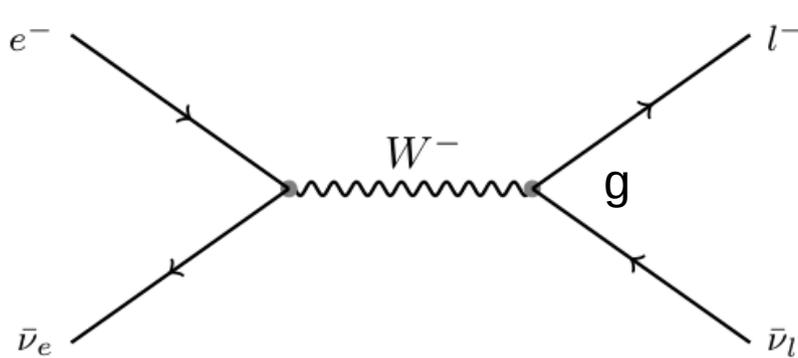
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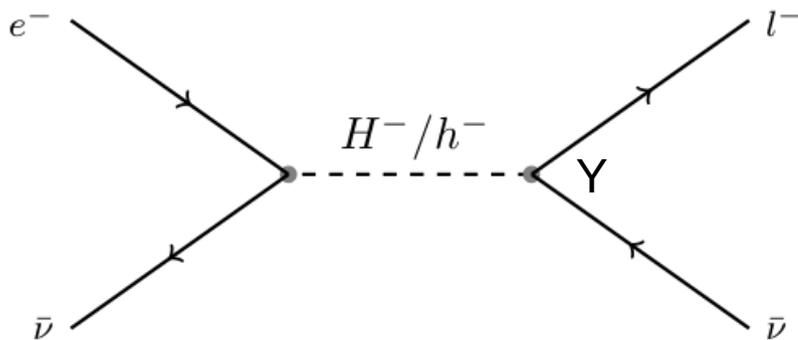


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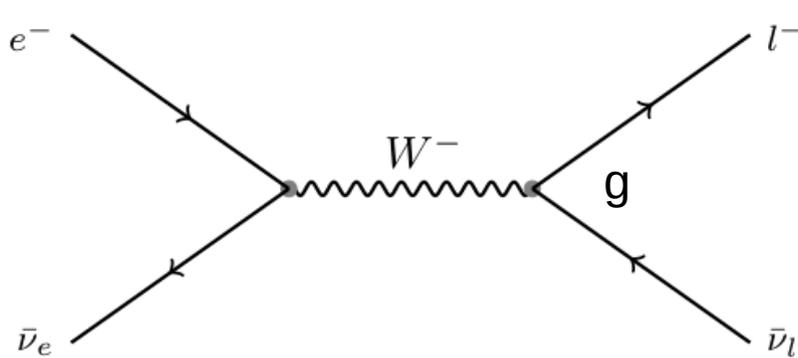
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Zee burst

Glashow-Like Signatures

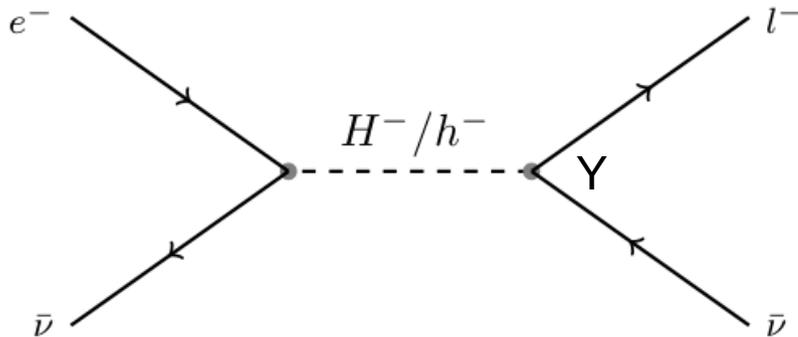


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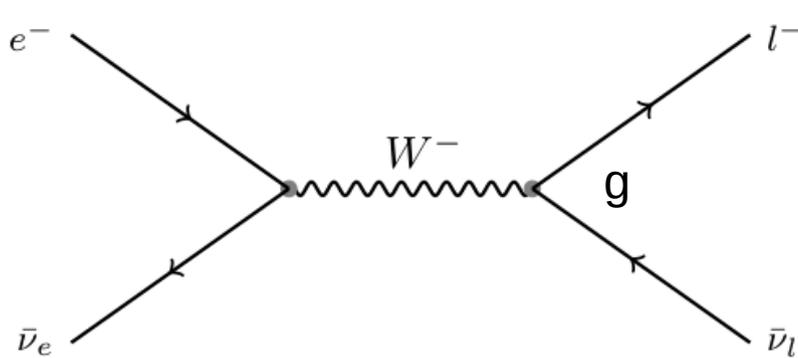
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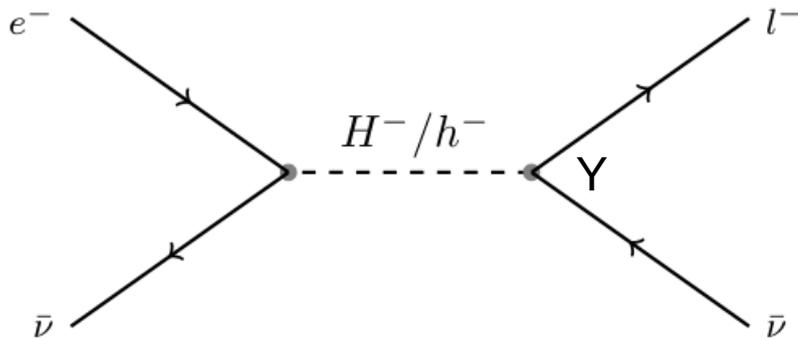


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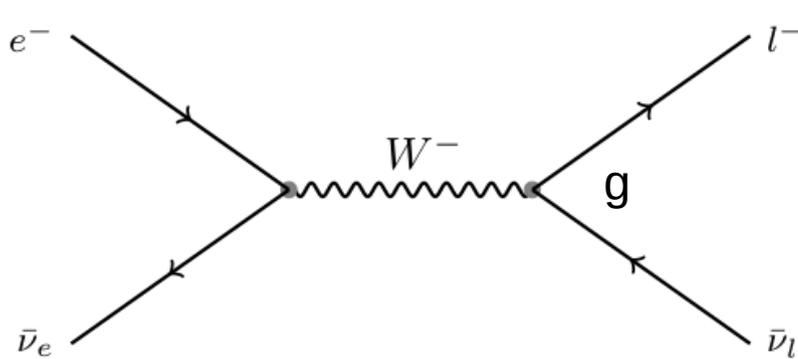


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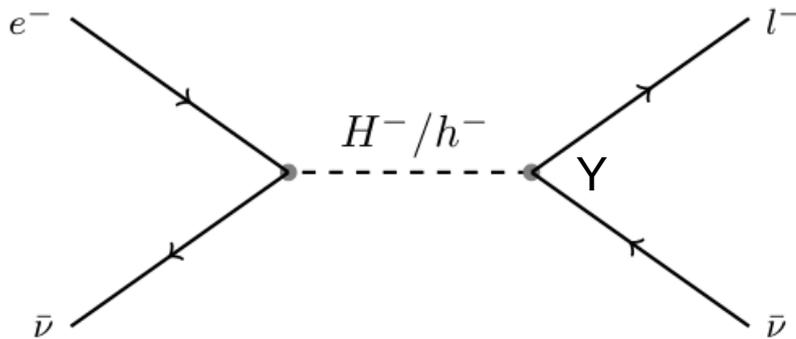


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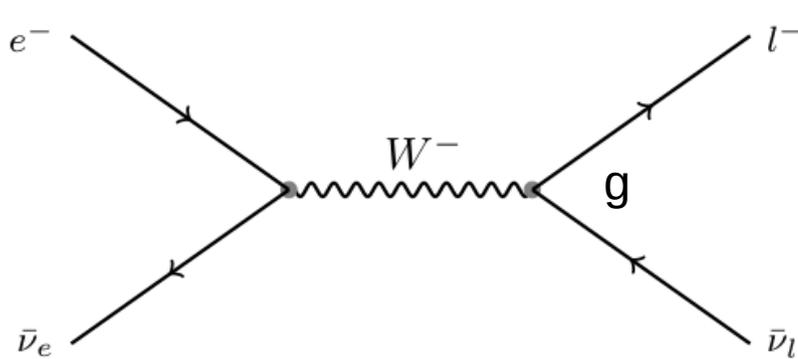
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Glashow-Like Signatures

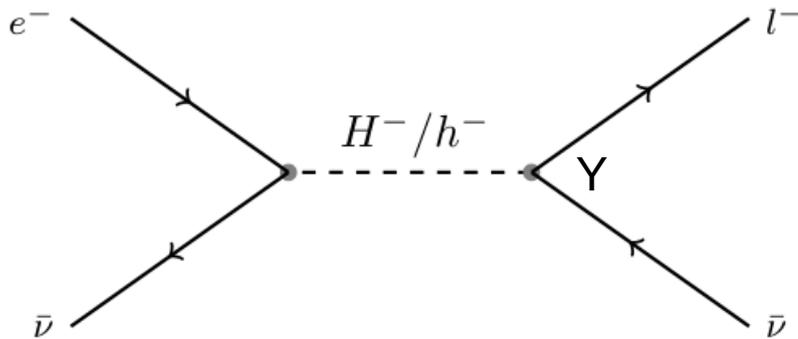


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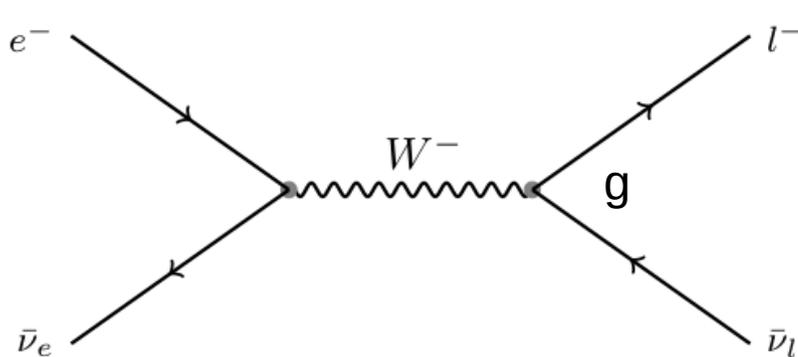
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m=80.4 GeV

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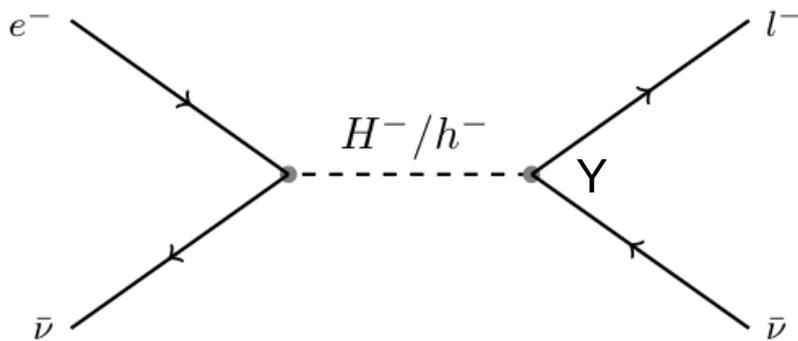


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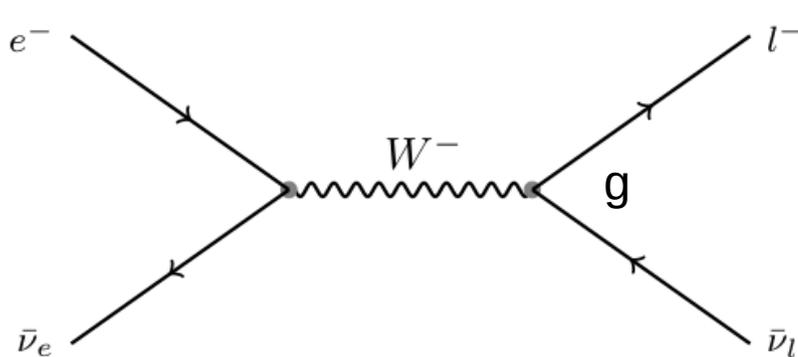
$$\Gamma_X = \sum_{\alpha\beta} |Y_{\alpha\beta}|^2 \sin^2 \varphi m_X / 16\pi$$

$$E_\nu = m^2 / 2m_e \approx 6.3 \text{ PeV}, 9.8 \text{ PeV}$$

m=80.4 GeV

m=100 GeV

Glashow-Like Signatures

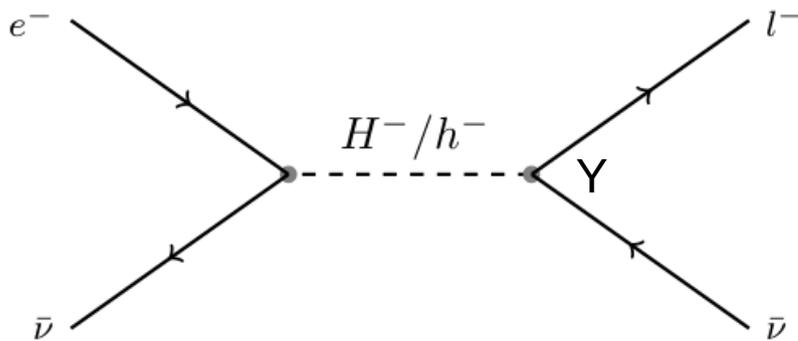


$$\sigma_{\text{Glashow}}(s) \sim \Gamma_W^2 \text{BR}(W^- \rightarrow \bar{\nu}_e e^-) \text{BR}(W^- \rightarrow \text{All}) \times \frac{s/m_W^2}{(s - m_W^2)^2 + (m_W \Gamma_W)^2}$$

@ resonance, becomes dominant

S. L. Glashow 1960

$$-\mathcal{L}_Y \supset f_{\alpha\beta} L_\alpha^i L_\beta^j \epsilon_{ij} \eta^+ + \tilde{Y}_{\alpha\beta} \tilde{H}_1^i L_\alpha^j \ell_\beta^c \epsilon_{ij} + Y_{\alpha\beta} \tilde{H}_2^i L_\alpha^j \ell_\beta^c \epsilon_{ij} + \text{H.c.}$$



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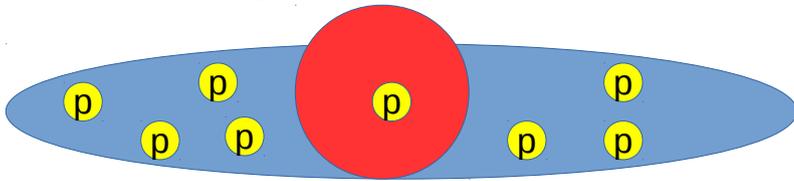
Where to find these High Energy neutrinos?



Astrophysical Neutrino Sources

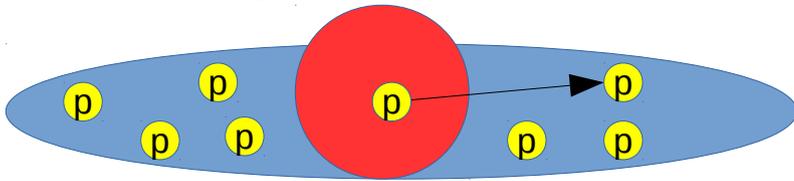
Astrophysical Neutrino Sources

hadro-nuclear
production



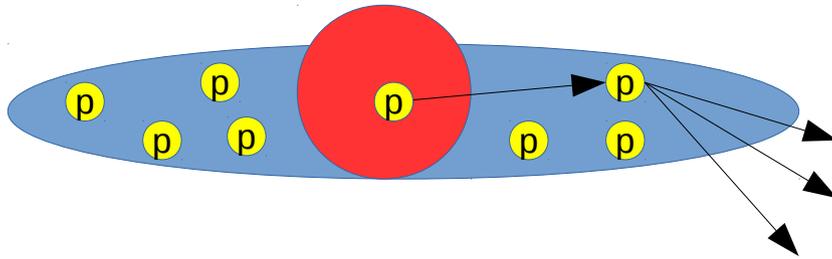
Astrophysical Neutrino Sources

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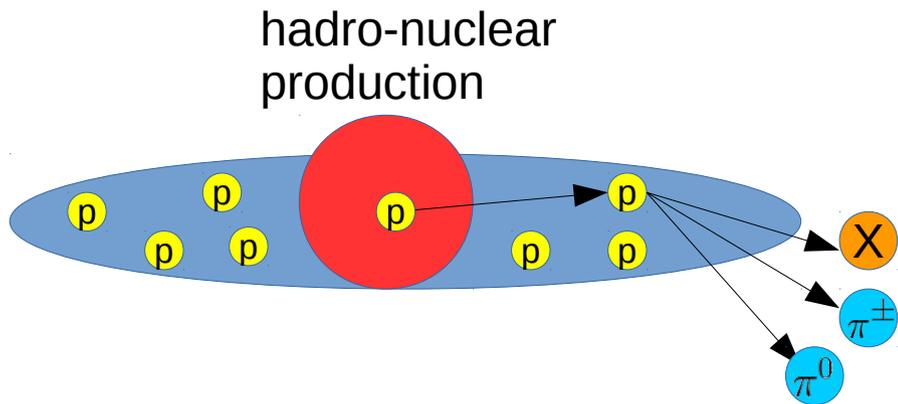


Astrophysical Neutrino Sources

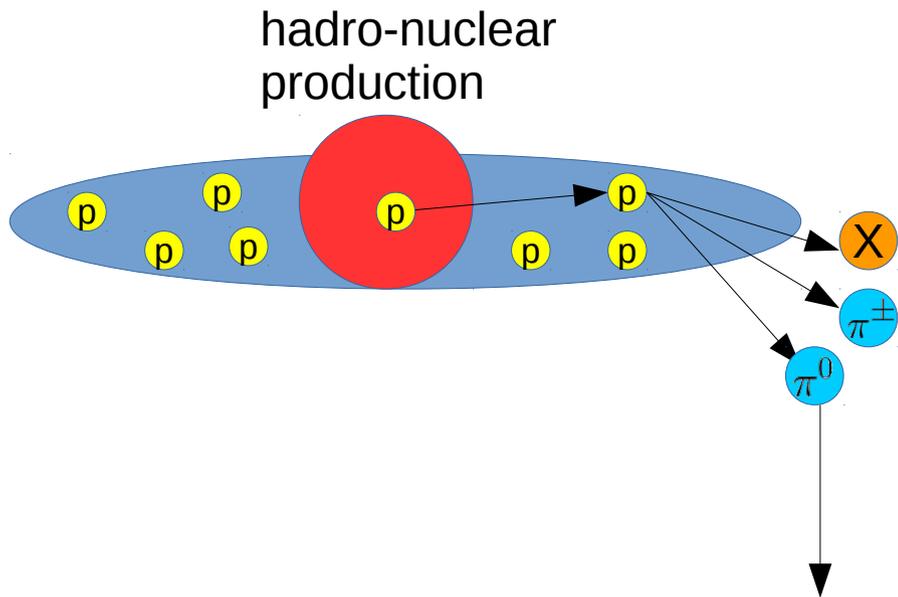
hadro-nuclear
production



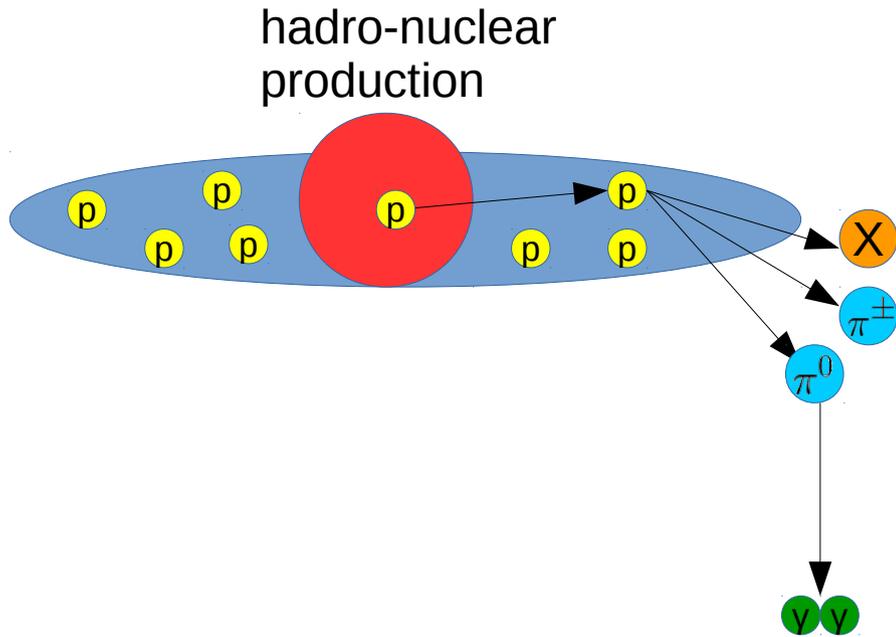
Astrophysical Neutrino Sources



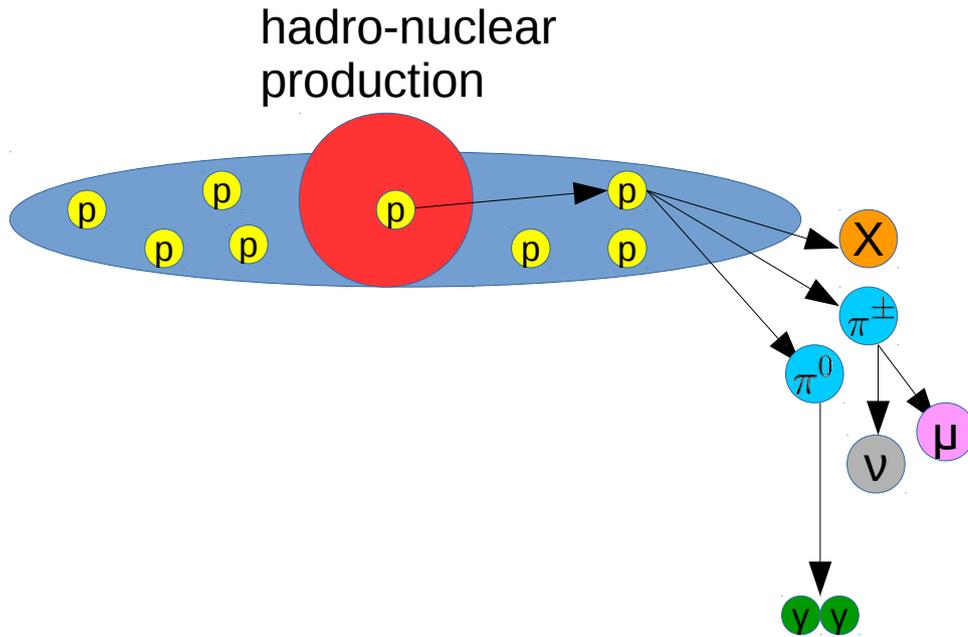
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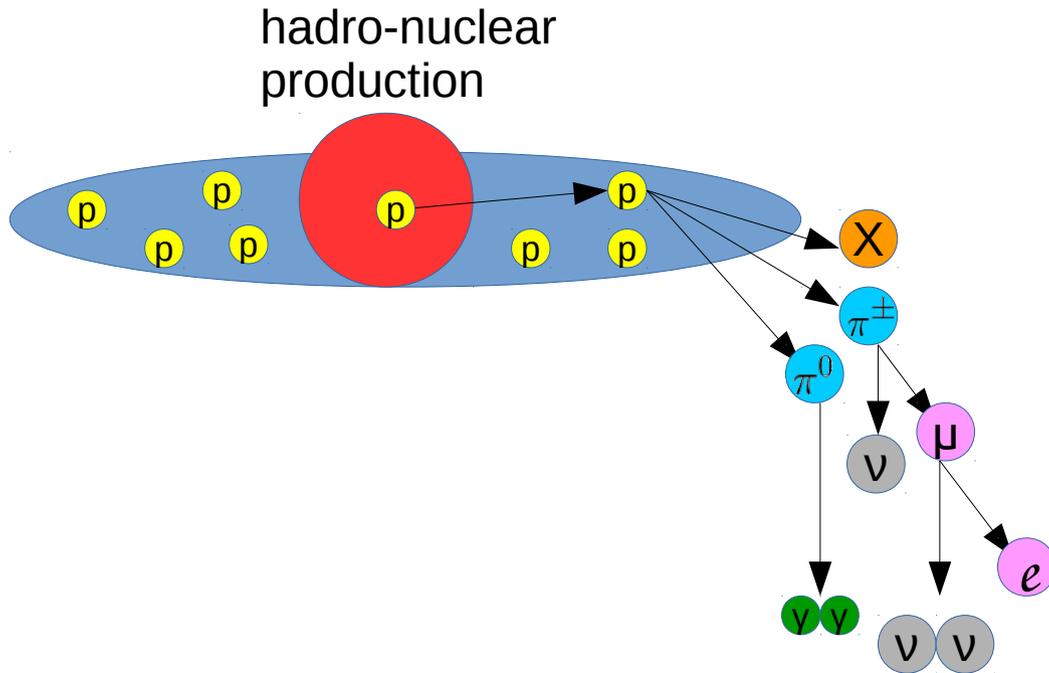
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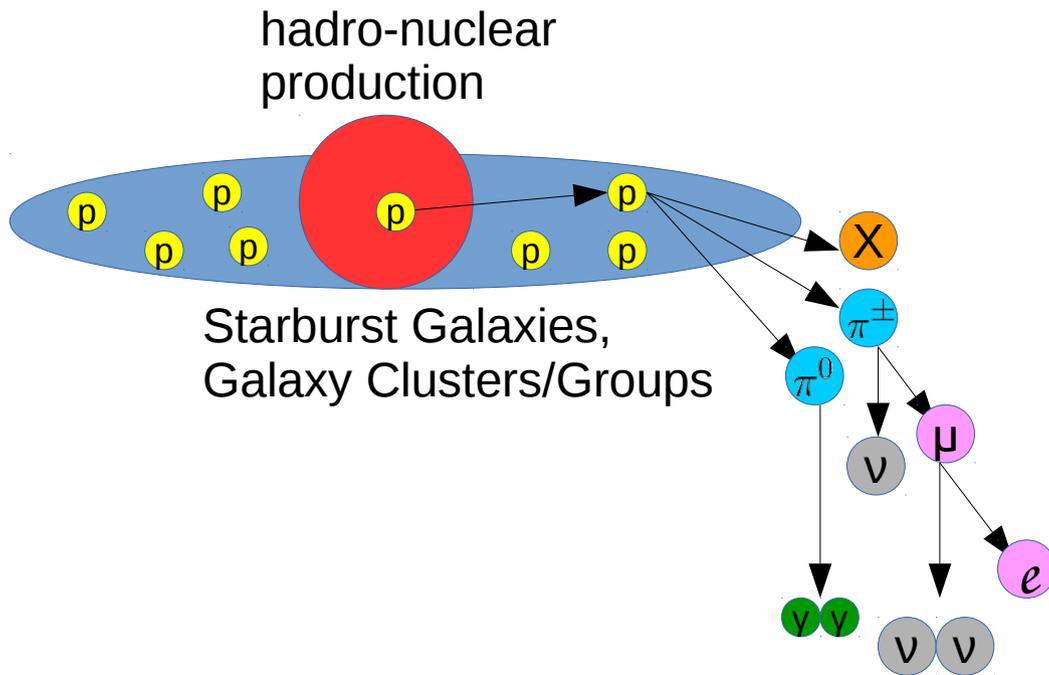
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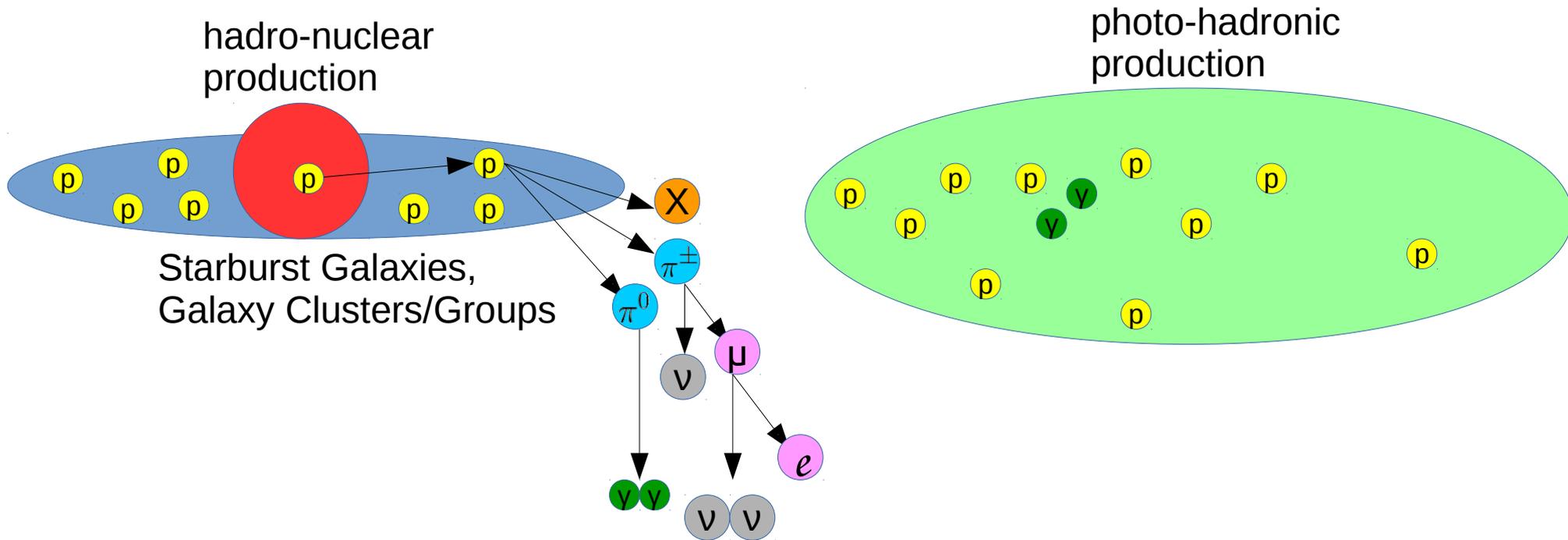
Astrophysical Neutrino Sources



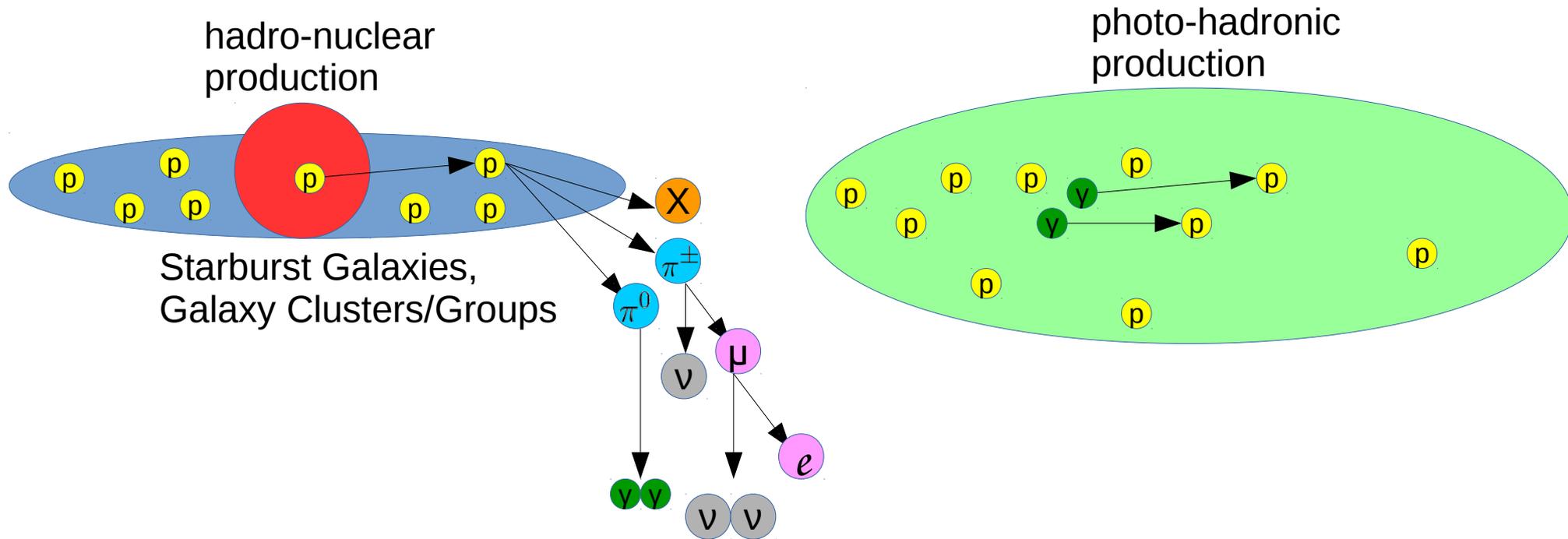
Astrophysical Neutrino Sources



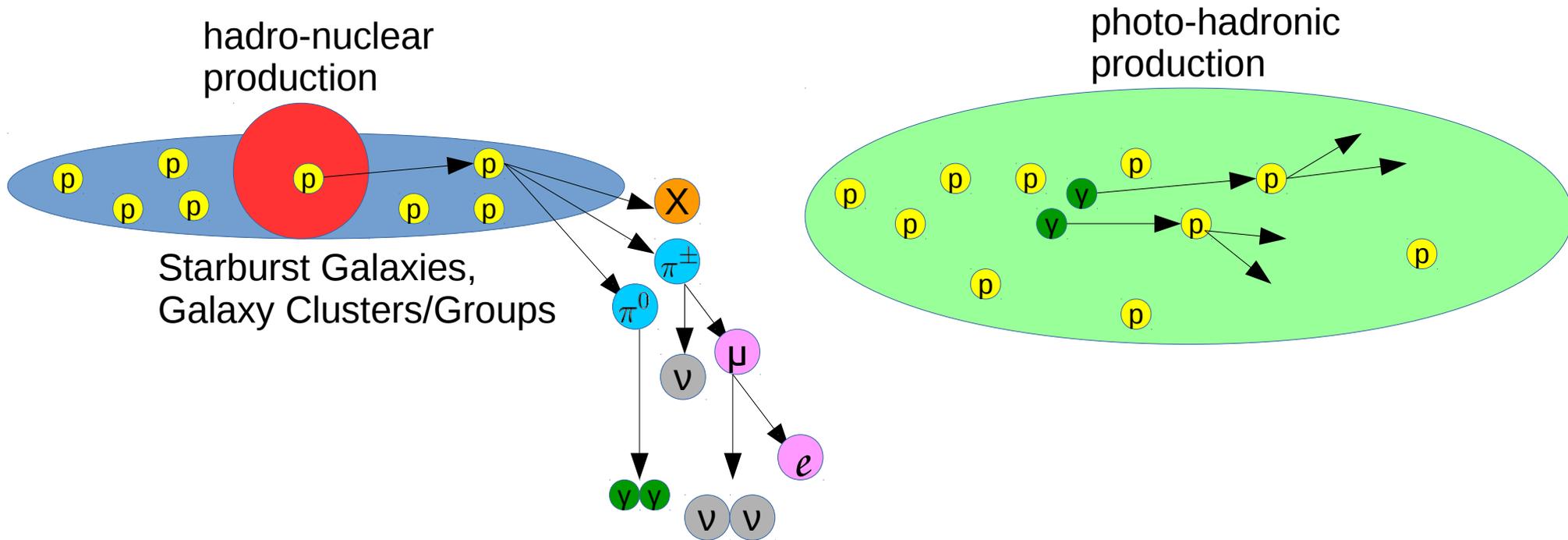
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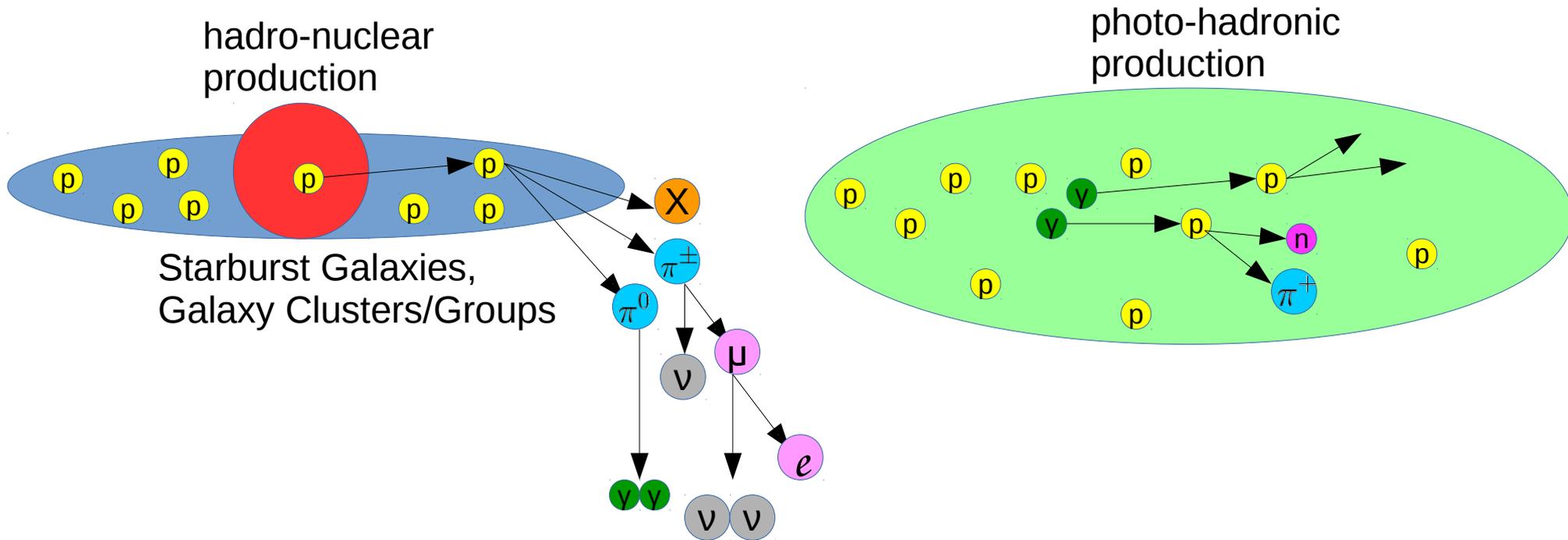
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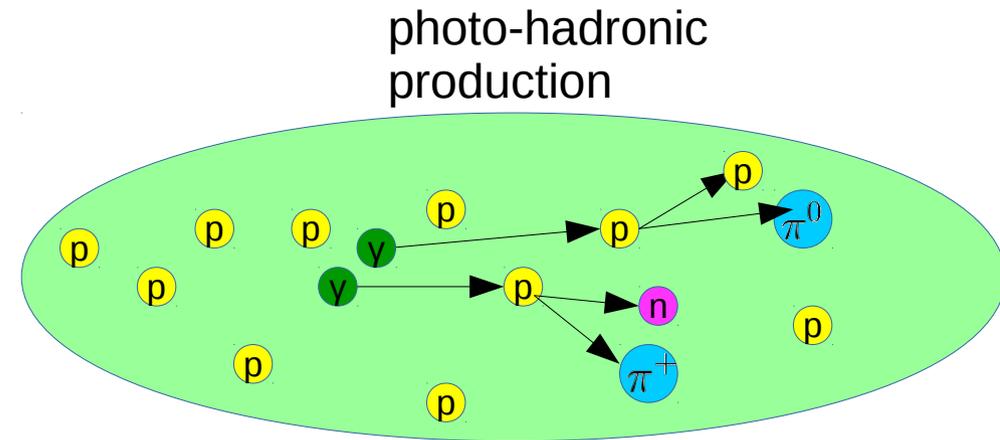
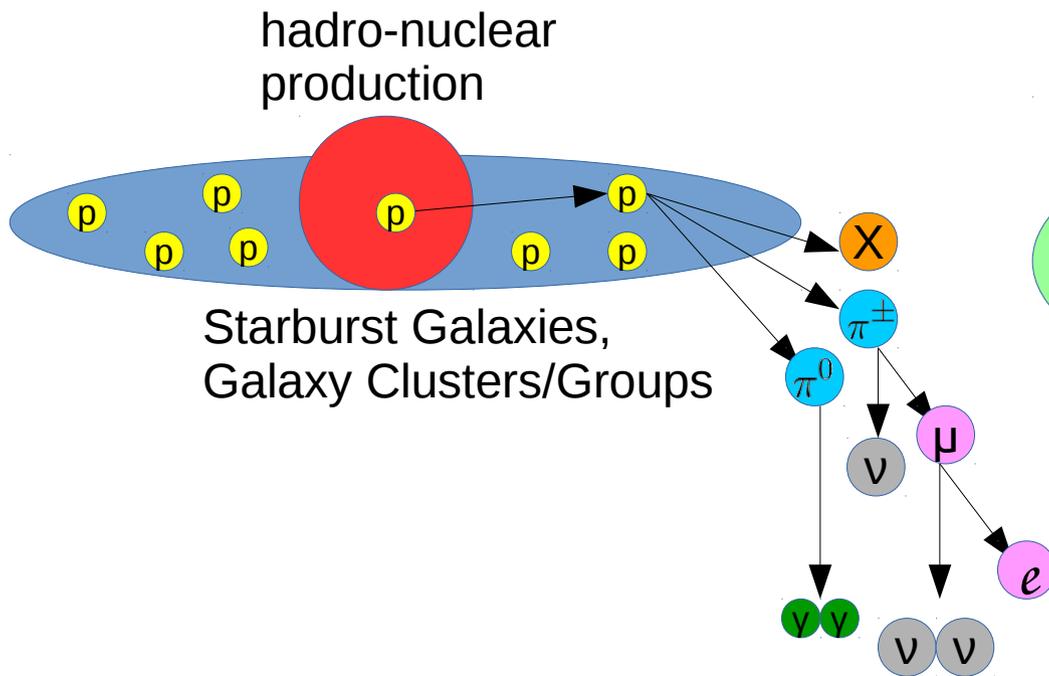
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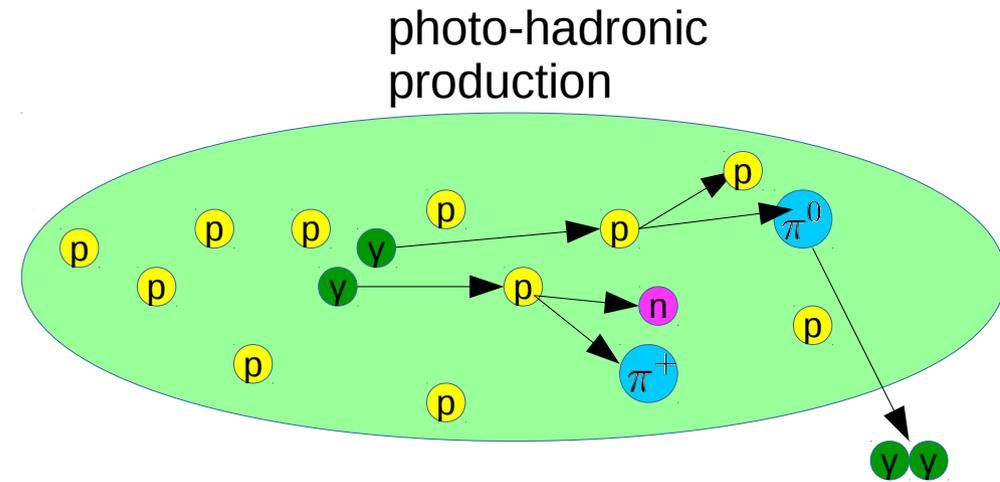
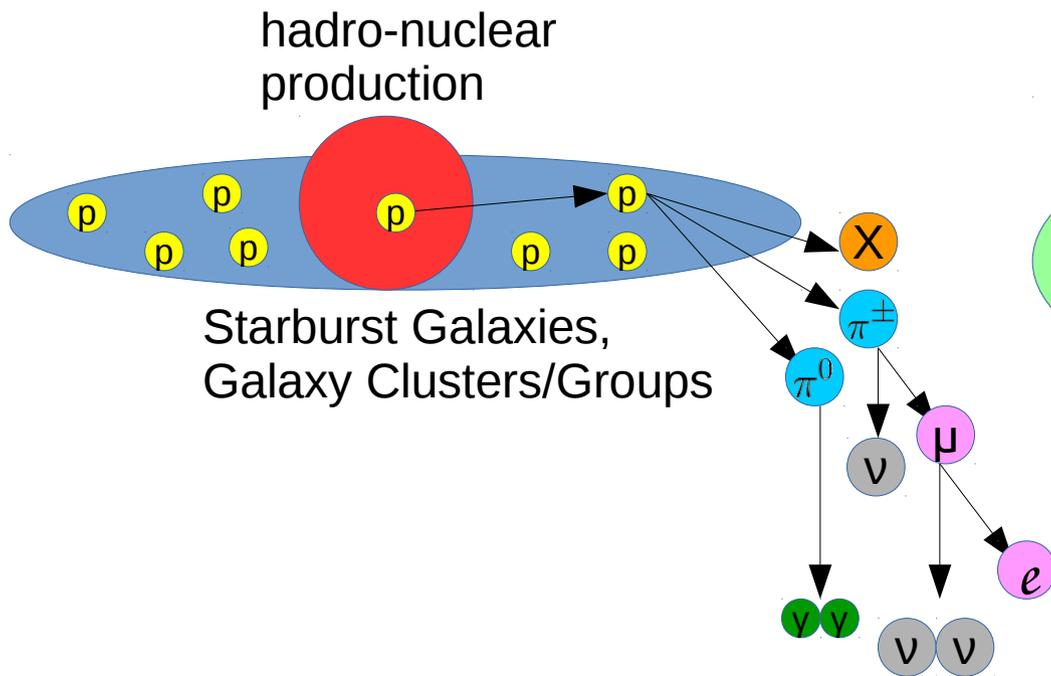
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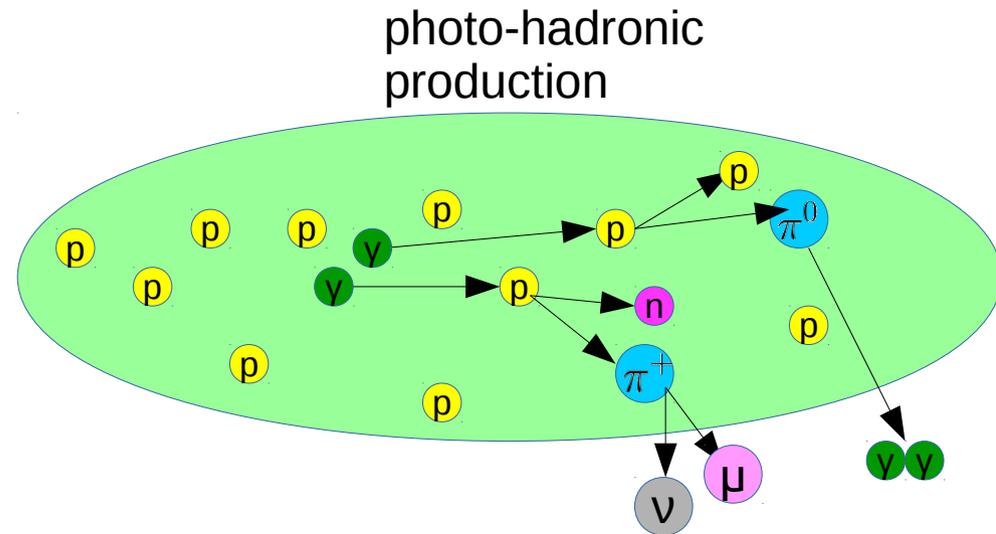
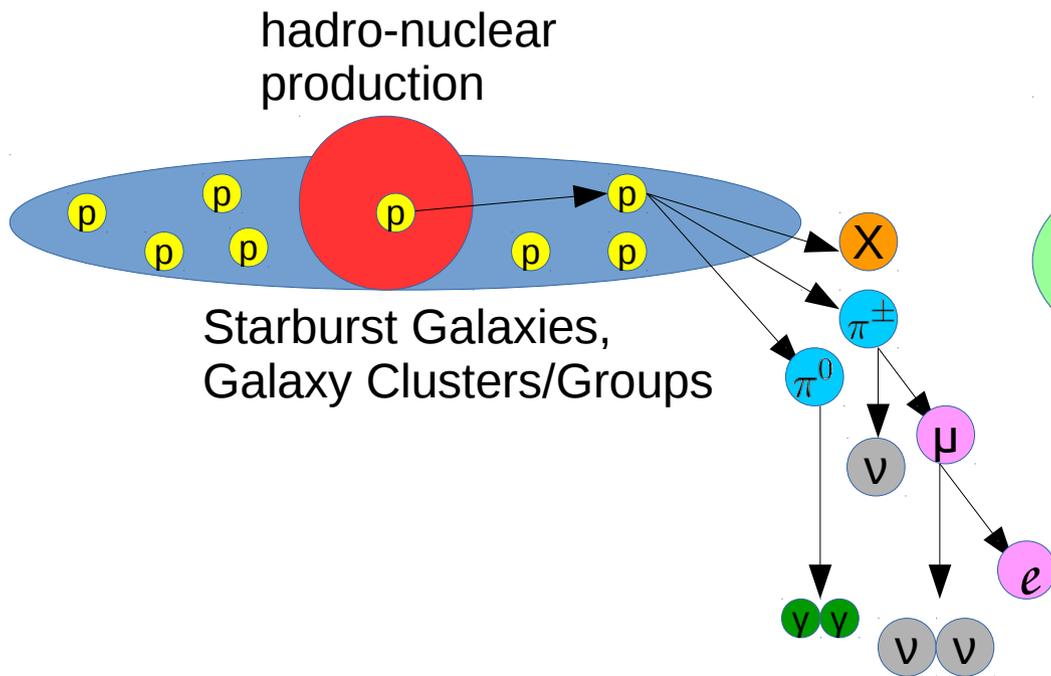
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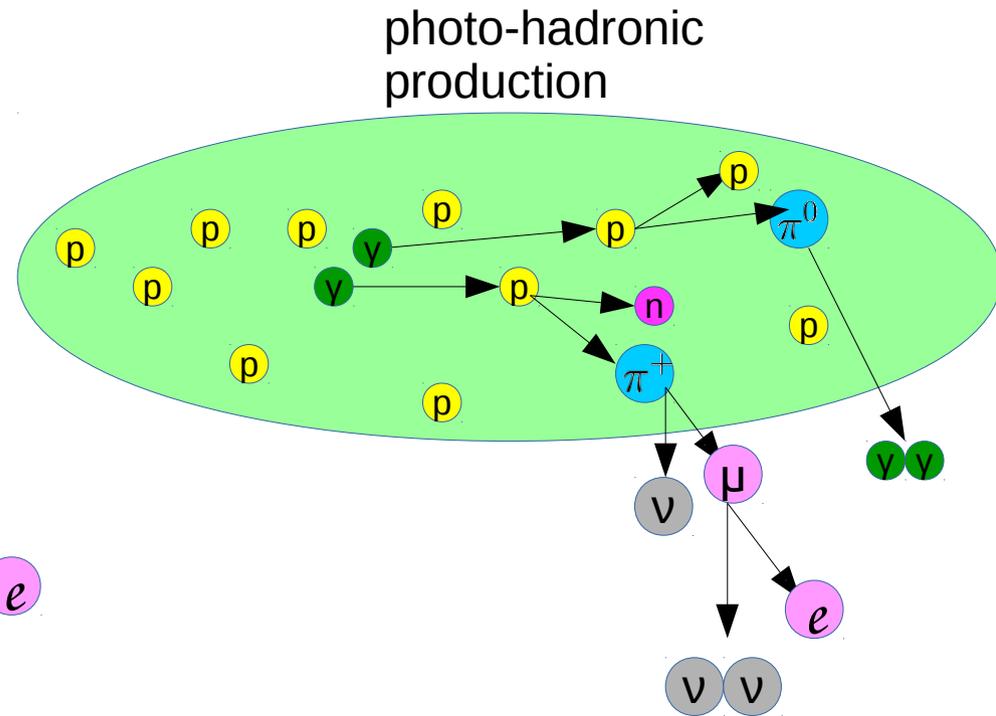
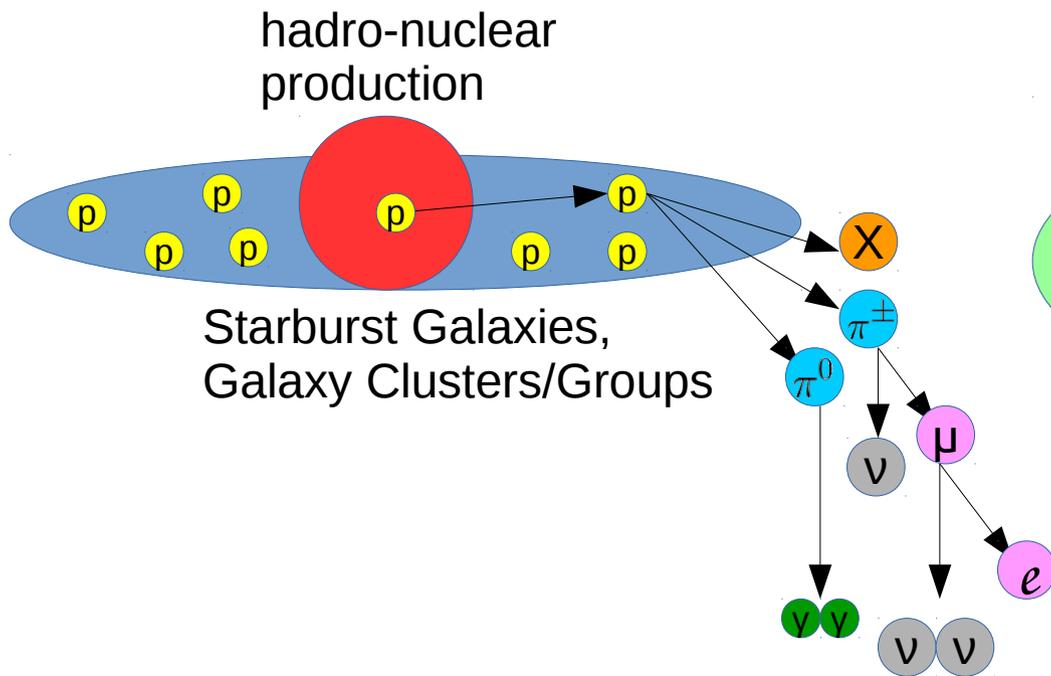
Astrophysical Neutrino Sources



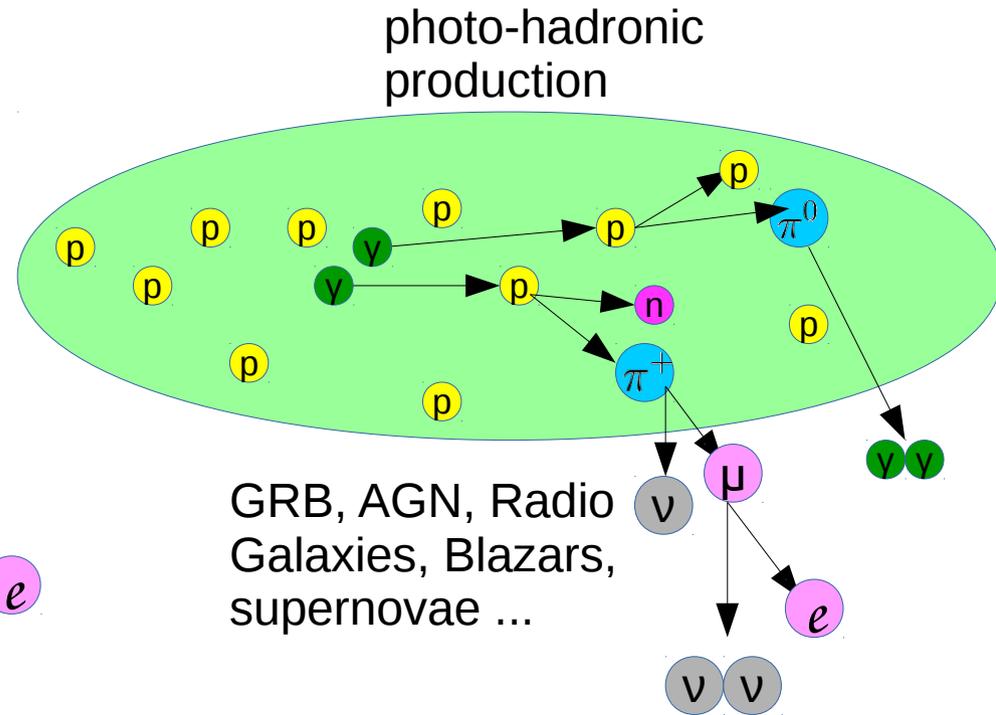
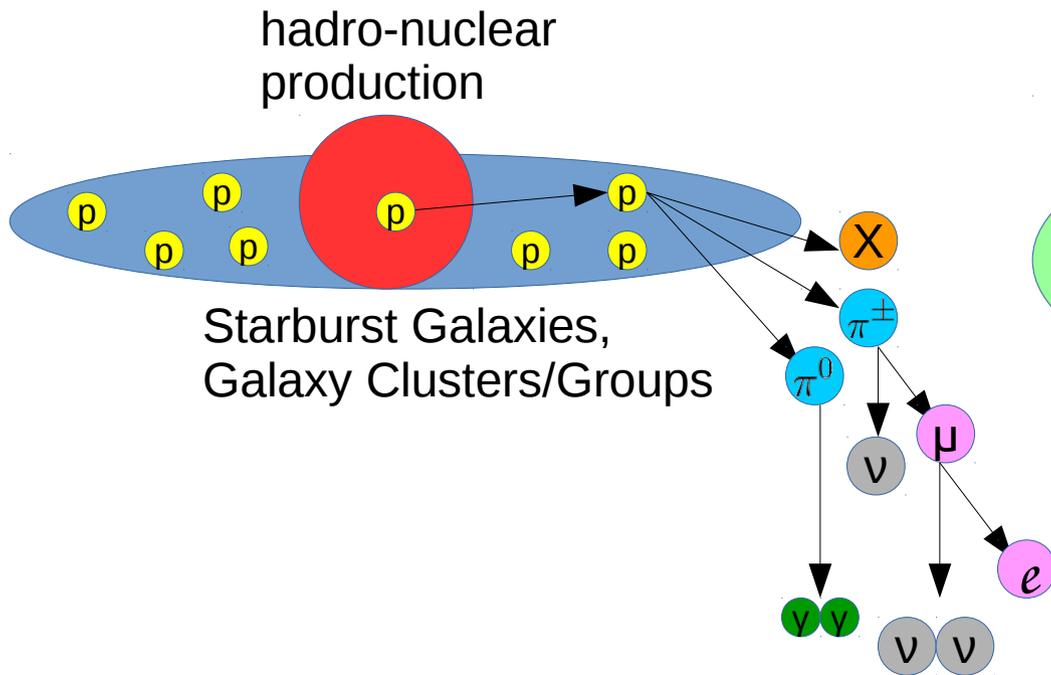
Astrophysical Neutrino Sources



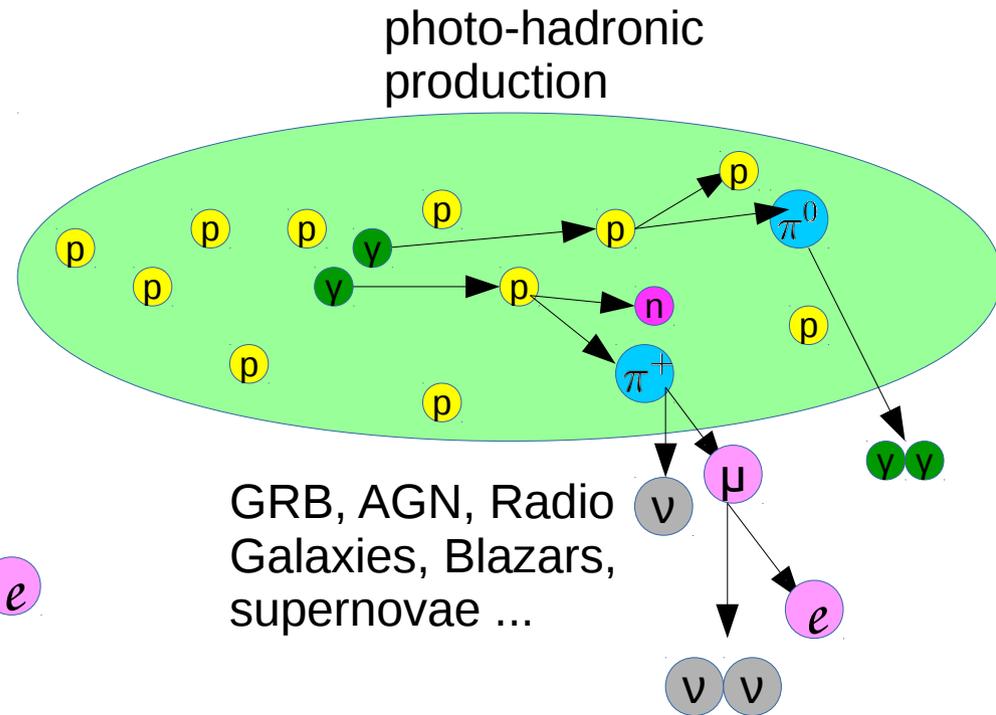
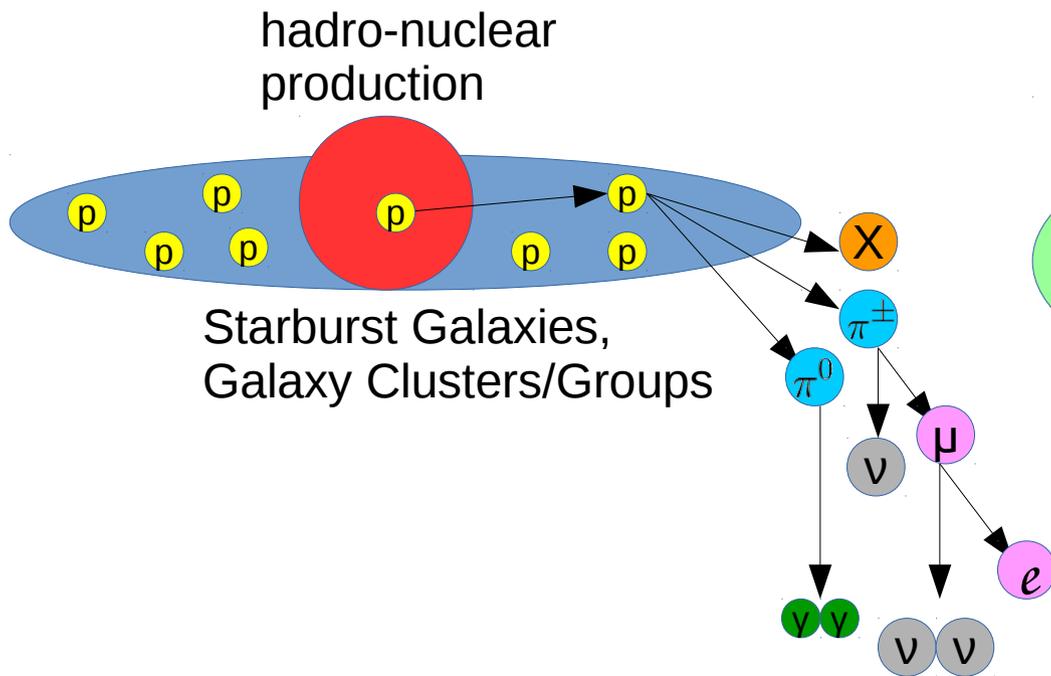
Astrophysical Neutrino Sources



Astrophysical Neutrino Sources



Astrophysical Neutrino Sources

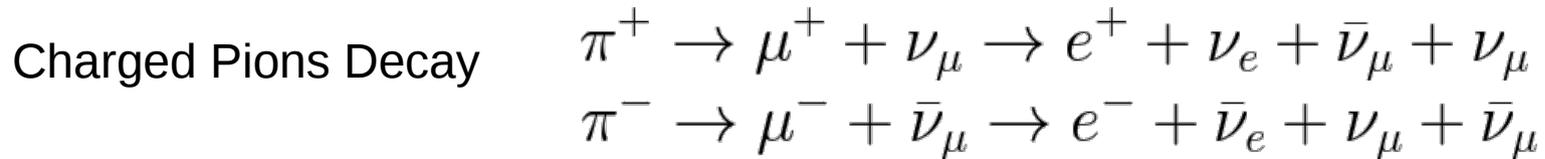
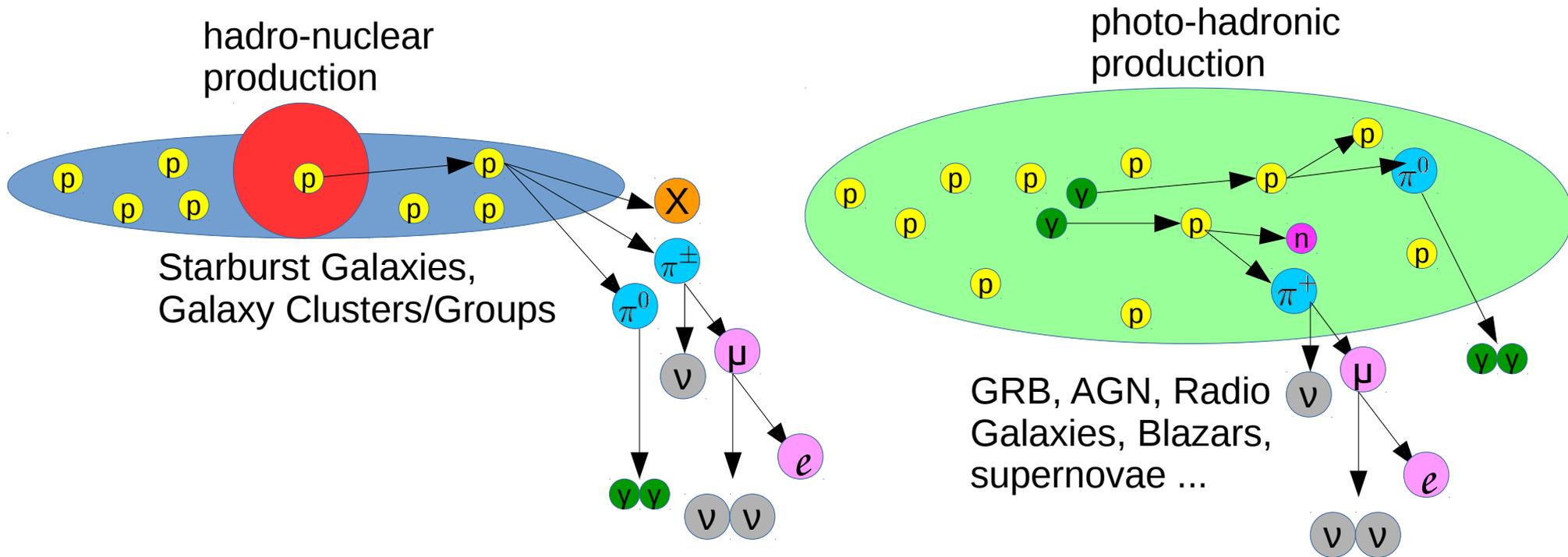


Charged Pions Decay

$$\pi^+ \rightarrow \mu^+ + \nu_\mu \rightarrow e^+ + \nu_e + \bar{\nu}_\mu + \nu_\mu$$

$$\pi^- \rightarrow \mu^- + \bar{\nu}_\mu \rightarrow e^- + \bar{\nu}_e + \nu_\mu + \bar{\nu}_\mu$$

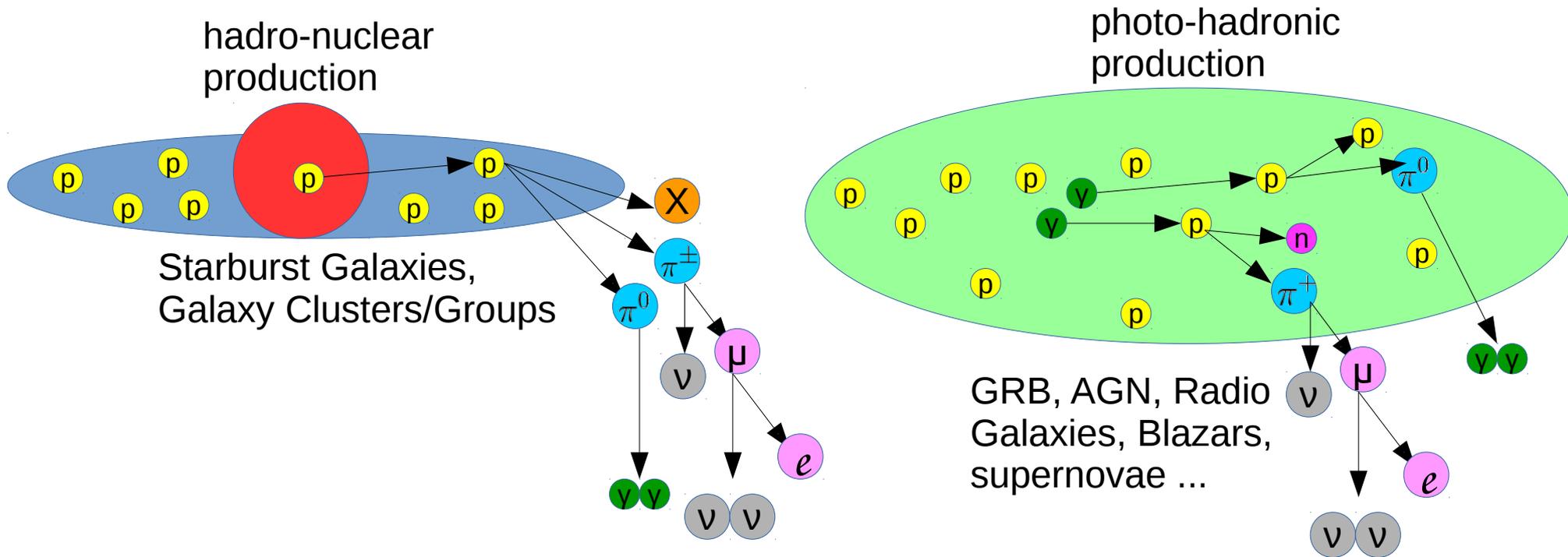
Astrophysical Neutrino Sources



Neutrinos typically have 1-5% of proton energy

Maximally: $E_{\text{GZK}} \sim 5 \times 10^4 \text{ PeV}$

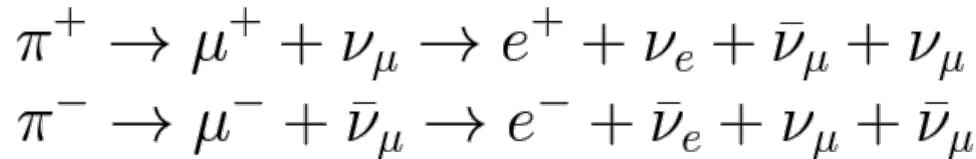
Astrophysical Neutrino Sources



Starburst Galaxies,
Galaxy Clusters/Groups

GRB, AGN, Radio
Galaxies, Blazars,
supernovae ...

Charged Pions Decay

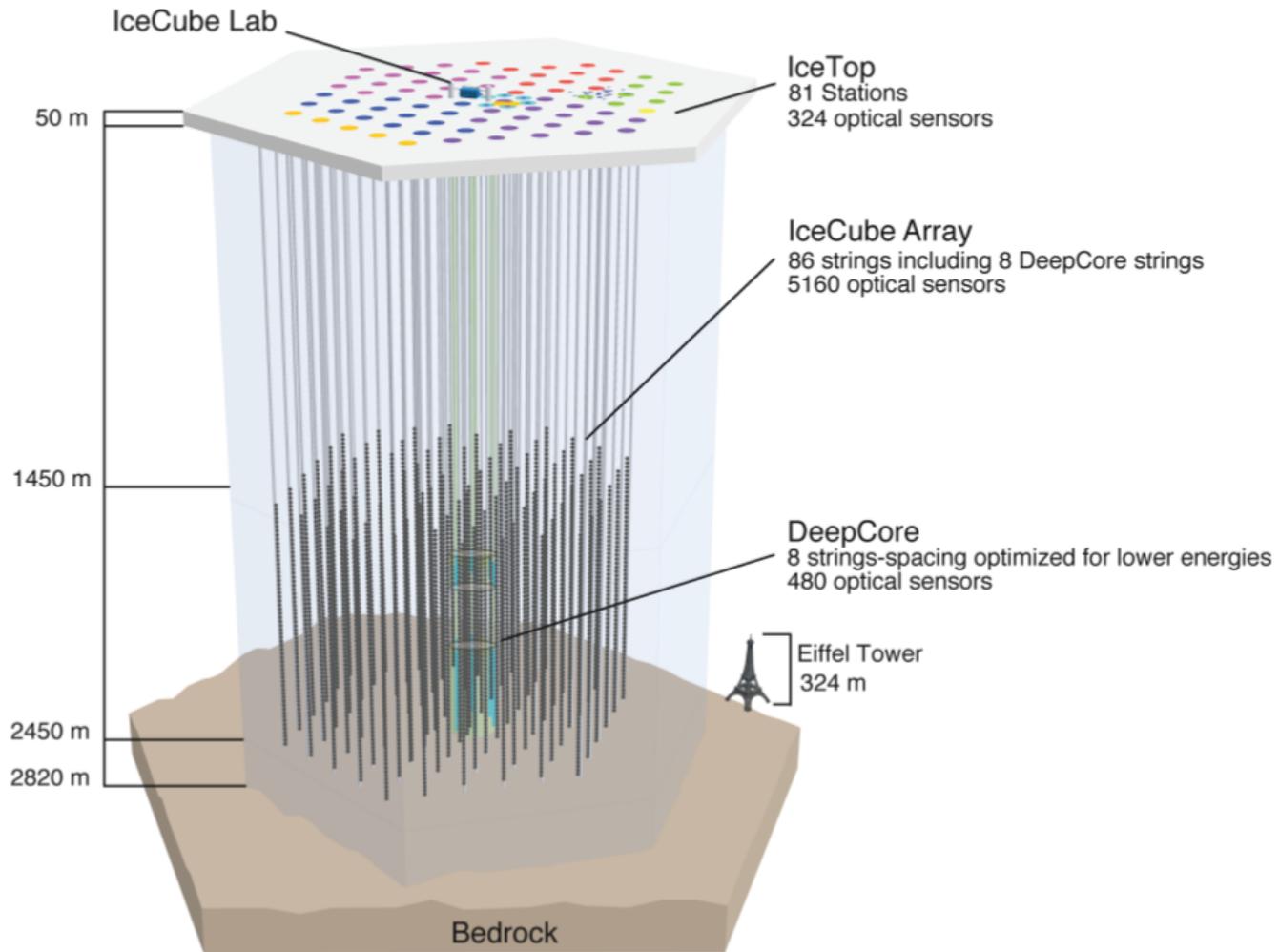


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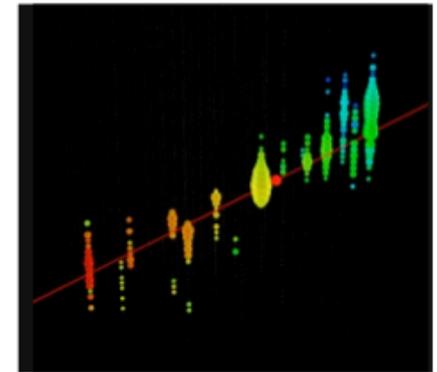
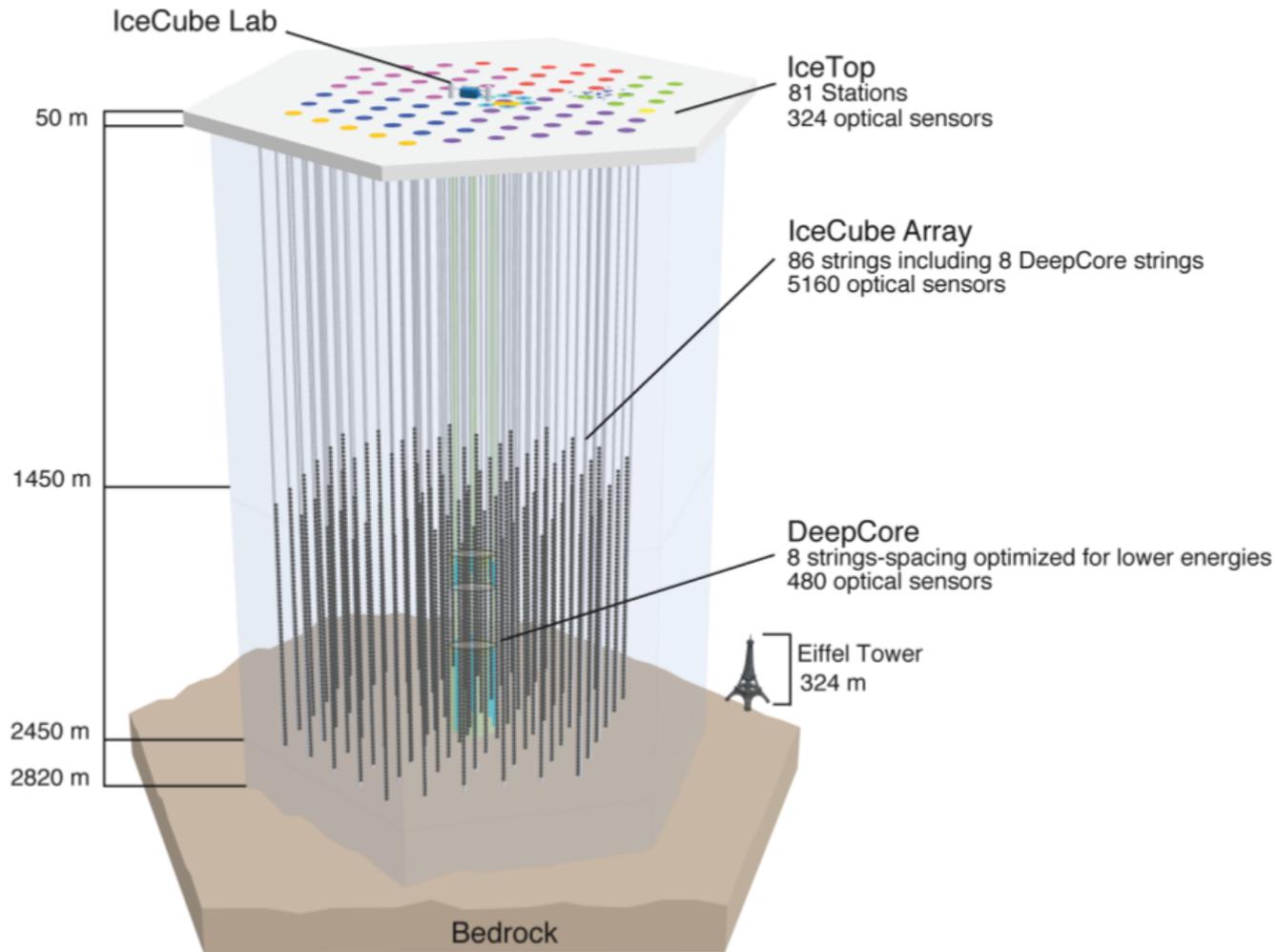
Maximally: $E_{GZK} \sim 5 \times 10^4 \text{ PeV}$

How do we detect them?

IceCube Detector

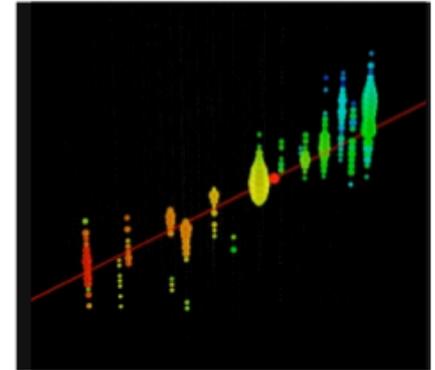
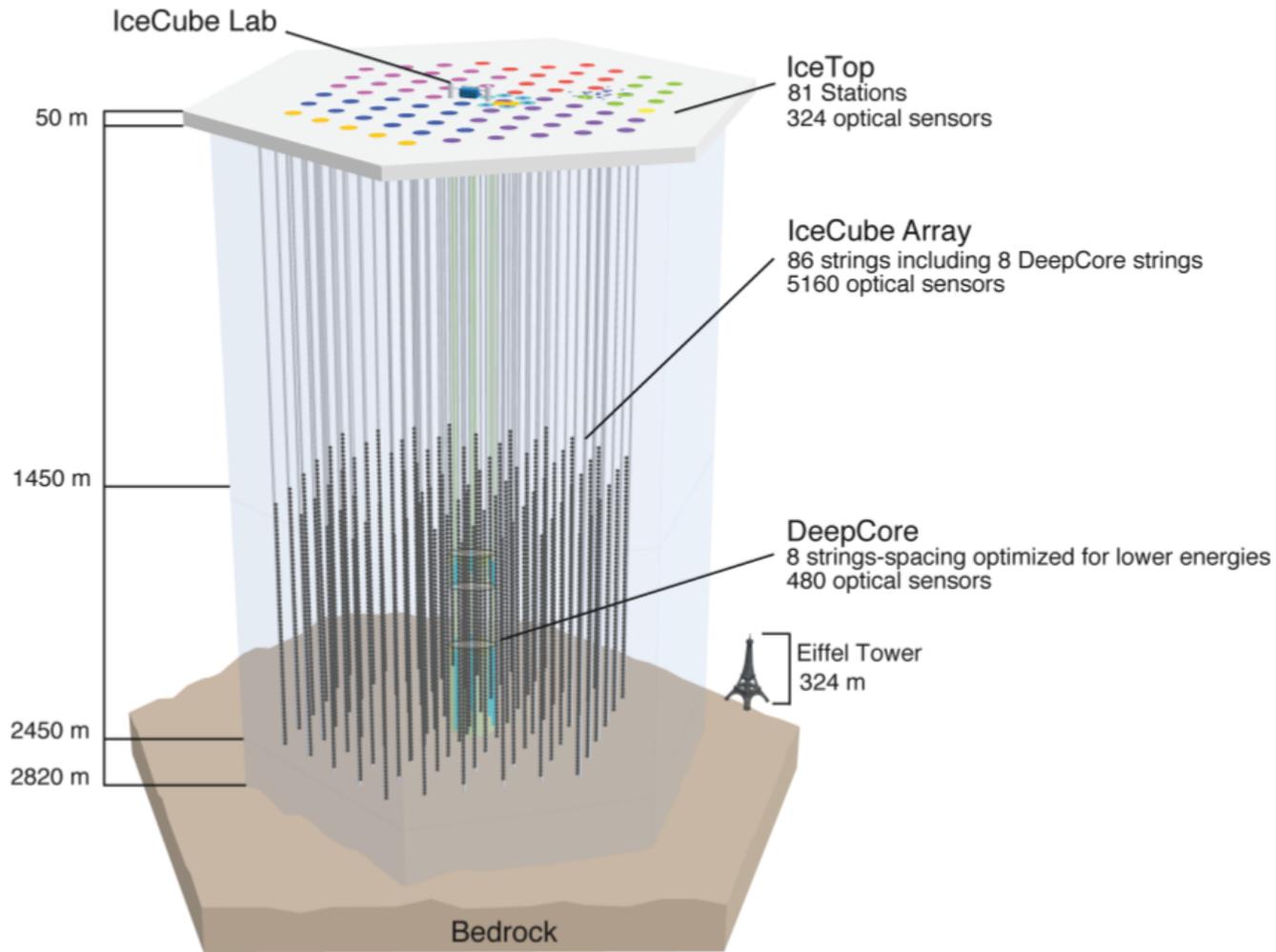


IceCube Detector

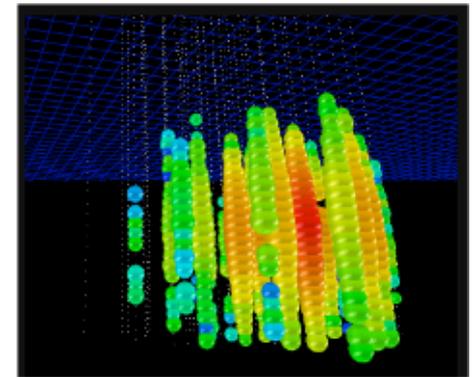


track

IceCube Detector

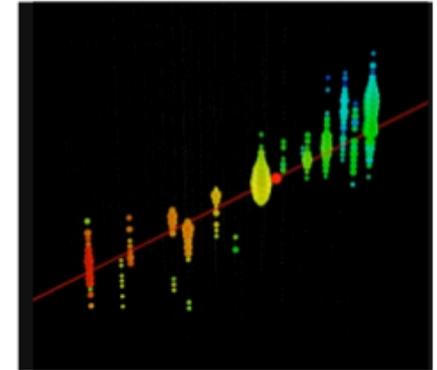
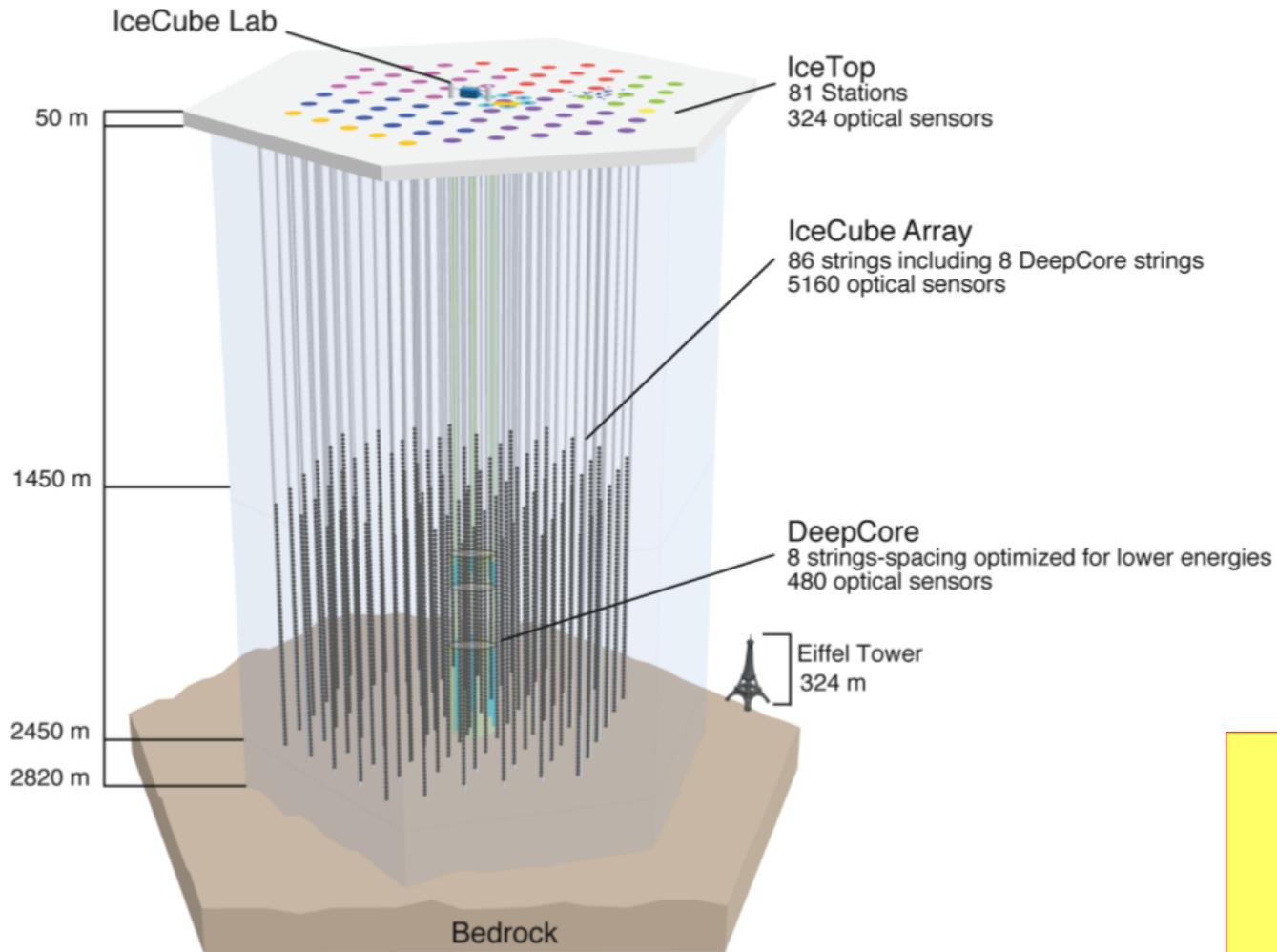


track

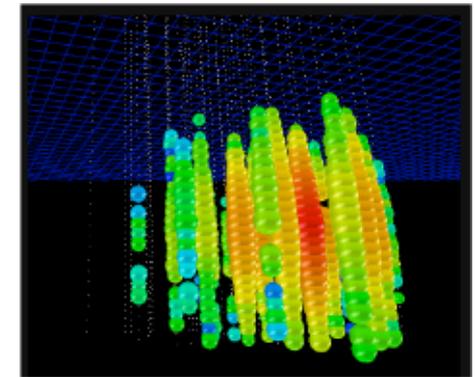


cascade

IceCube Detector



track



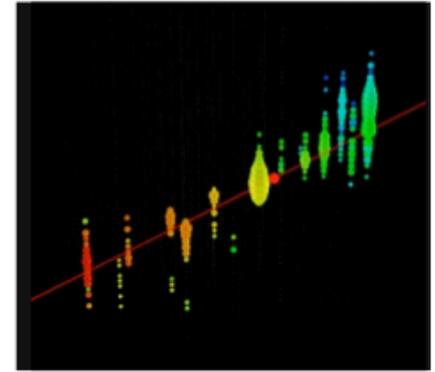
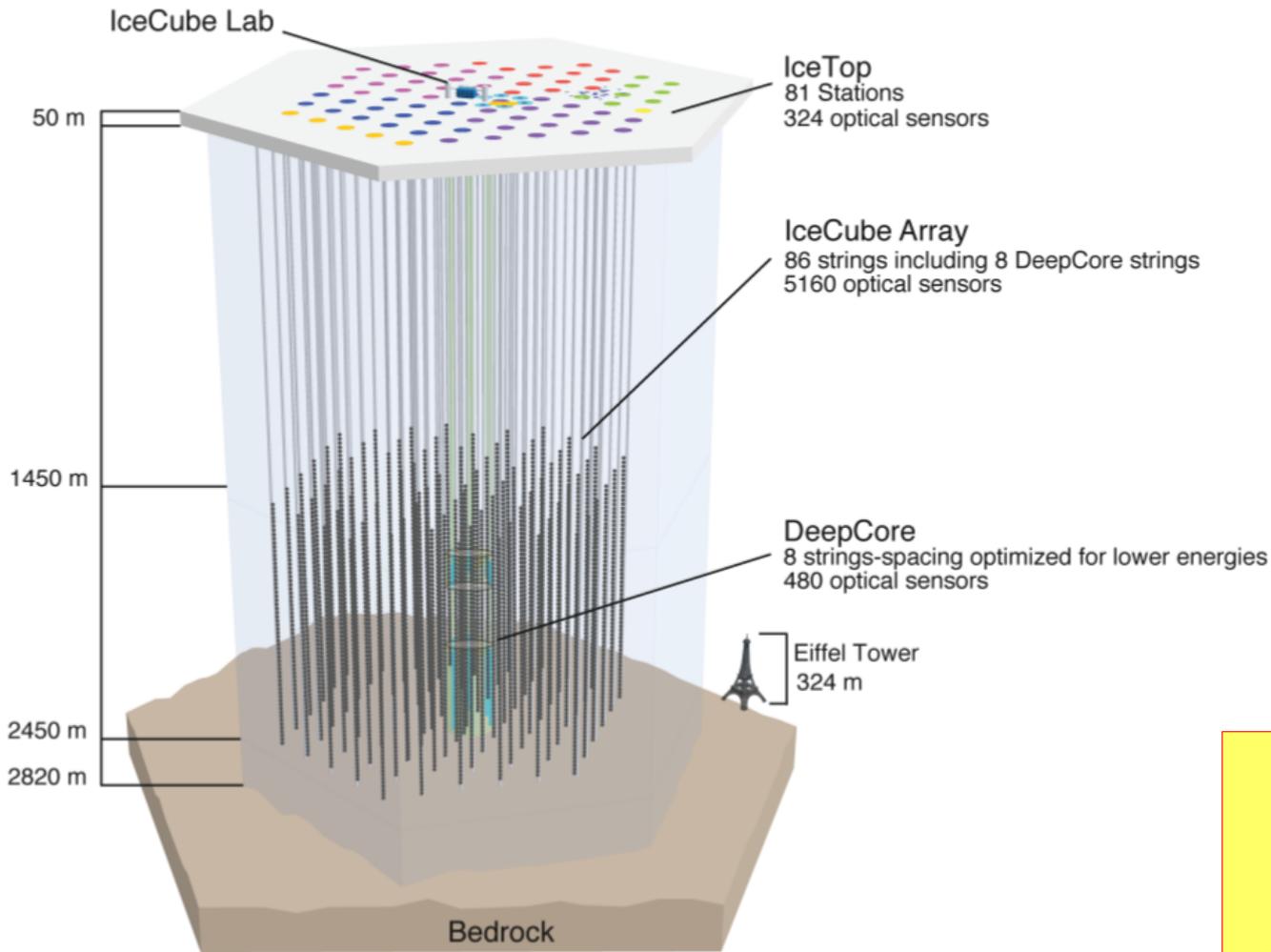
cascade

Mechanism:

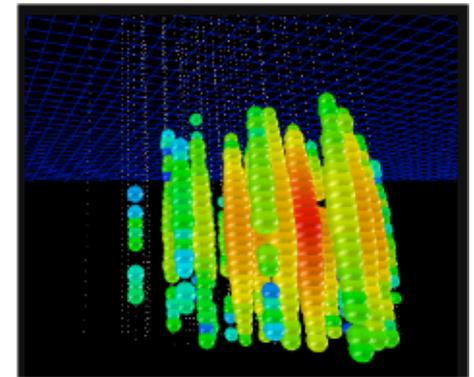
$$\nu_l + N \rightarrow \begin{cases} l + X & (CC) \\ \nu_l + X & (NC) \end{cases}$$

Cherenkov radiation from interaction products: leptons and hadrons

IceCube Detector



track



cascade

Mechanism:

$$\nu_l + N \rightarrow \begin{cases} l + X & (CC) \\ \nu_l + X & (NC) \end{cases}$$

Cherenkov radiation from interaction products: leptons and hadrons

ν_e interactions dominates in special case



IC signal simulation

IC signal simulation

$$N_i^{\text{HESE}} = \int d\Omega \int_{E_{i,\min}}^{E_{i,\max}} dE \sum_{\ell=e,\mu,\tau} \Phi_{\nu_\ell}(E) \cdot T \cdot A_{\nu_\ell}(E, \Omega)$$

IC signal simulation

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$A_{\nu_\ell}(E, \Omega)$: HESE effective area, sum of cross sections for all the particles in the detector, an effective total cross section

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T : Exposure time is 2635 days

IC signal simulation

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$$\Phi(E_\nu) = \Phi_0(E_\nu/100 \text{ TeV})^{-\gamma}$$

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$$\begin{aligned}\Phi(E_\nu) &= \Phi_0 (E_\nu / 100 \text{ TeV})^{-\gamma} \\ \Phi_0 &= 6.45_{-0.46}^{+1.46} \text{ GeV}^{-1} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1} \\ \gamma &= 2.89_{-0.19}^{+0.20}\end{aligned}$$

IC signal simulation

$$N_i^{\text{HESE}} = \int d\Omega \int_{E_{i,\min}}^{E_{i,\max}} dE \sum_{\ell=e,\mu,\tau} \Phi_{\nu_\ell}(E) \cdot T \cdot A_{\nu_\ell}(E, \Omega)$$

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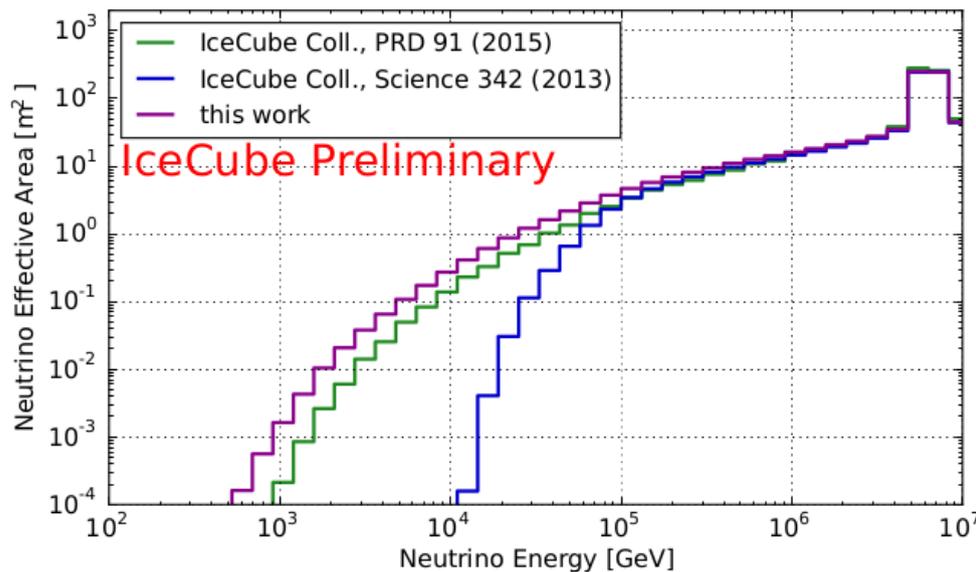
T : Exposure time is 2635 days

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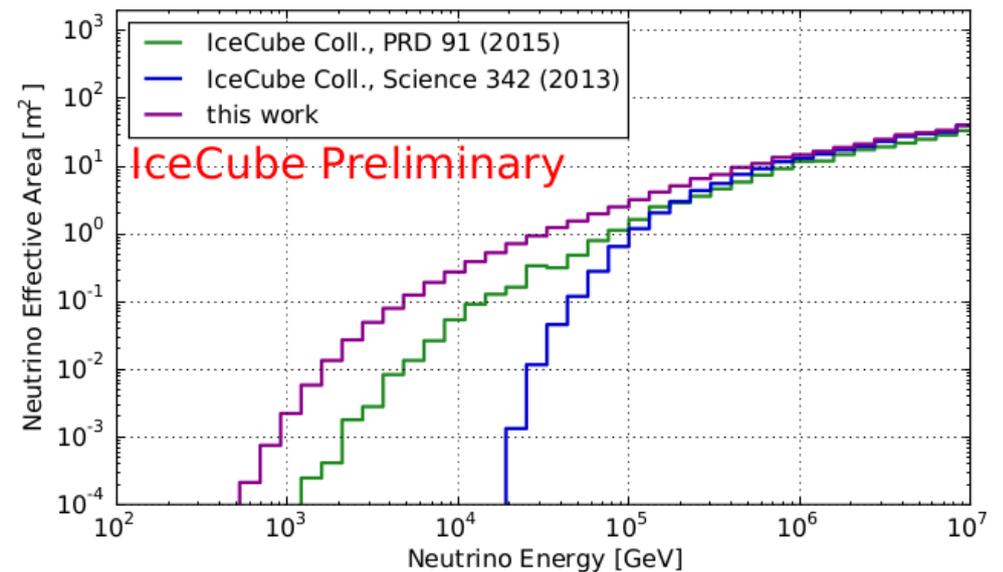
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HESE e neutrino effective area



HESE muon neutrino effective area



IC signal simulation

$$N_i^{\text{HESE}} = \int d\Omega \int_{E_{i,\min}}^{E_{i,\max}} dE \sum_{\ell=e,\mu,\tau} \Phi_{\nu_\ell}(E) \cdot T \cdot A_{\nu_\ell}(E, \Omega)$$

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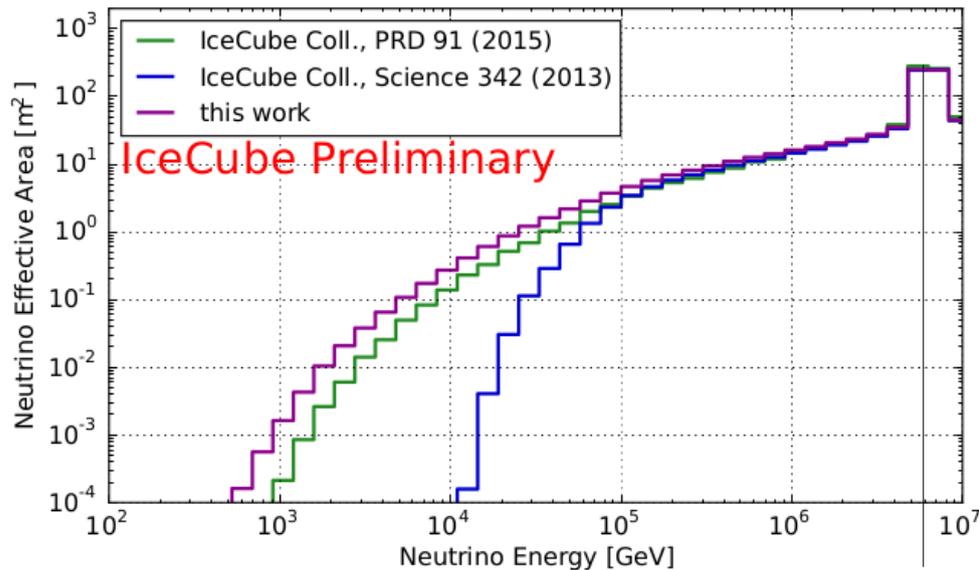
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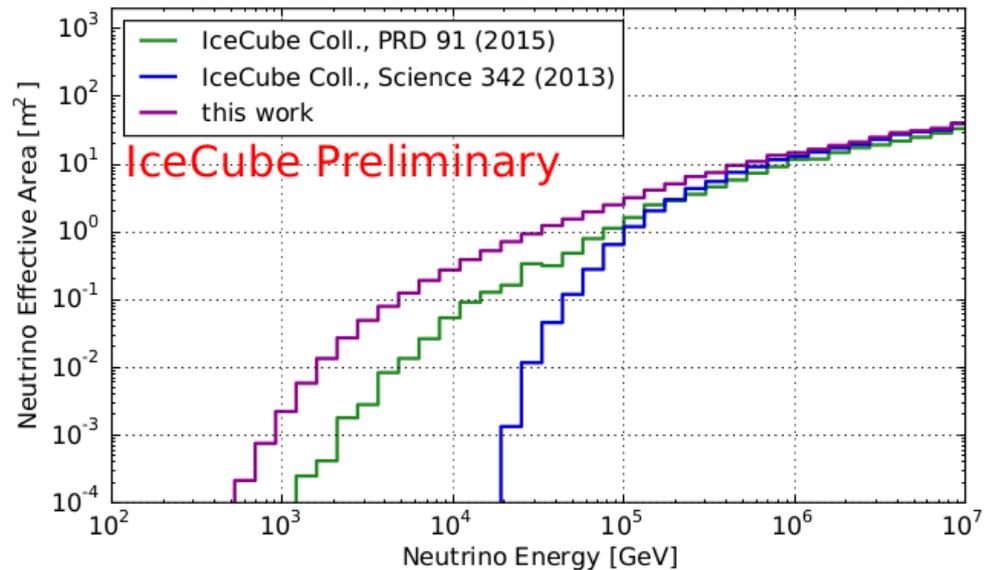
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HESE e neutrino effective area



6.3 PeV

HESE muon neutrino effective area



IC signal simulation

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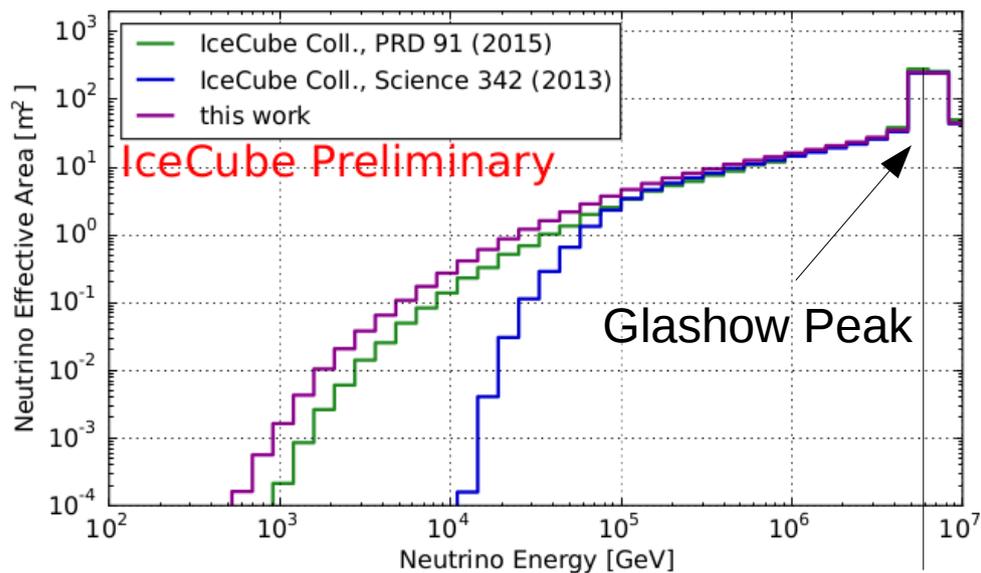
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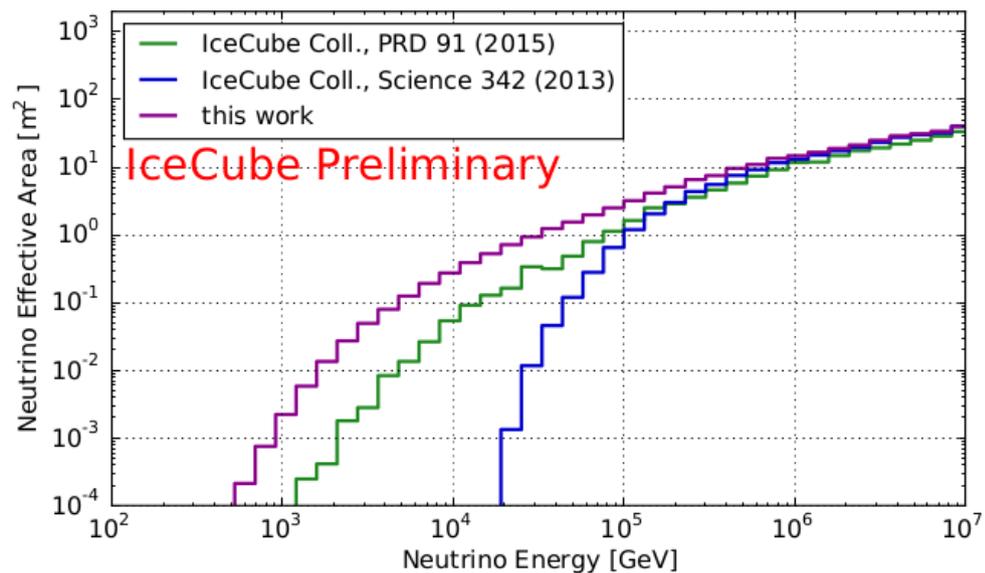
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HESE e neutrino effective area

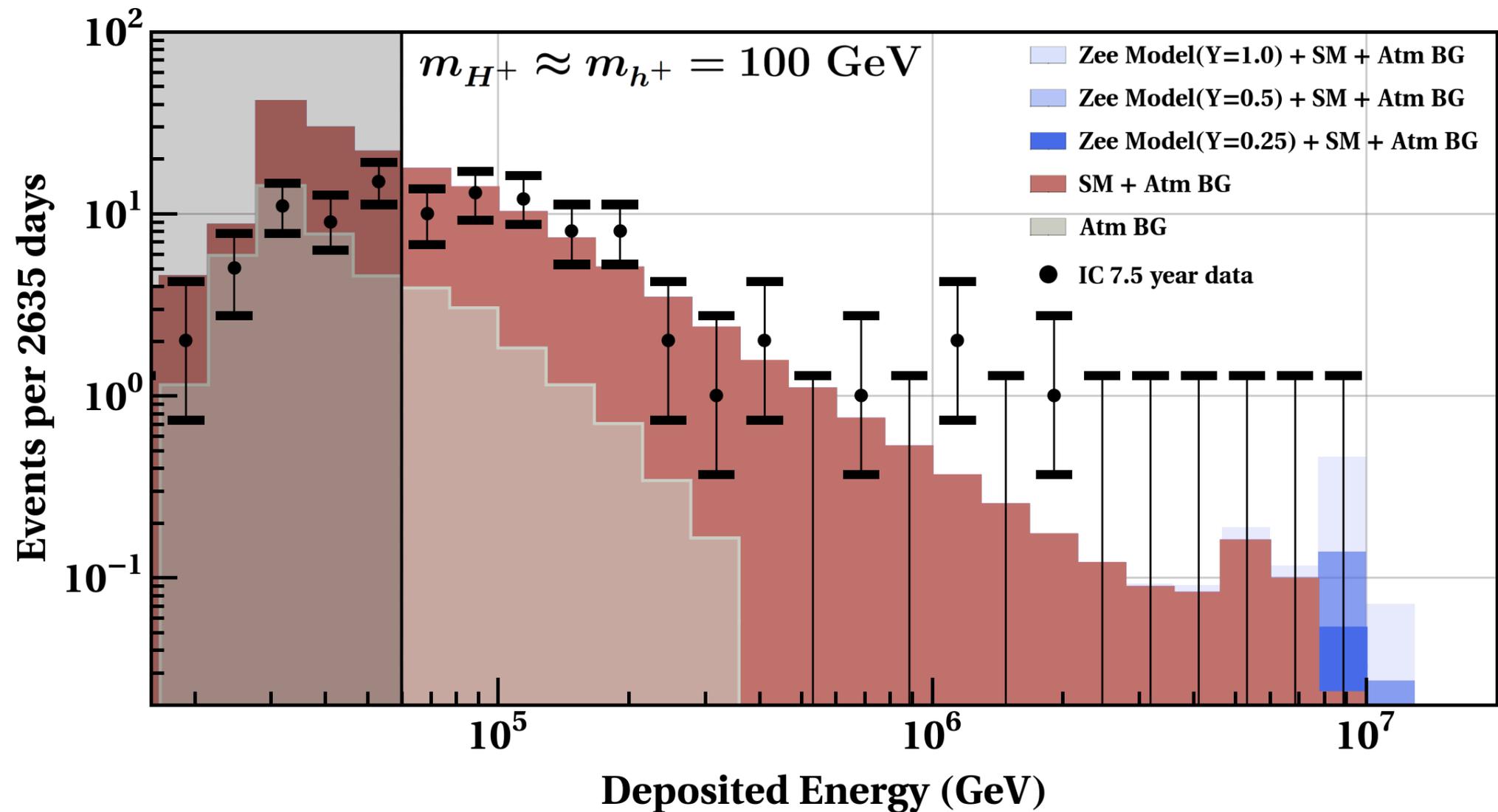


6.3 PeV

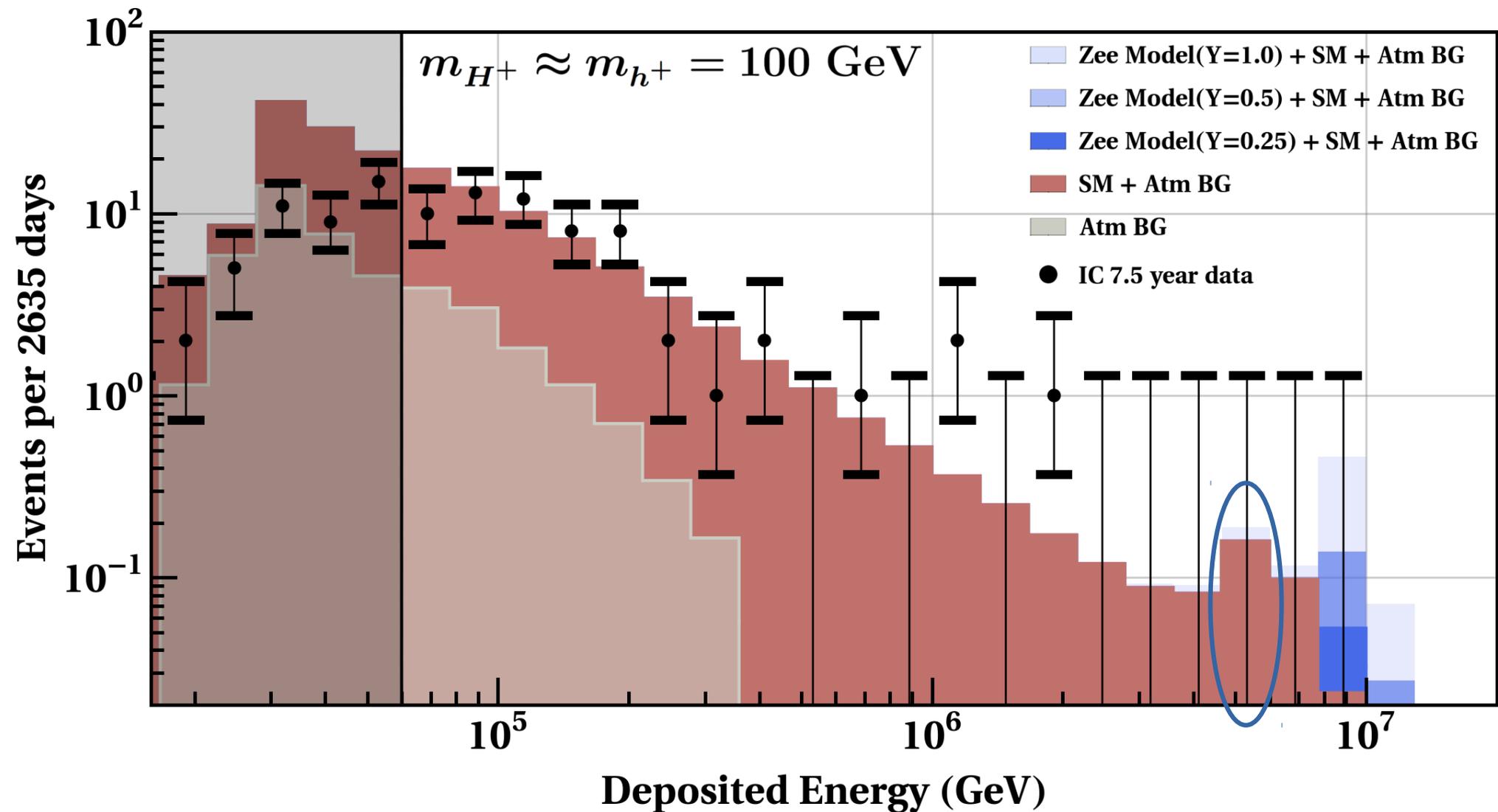
HESE muon neutrino effective area



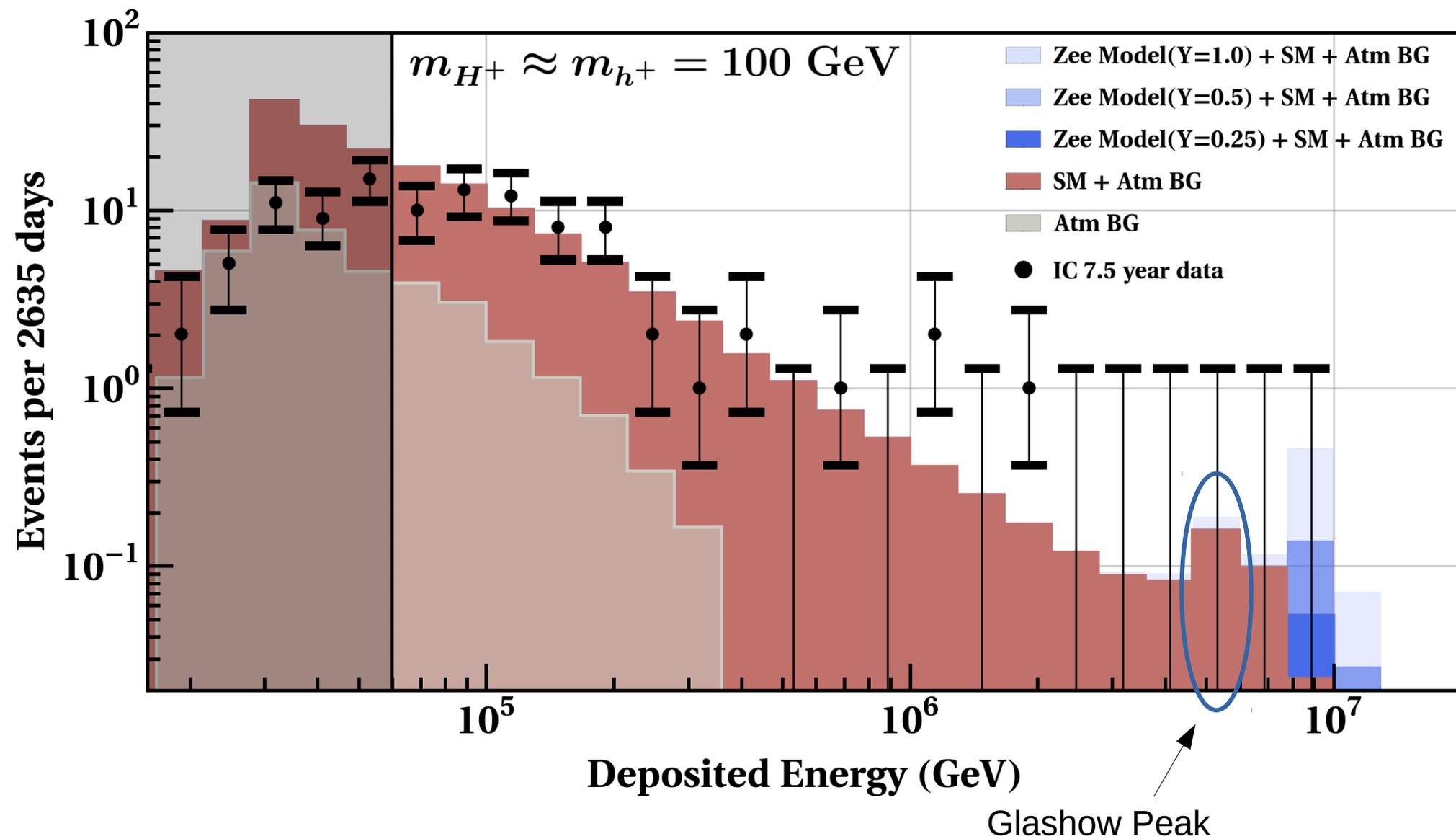
Spectrum plot



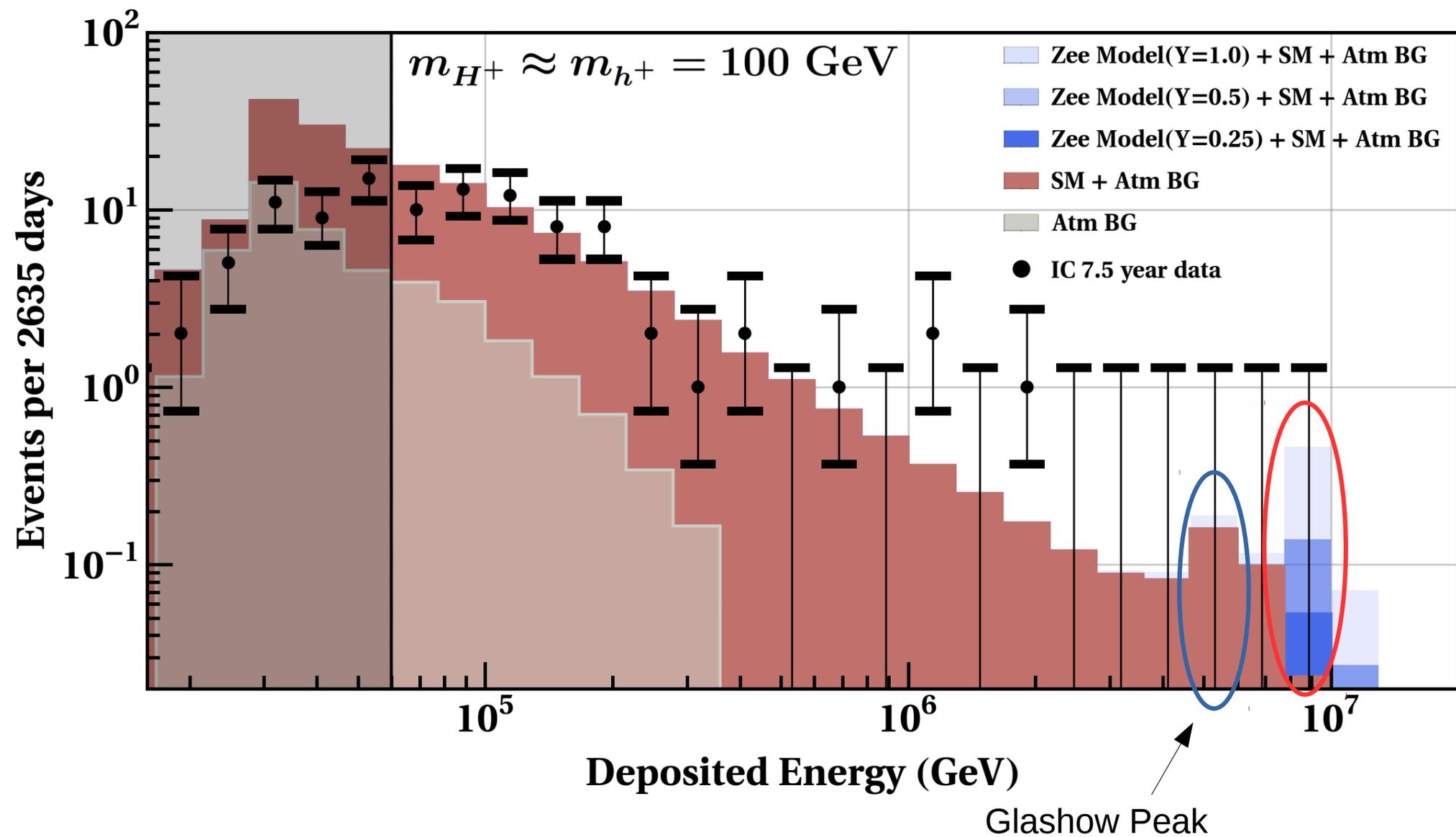
Spectrum plot



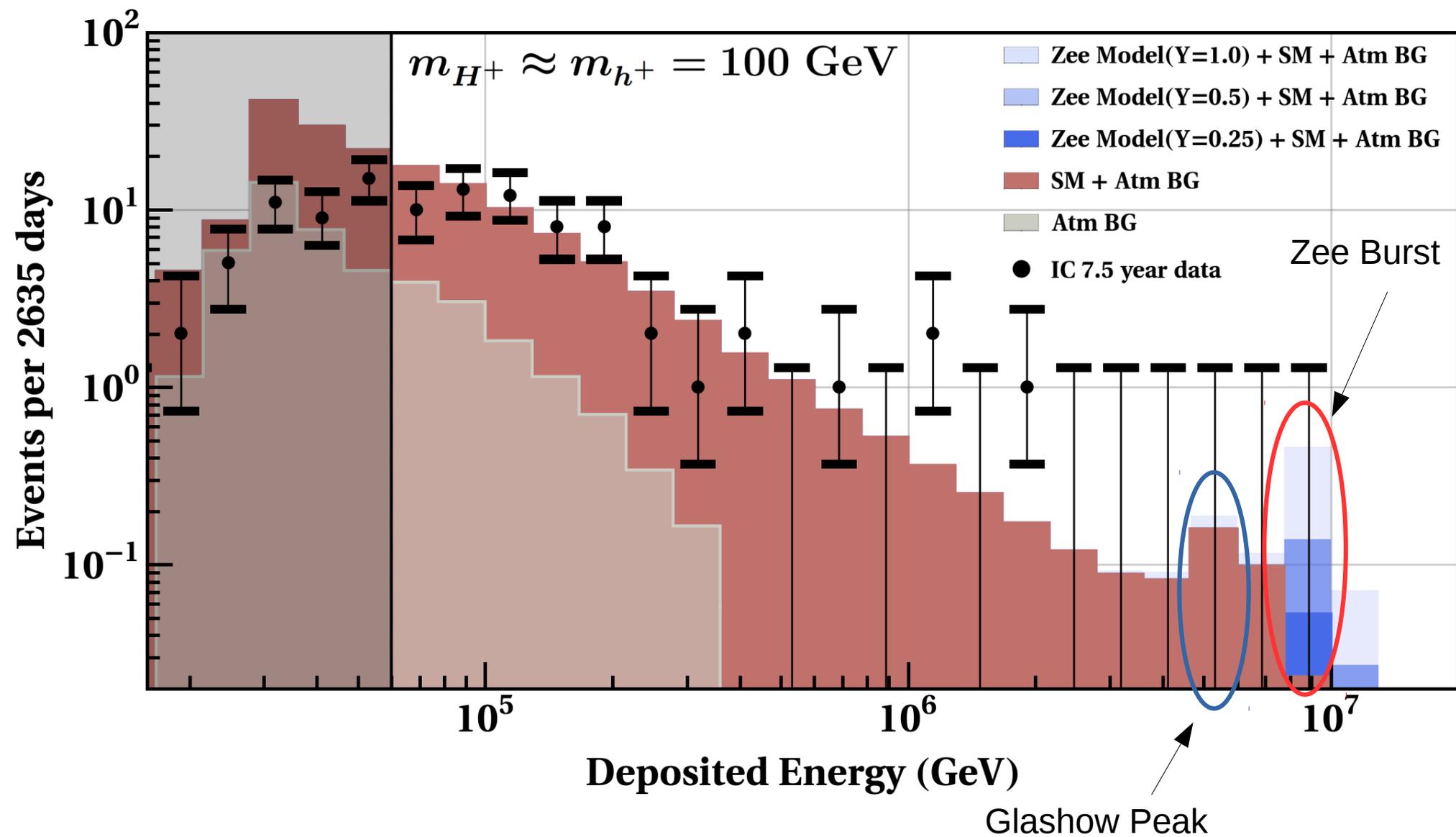
Spectrum plot



Spectrum plot



Spectrum plot





NSI from Zee Model

NSI from Zee Model

$$-\mathcal{L}_Y \supset f_{\alpha\beta} L_\alpha^i L_\beta^j \epsilon_{ij} \eta^+ + \tilde{Y}_{\alpha\beta} \tilde{H}_1^i L_\alpha^j \ell_\beta^c \epsilon_{ij} + Y_{\alpha\beta} \tilde{H}_2^i L_\alpha^j \ell_\beta^c \epsilon_{ij} + \text{H.c.}$$

NSI from Zee Model

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Effectively, we have:

NSI from Zee Model

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Effectively, we have:

$$\mathcal{L}_{\text{eff}} = -\frac{1}{2} Y_{\alpha\rho} Y_{\beta\sigma}^* \left(\frac{\sin^2 \varphi}{m_{h^+}^2} + \frac{\cos^2 \varphi}{m_{H^+}^2} \right) (\bar{\nu}_\alpha \gamma^\mu P_L \nu_\beta) (\bar{\ell}_\sigma \gamma_\mu P_R \ell_\rho)$$

NSI from Zee Model

$$-\mathcal{L}_Y \supset f_{\alpha\beta} L_\alpha^i L_\beta^j \epsilon_{ij} \eta^+ + \tilde{Y}_{\alpha\beta} \tilde{H}_1^i L_\alpha^j \ell_\beta^c \epsilon_{ij} + Y_{\alpha\beta} \tilde{H}_2^i L_\alpha^j \ell_\beta^c \epsilon_{ij} + \text{H.c.}$$

Effectively, we have:

$$\mathcal{L}_{\text{eff}} = -\frac{1}{2} Y_{\alpha\rho} Y_{\beta\sigma}^* \left(\frac{\sin^2 \varphi}{m_{h^+}^2} + \frac{\cos^2 \varphi}{m_{H^+}^2} \right) (\bar{\nu}_\alpha \gamma^\mu P_L \nu_\beta) (\bar{\ell}_\sigma \gamma_\mu P_R \ell_\rho)$$

$$\varepsilon_{\alpha\beta} = \frac{Y_{\alpha e} Y_{\beta e}^*}{4\sqrt{2}G_F} \left(\frac{\sin^2 \varphi}{m_{h^+}^2} + \frac{\cos^2 \varphi}{m_{H^+}^2} \right)$$

NSI from Zee Model

$$-\mathcal{L}_Y \supset f_{\alpha\beta} L_\alpha^i L_\beta^j \epsilon_{ij} \eta^+ + \tilde{Y}_{\alpha\beta} \tilde{H}_1^i L_\alpha^j \ell_\beta^c \epsilon_{ij} + Y_{\alpha\beta} \tilde{H}_2^i L_\alpha^j \ell_\beta^c \epsilon_{ij} + \text{H.c.}$$

Effectively, we have:

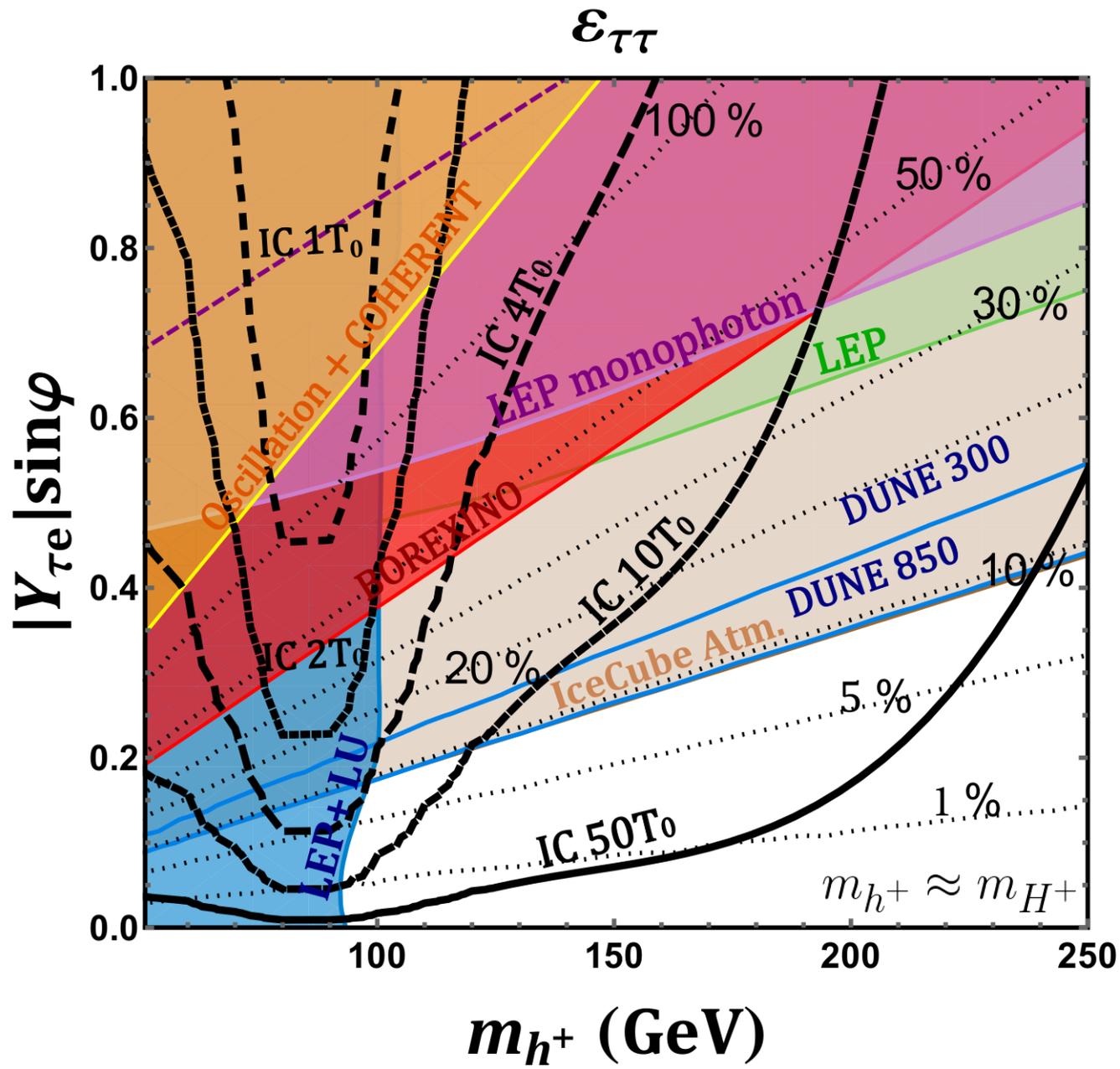
$$\mathcal{L}_{\text{eff}} = -\frac{1}{2} Y_{\alpha\rho} Y_{\beta\sigma}^* \left(\frac{\sin^2 \varphi}{m_{h^+}^2} + \frac{\cos^2 \varphi}{m_{H^+}^2} \right) (\bar{\nu}_\alpha \gamma^\mu P_L \nu_\beta) (\bar{\ell}_\sigma \gamma_\mu P_R \ell_\rho)$$

$$\varepsilon_{\alpha\beta} = \frac{Y_{\alpha e} Y_{\beta e}^*}{4\sqrt{2}G_F} \left(\frac{\sin^2 \varphi}{m_{h^+}^2} + \frac{\cos^2 \varphi}{m_{H^+}^2} \right)$$

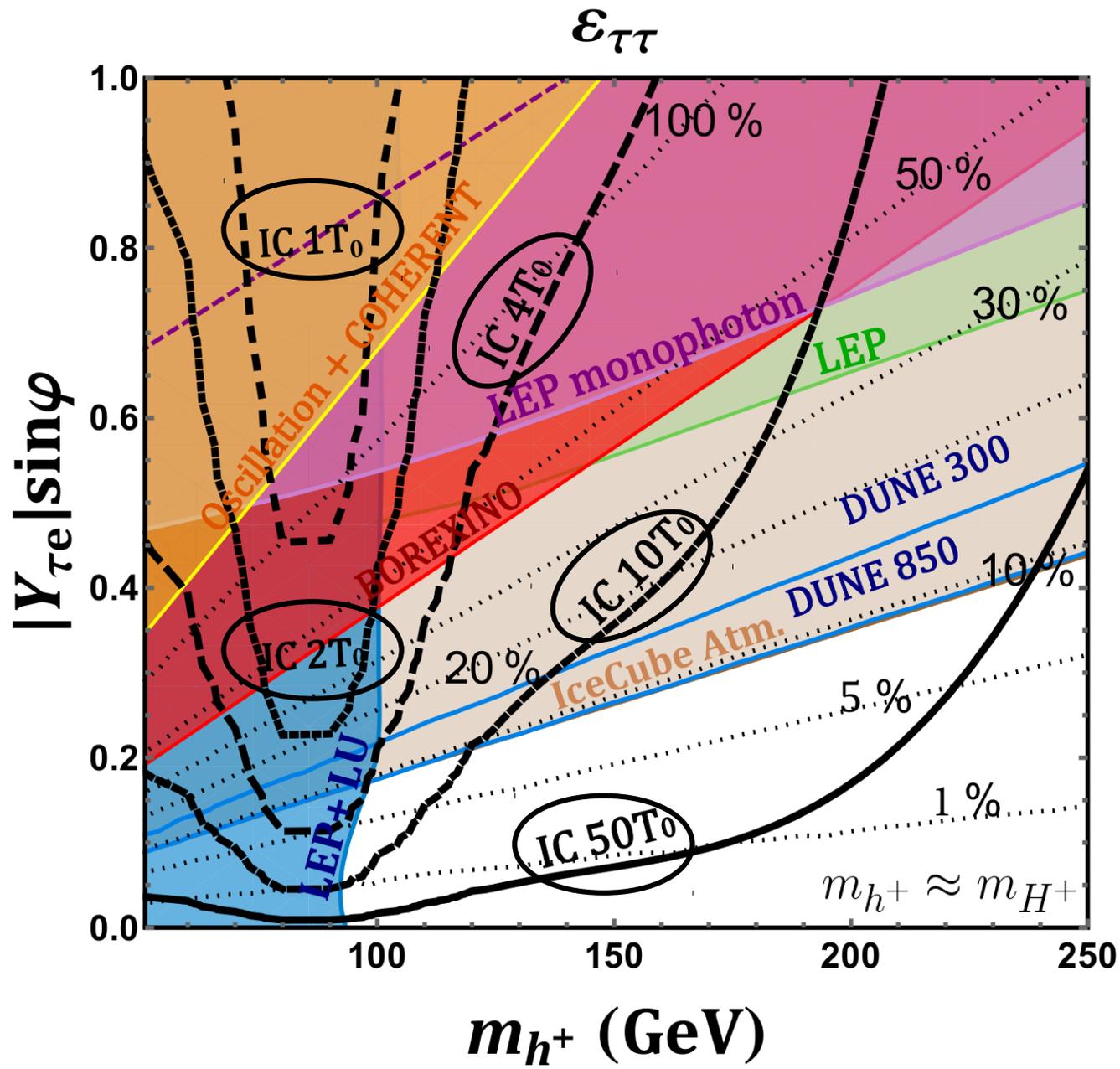
Condition for Maximum contribution to NSI and to Zee burst:

$$m_h^+ \approx m_H^+, \quad \sin\varphi = \cos\varphi$$

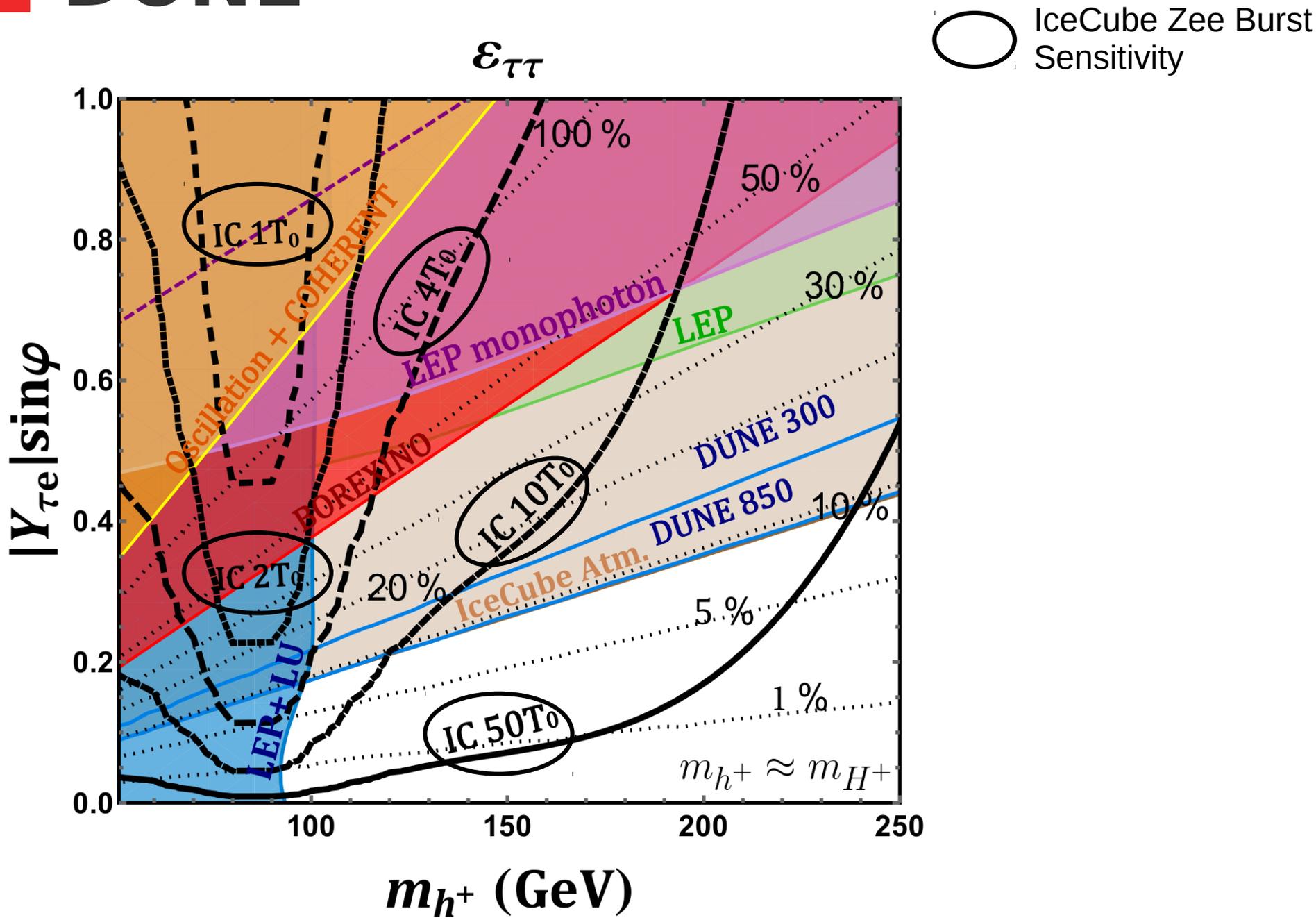
Sensitivity of IceCube and DUNE



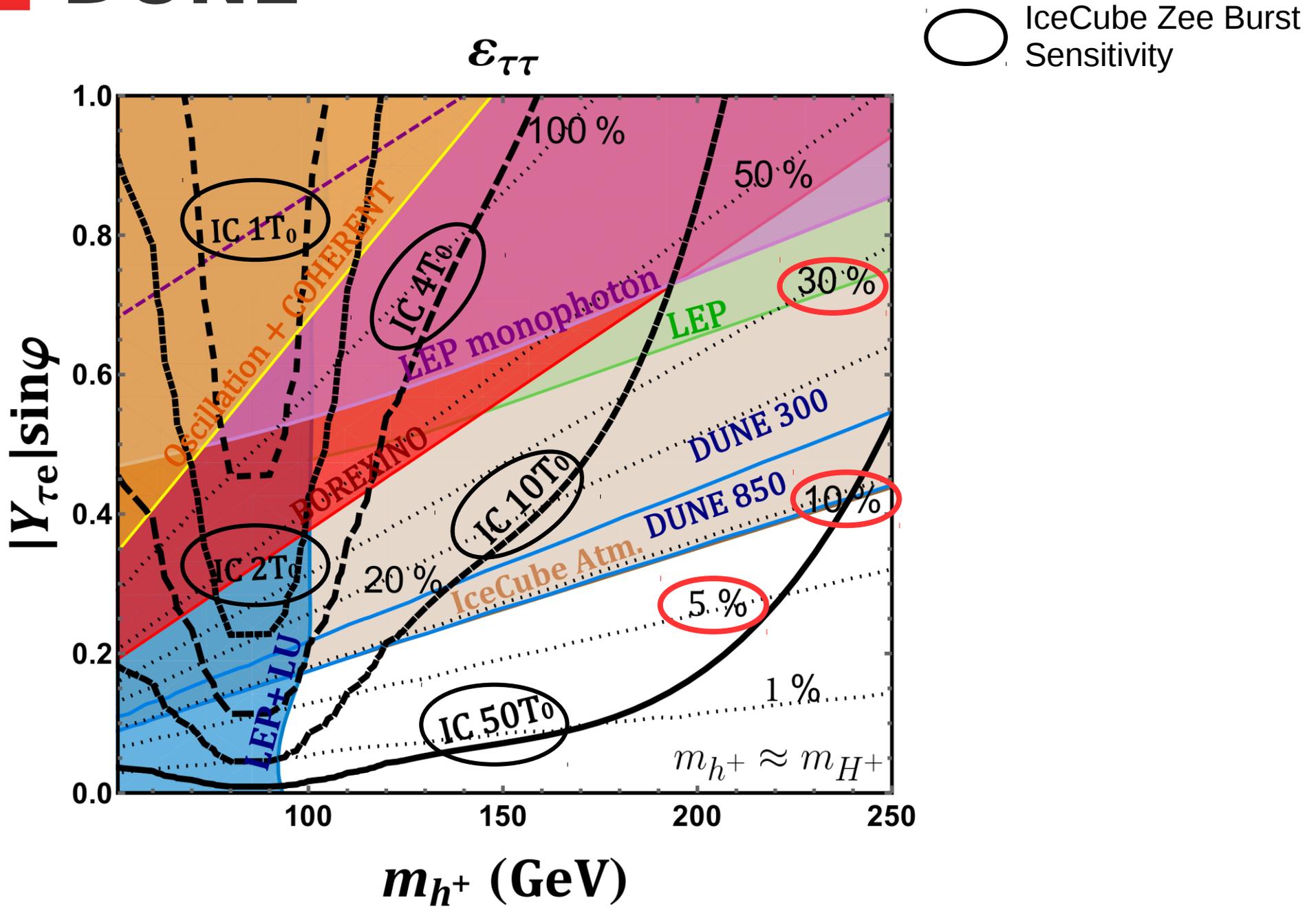
Sensitivity of IceCube and DUNE



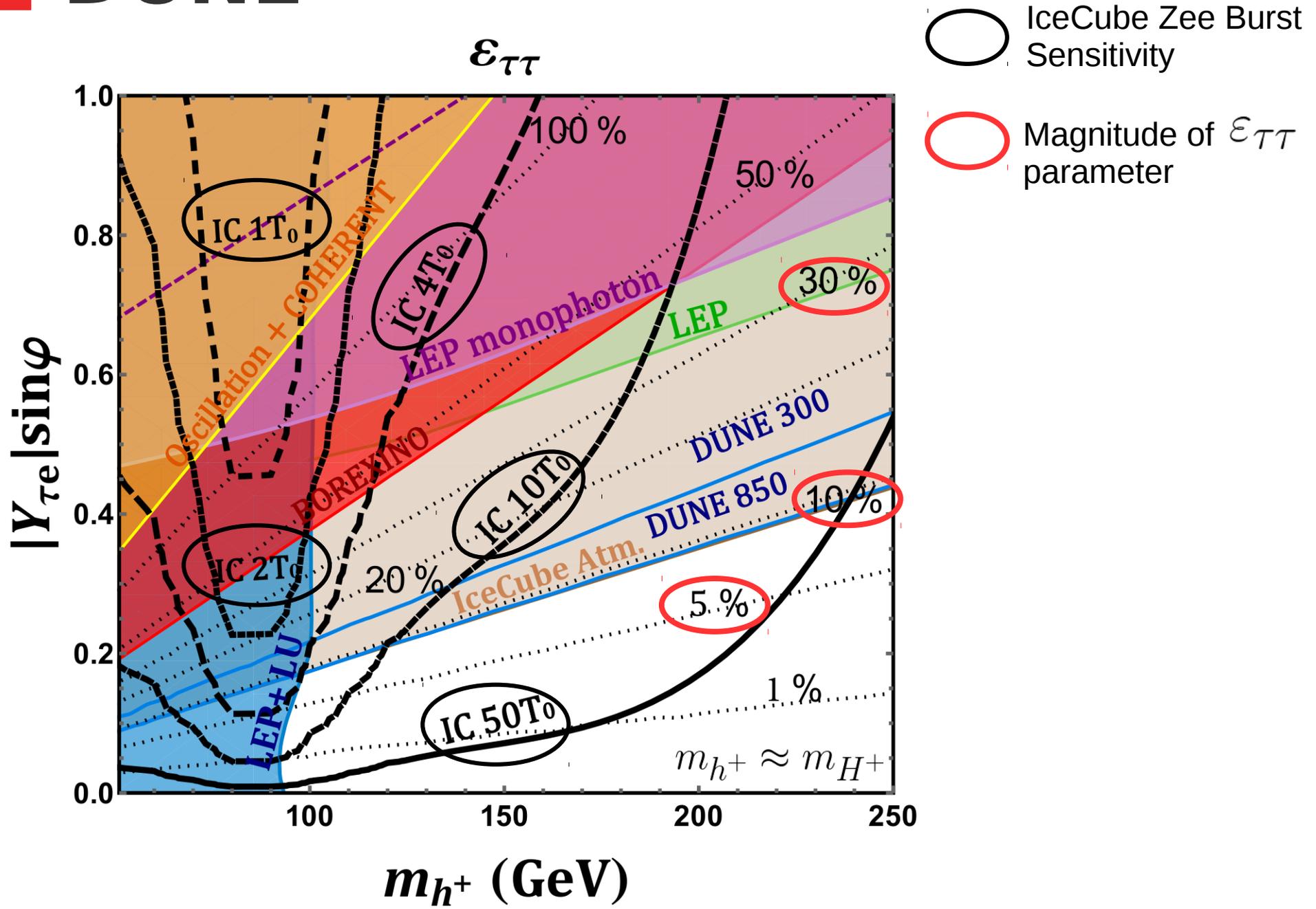
Sensitivity of IceCube and DUNE



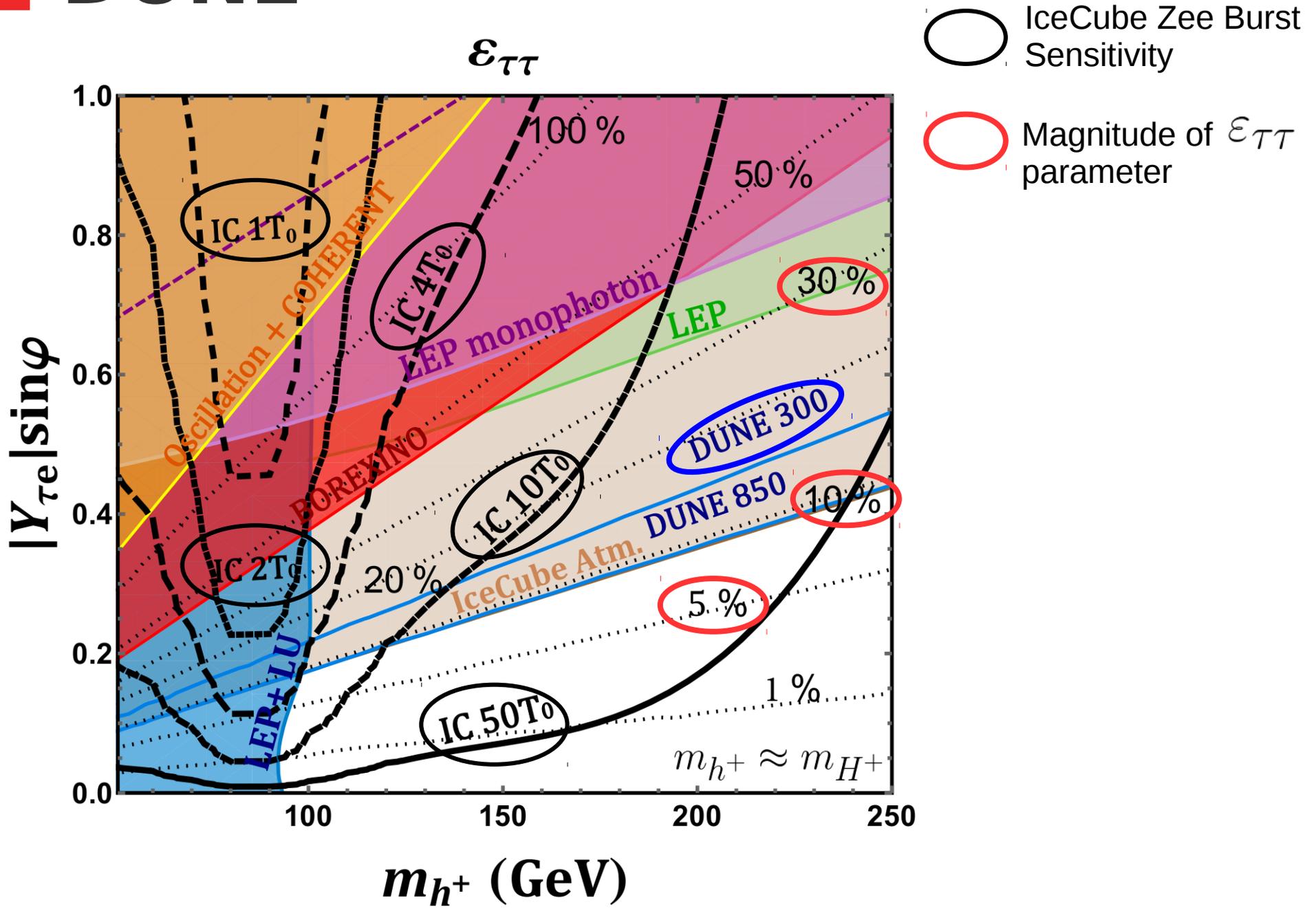
Sensitivity of IceCube and DUNE



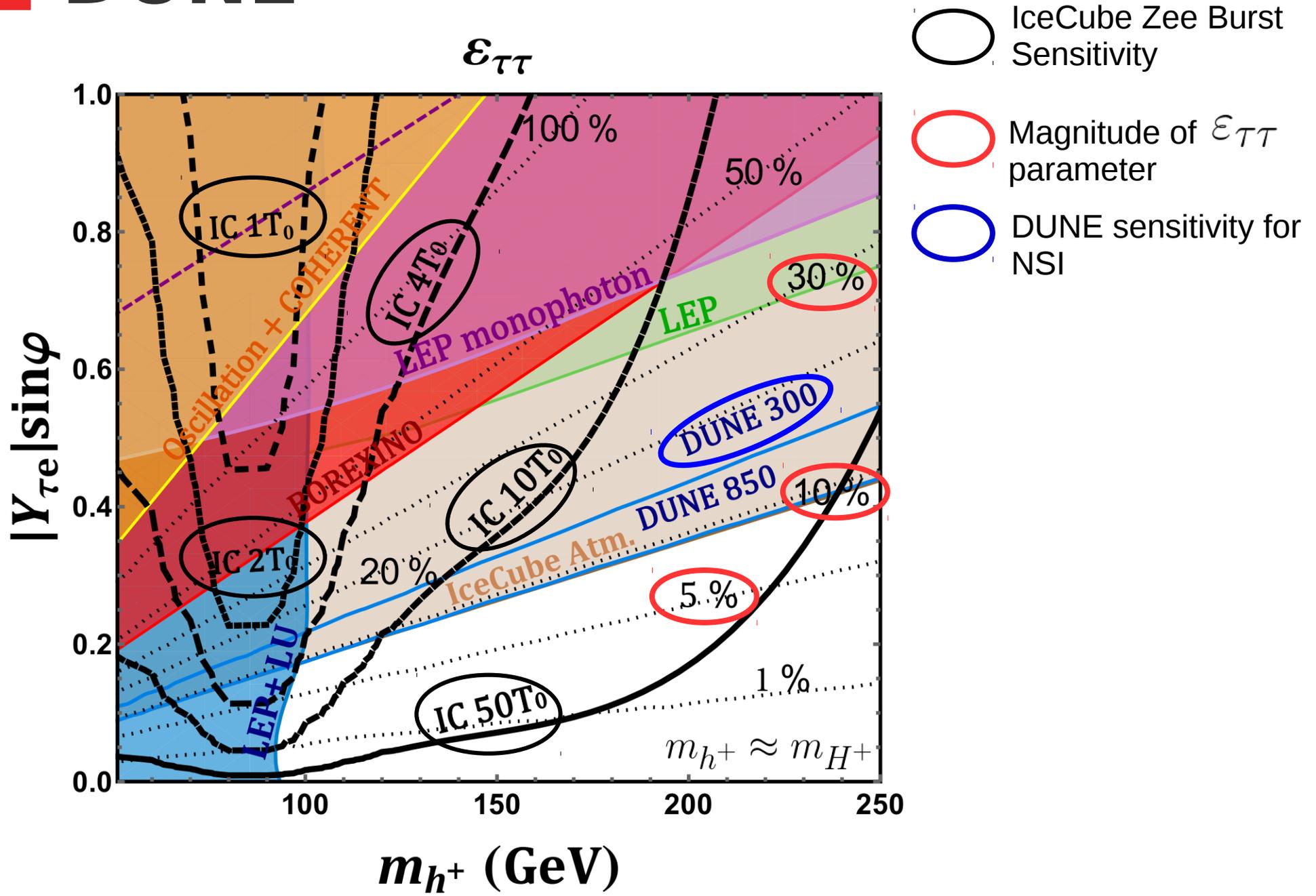
Sensitivity of IceCube and DUNE



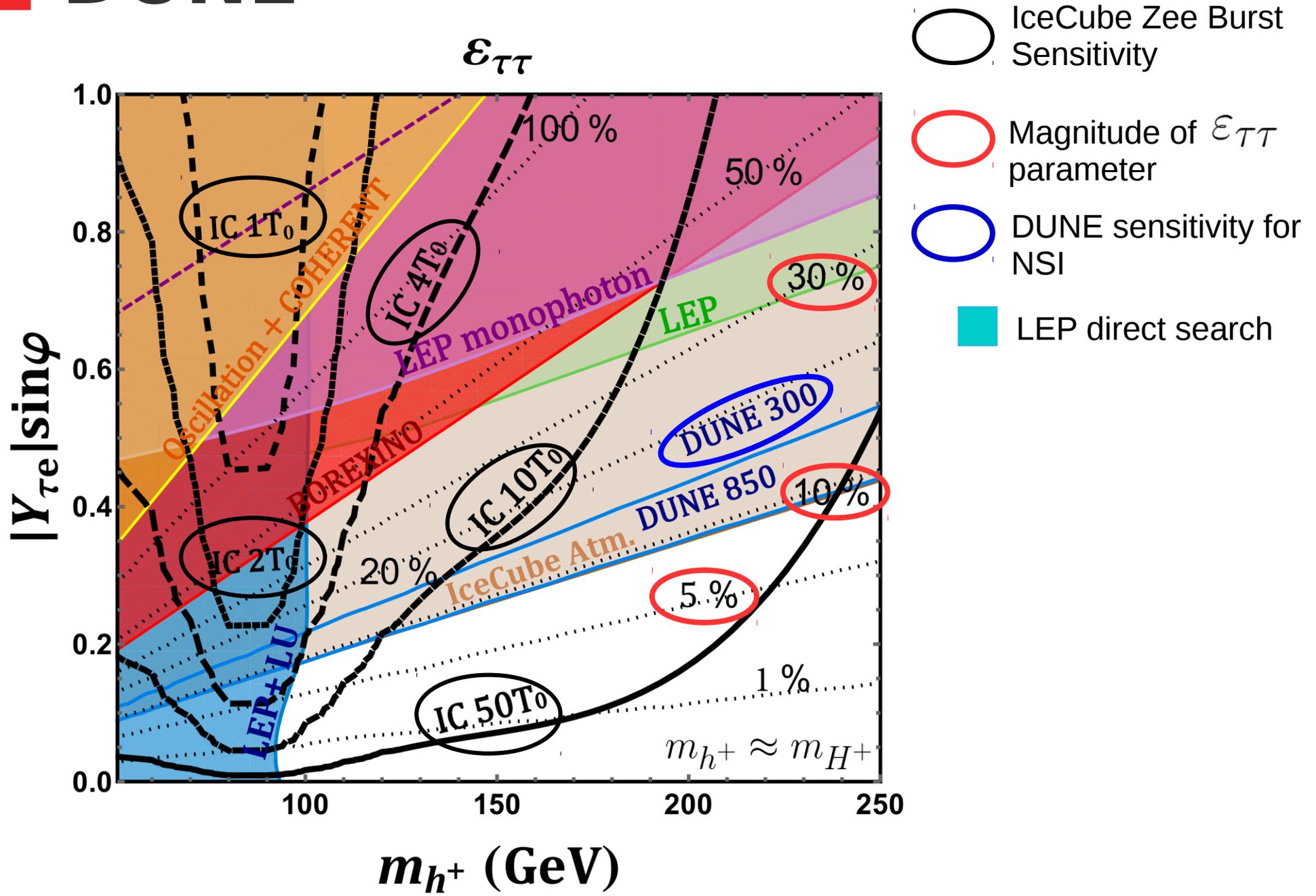
Sensitivity of IceCube and DUNE



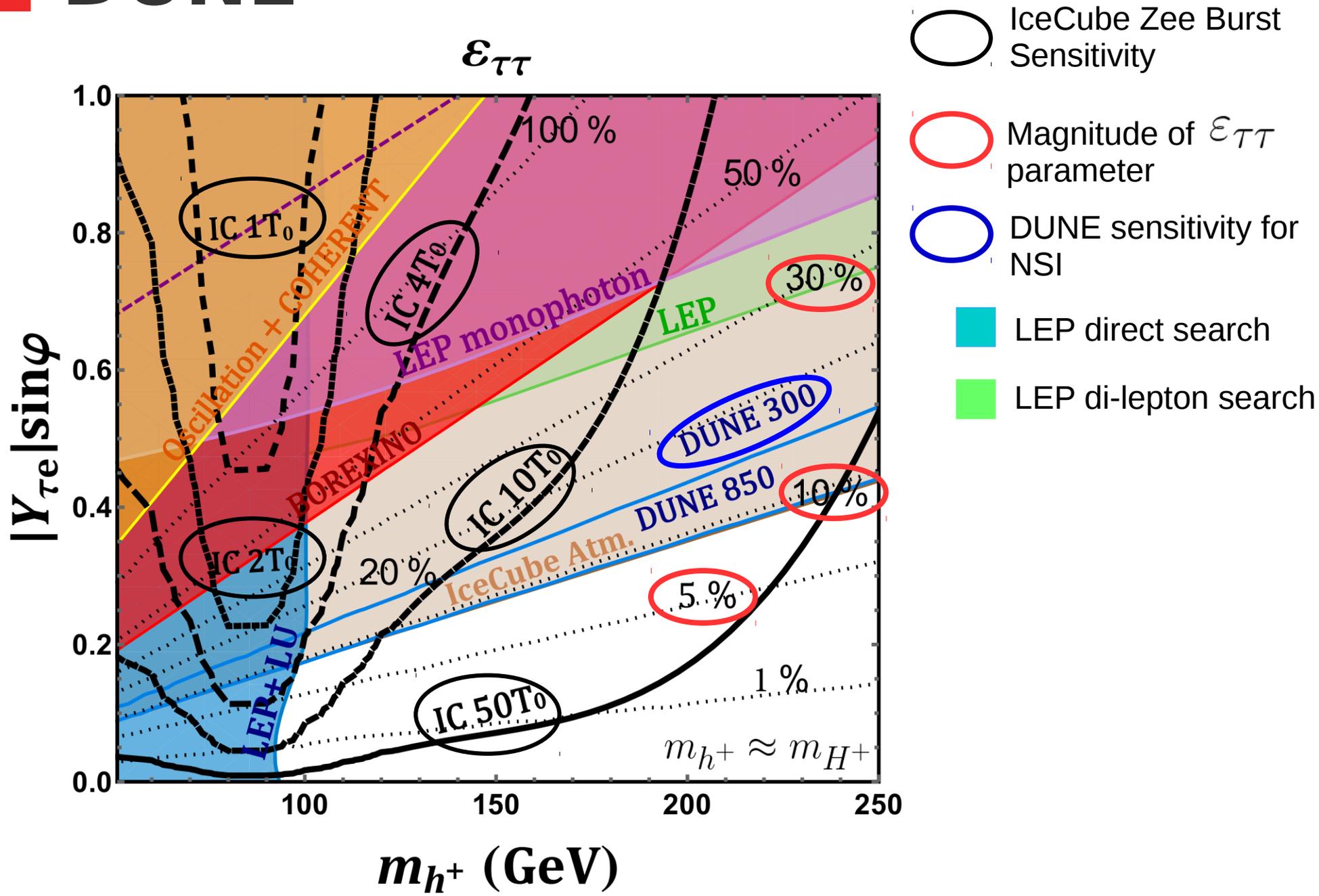
Sensitivity of IceCube and DUNE



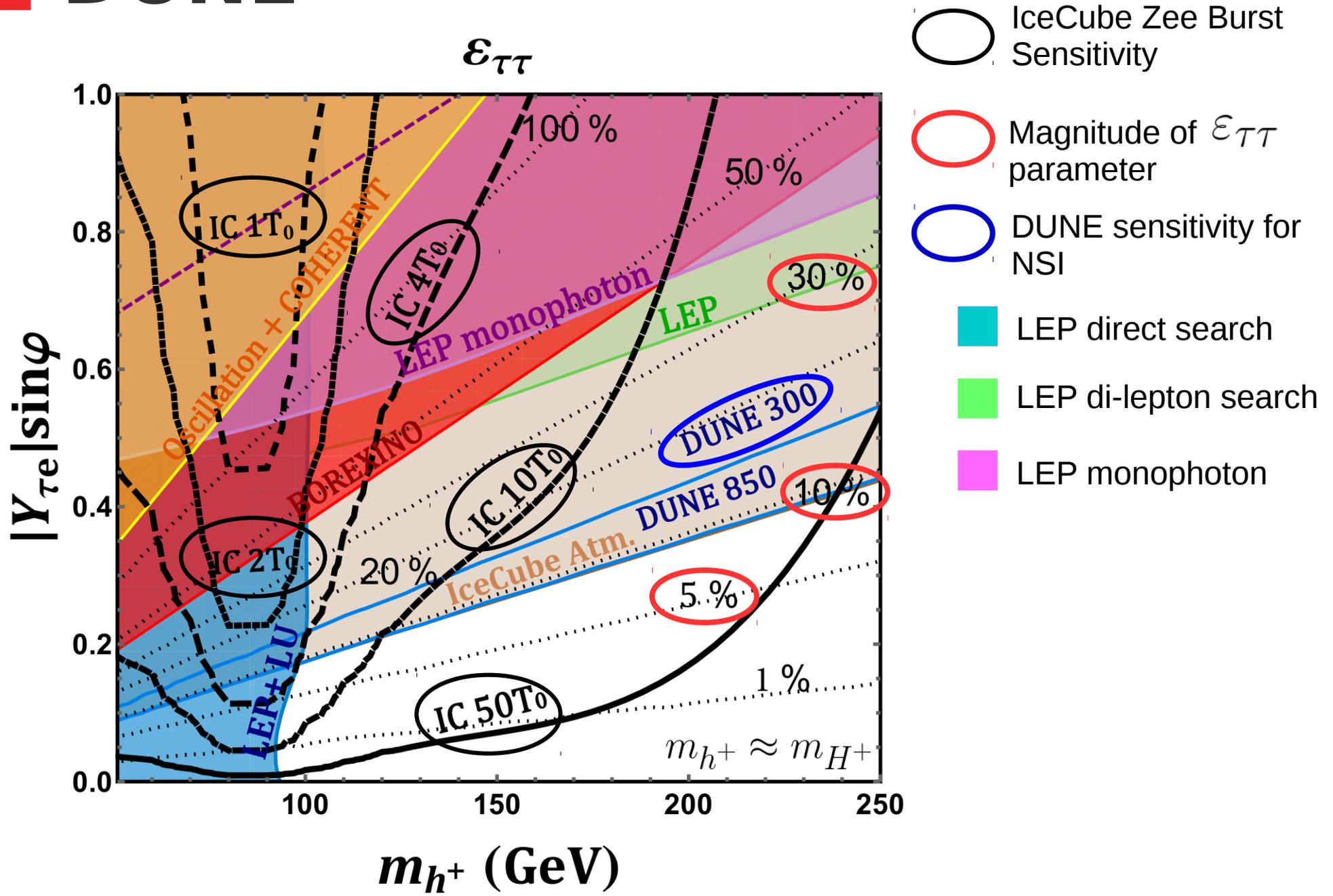
Sensitivity of IceCube and DUNE



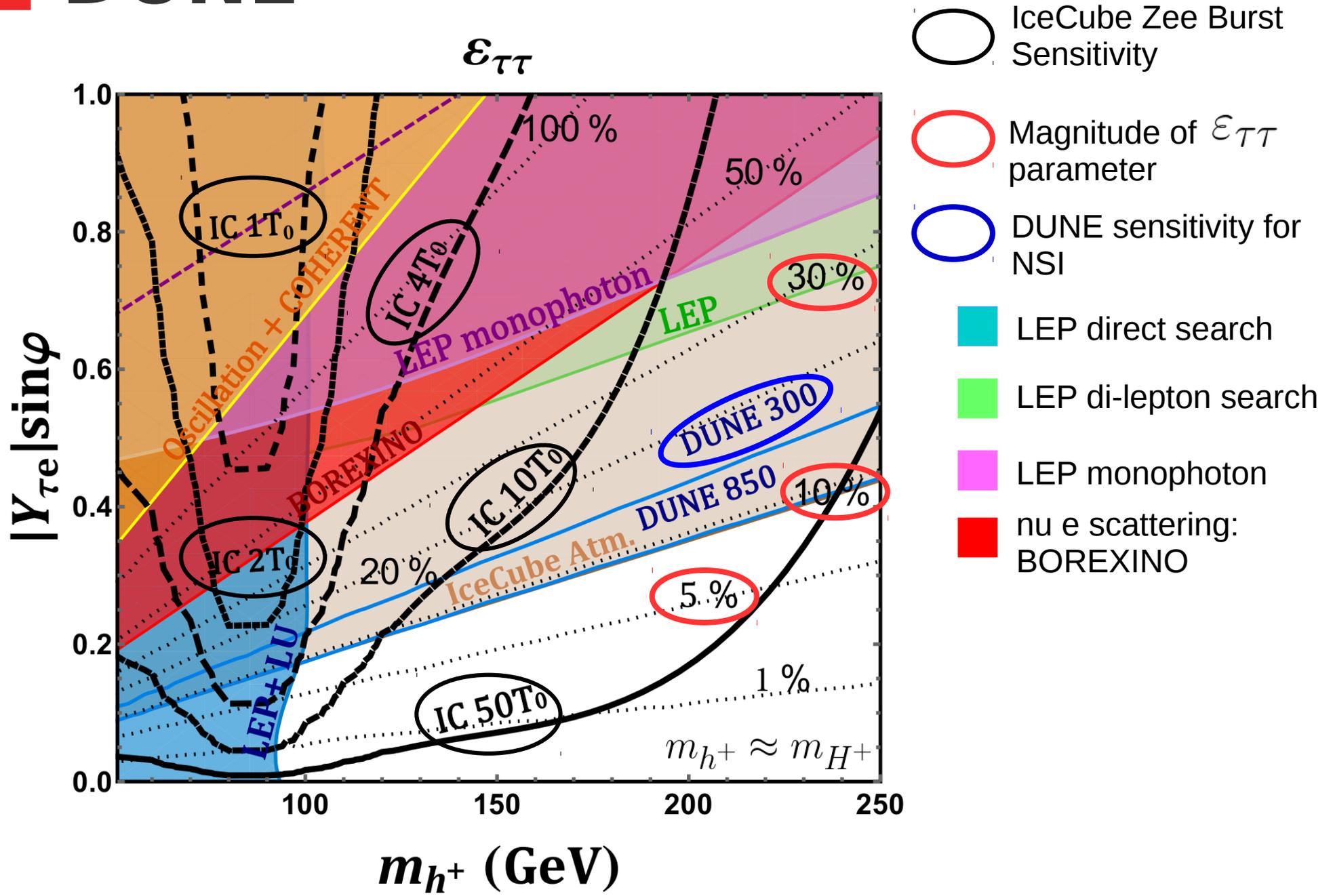
Sensitivity of IceCube and DUNE



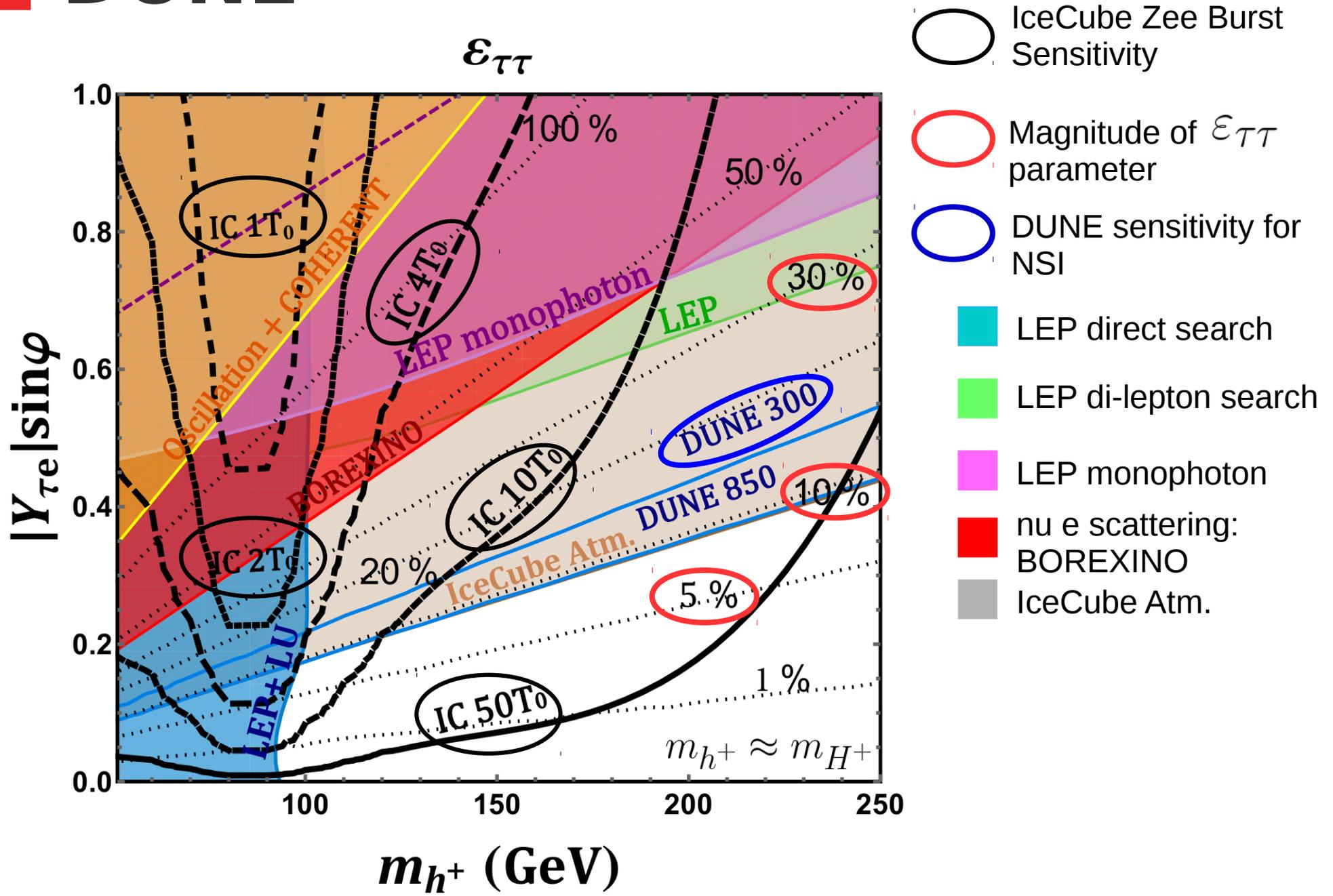
Sensitivity of IceCube and DUNE



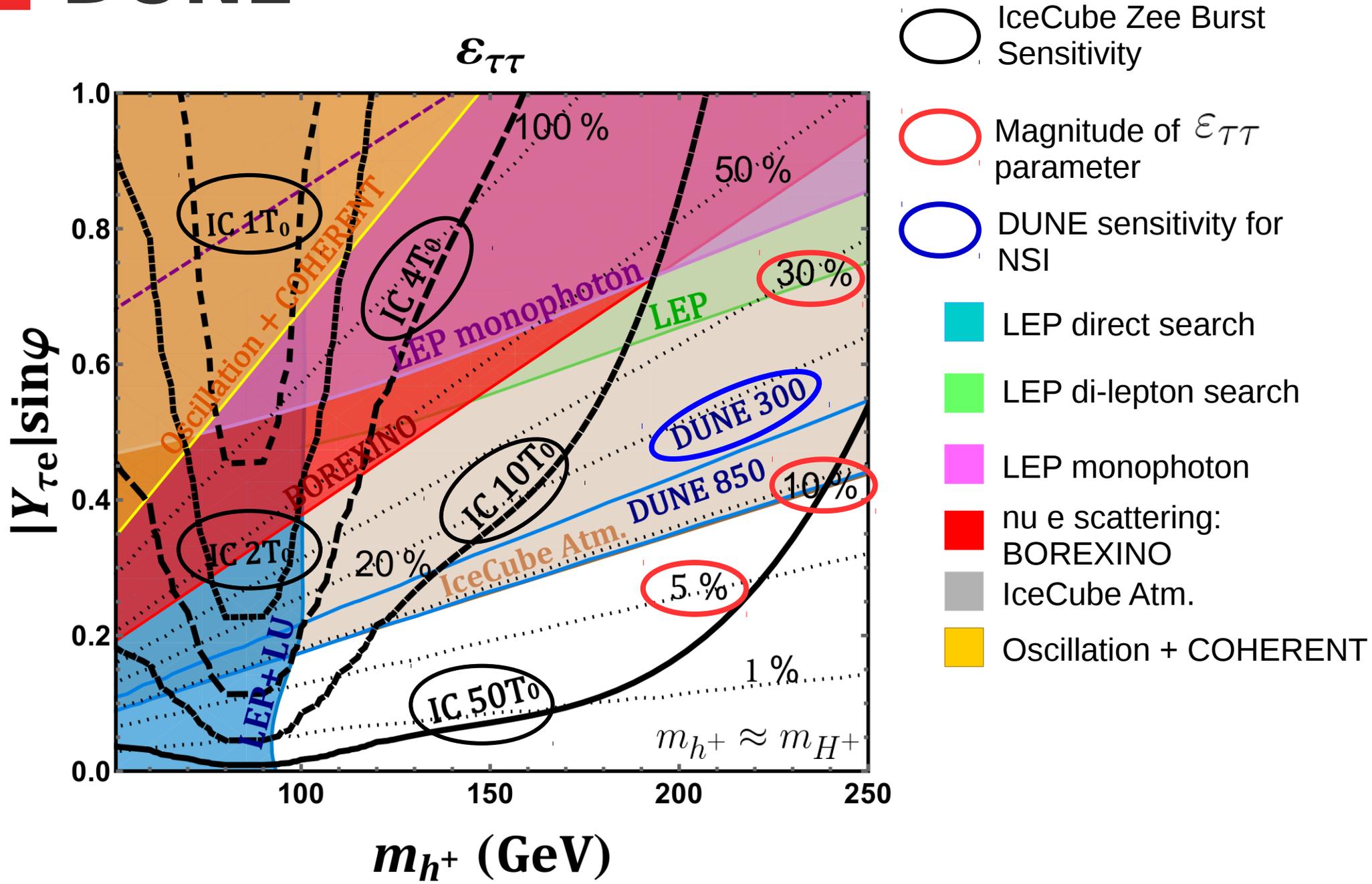
Sensitivity of IceCube and DUNE



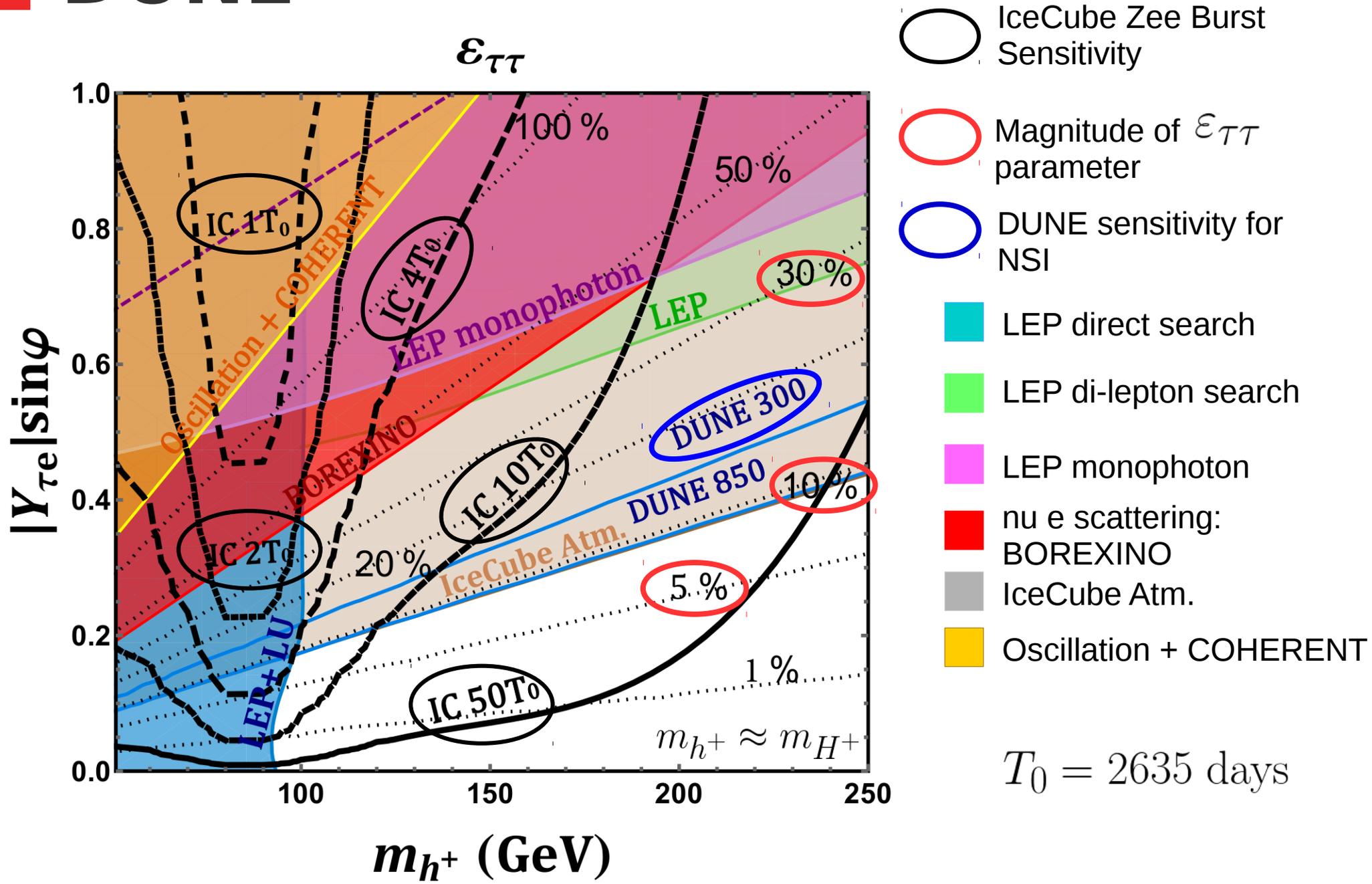
Sensitivity of IceCube and DUNE



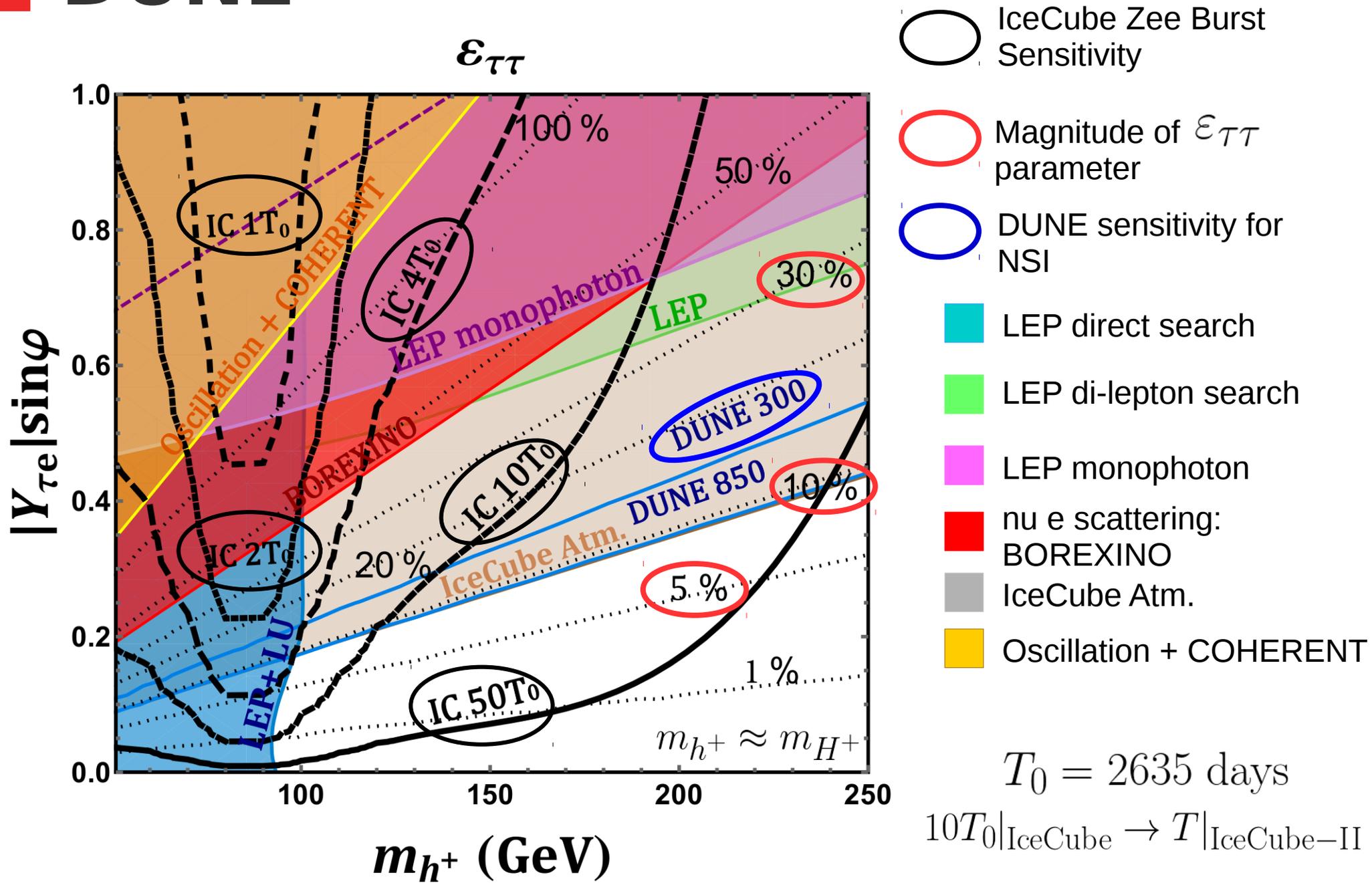
Sensitivity of IceCube and DUNE



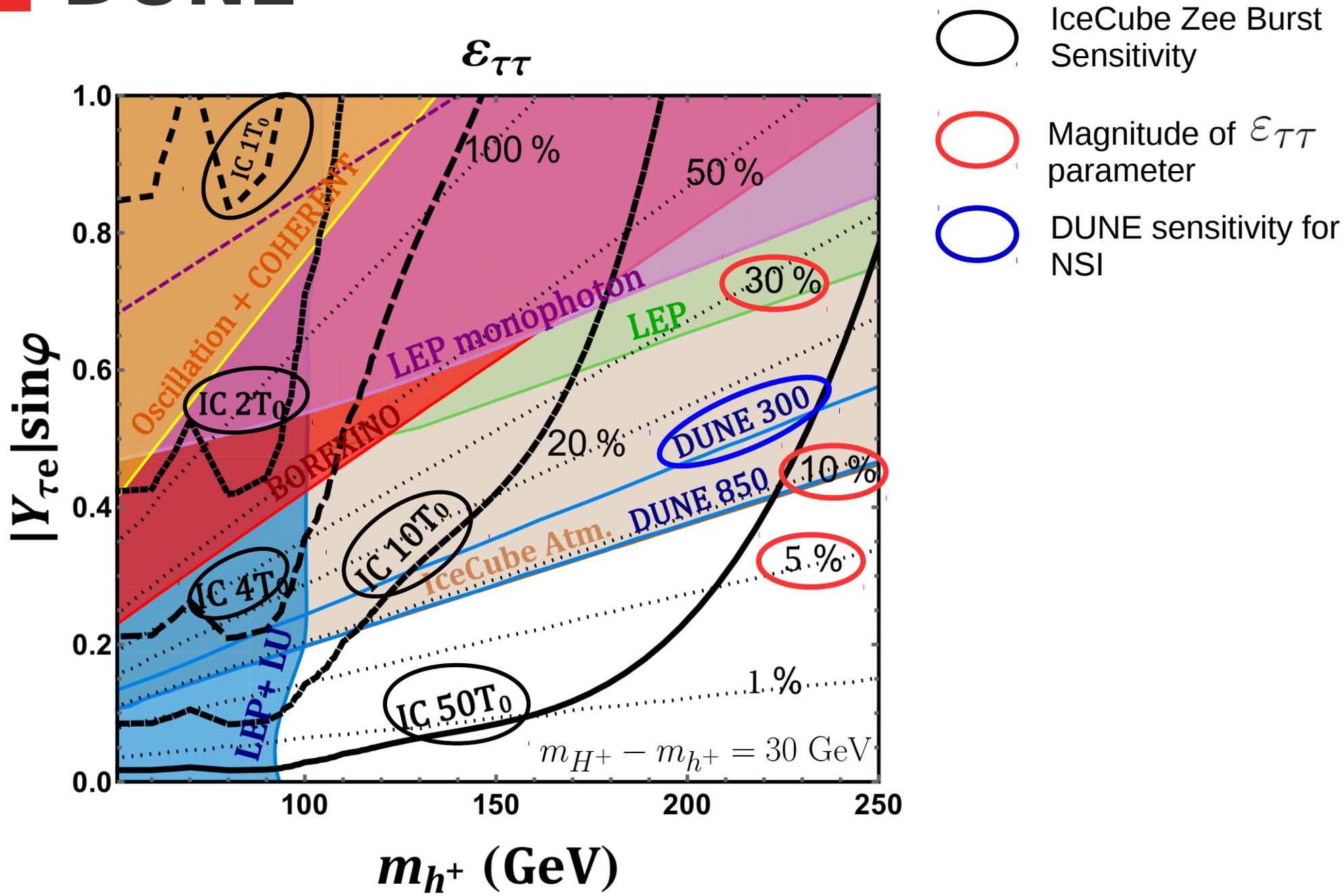
Sensitivity of IceCube and DUNE



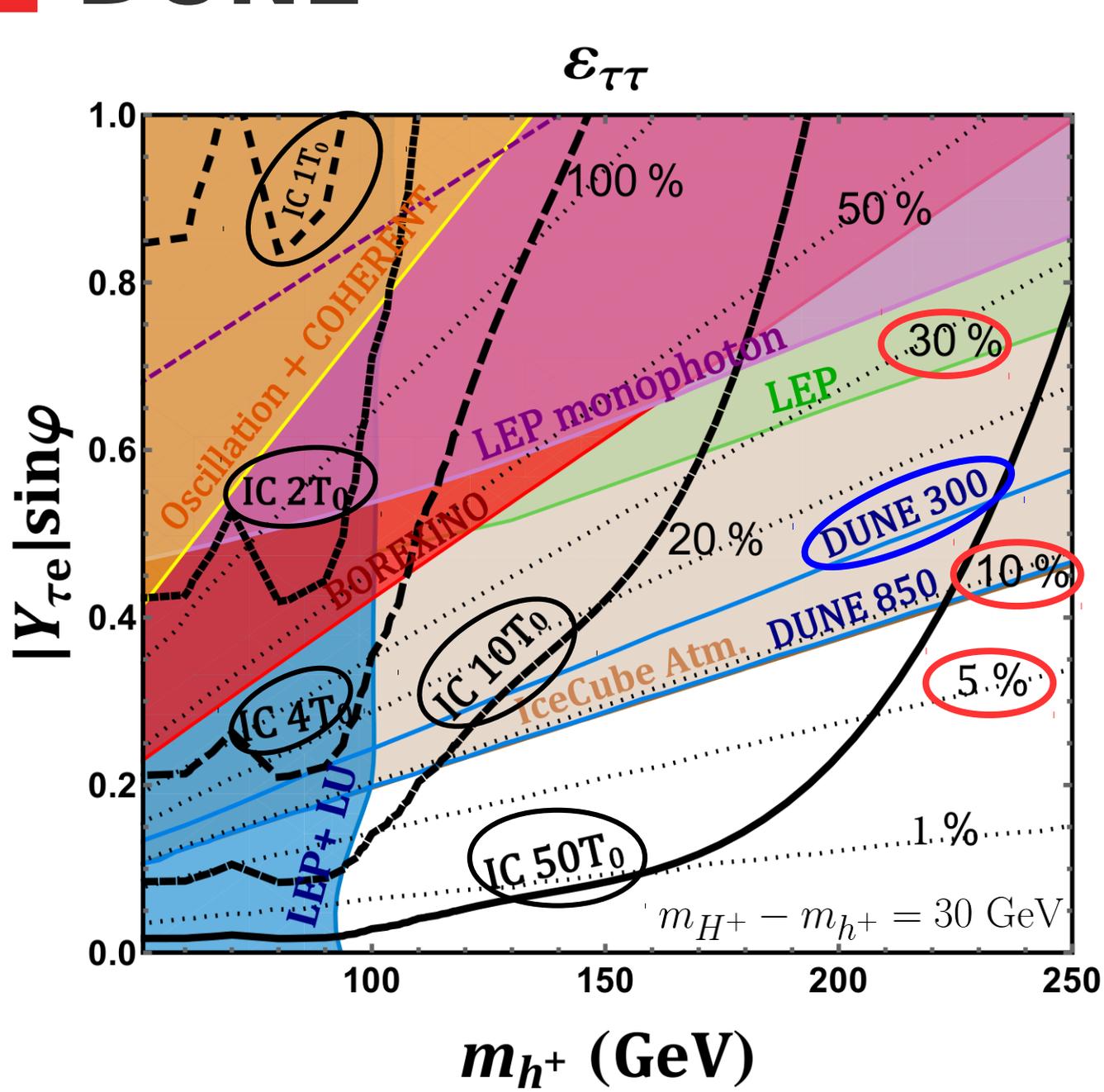
Sensitivity of IceCube and DUNE



Sensitivity of IceCube and DUNE



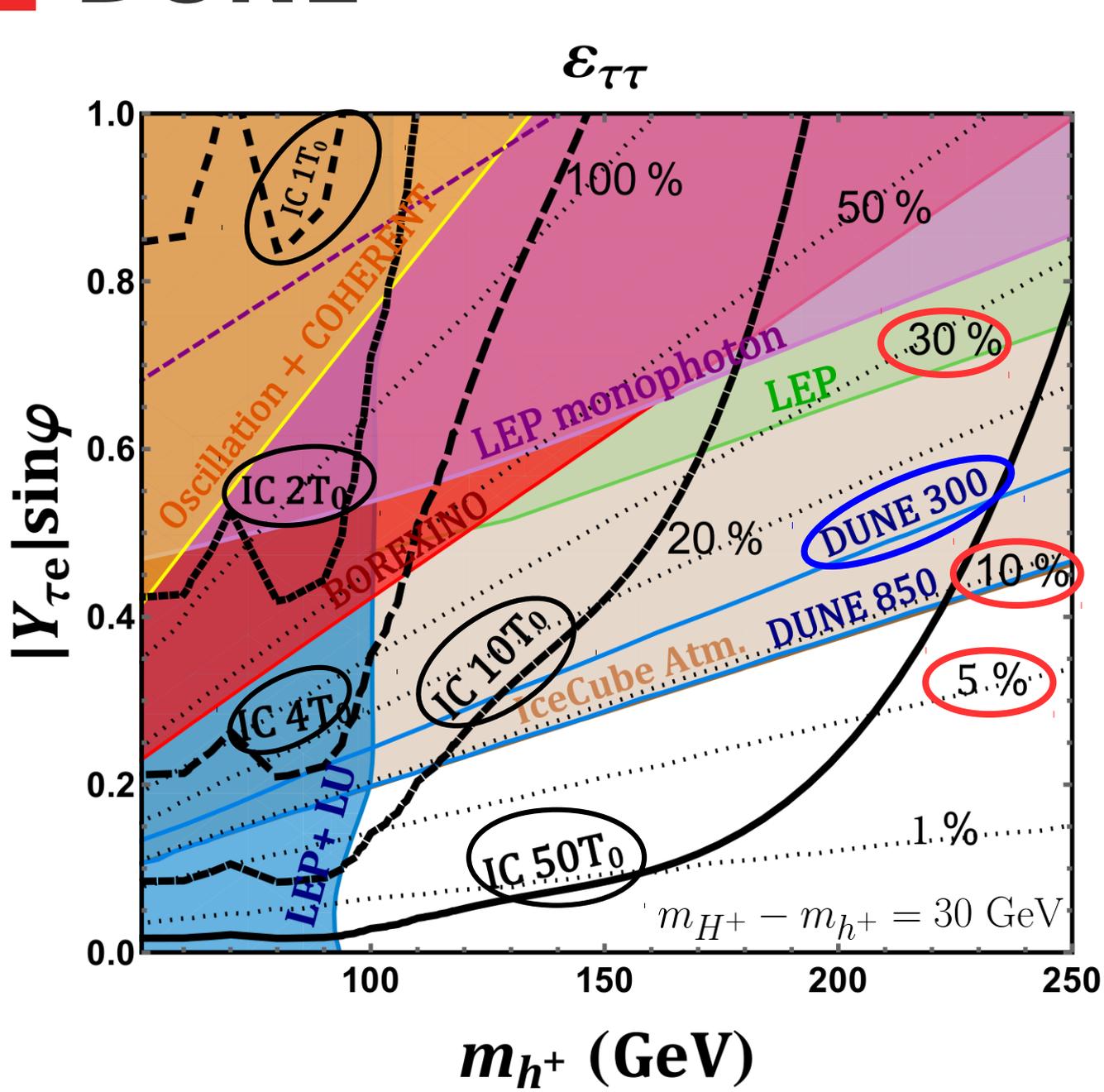
Sensitivity of IceCube and DUNE



- IceCube Zee Burst Sensitivity
- Magnitude of $\epsilon_{\tau\tau}$ parameter
- DUNE sensitivity for NSI

Double-dip feature is due to the double peak cross section feature:

Sensitivity of IceCube and DUNE



- IceCube Zee Burst Sensitivity
- Magnitude of $\epsilon_{\tau\tau}$ parameter
- DUNE sensitivity for NSI

Double-dip feature is due to the double peak cross section feature:

$$m_{H^+} - m_{h^+} = 30 \text{ GeV}$$



Conclusion

- We proposed a new way to probe light charged scalars using a Glashow-like resonance in the UHE neutrino data (IceCube).
- The same interactions for Glashow-like resonance also give rise to observable NSI effect.
- UHE neutrinos provide a complementary probe of NSI.
- We have used the popular Zee model of radiative neutrino mass as a demonstration.
- Further extensions to other models are possible and promising.



Thank you!